

How does the Einstein equation look like when formulated using the Vielbein?

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Consider the functional derivative of the Lagrange density:

$$\frac{\delta \mathcal{L}}{\delta V^\mu{}_\alpha} \quad , \quad \frac{\delta S}{\delta V^\mu{}_\alpha} = 0$$

Consider the parts that can be written in terms of the metric, like the Einstein-Hilbert term:

$$M_{\text{pl}}^2 \int d^4x \sqrt{g} R$$

Then:

$$\frac{\delta \mathcal{L}}{\delta V^\mu{}_\alpha} = \frac{\partial \mathcal{L}}{\partial g^{\lambda\kappa}} \frac{\partial g^{\lambda\kappa}}{\partial V^\mu{}_\alpha}$$

$$\begin{aligned} \frac{\partial g^{\lambda\kappa}}{\partial V^\mu{}_\alpha} &= \frac{\delta}{\delta V^\mu{}_\alpha} \left( V^\lambda{}_\rho V^\kappa{}_\sigma \eta^{\rho\sigma} \right) \\ &= \delta^\lambda{}_\mu \delta^\kappa{}_\alpha V^\rho{}_\nu \eta^{\rho\nu} \\ &\quad + V^\lambda{}_\beta \delta^\mu{}_\alpha \delta^\kappa{}_\nu \eta^{\beta\nu} \end{aligned}$$

And hence:

$$\frac{\partial g^{\lambda\kappa}}{\partial V^{\mu}_a} V_{\nu\alpha} = \delta^{\lambda}_{\mu} \delta^{\kappa}_{\nu} + \delta^{\lambda}_{\nu} \delta^{\kappa}_{\mu}$$

$$\frac{\partial \mathcal{L}}{\partial V^{\mu}_a} V_{\nu\alpha} = 2 \frac{\partial \mathcal{L}}{\partial g^{\alpha\mu}}$$



This enters in the  
energy-momentum  
tensor of GR

This will enter  
the generalization of the  
energy-momentum tensor  
when the vielbein is used

Word auf caution: It is not obvious that  
this construction leads to a symmetric  
energy-momentum tensor. But one can prove  
that it does.

## How to construct the covariant derivative for the vielbeins

We would like to construct a covariant derivative, such that

$$\mathcal{D}_\beta A_\alpha$$

transforms like a tensor

$$\mathcal{D}_\beta A_\alpha \rightarrow \Lambda_\rho^\delta \Lambda_\alpha^{\kappa\beta} \mathcal{D}_\delta A_\kappa$$

$$A_\alpha \rightarrow \Lambda(x)^\beta_\alpha A_\beta$$

The construction follows the same steps as in GR.

We can do the same approach for general representations, so consider:

$$\psi \rightarrow \mathcal{D}(\Lambda(x)) \psi$$

We would like to construct a covariant derivative such that

$$\nabla_\alpha \psi$$

transforms as

$$\nabla_\alpha \psi_n \rightarrow D(\Lambda)_{nm} \Lambda_\alpha^\beta (\nabla_\beta \psi_m)$$

$$\psi_n \rightarrow D(\Lambda)_{nm} \psi_m$$

Ultimately, one finds for the covariant derivatives:

$$\nabla_\alpha = \nabla_\alpha^\mu \left( \partial_\mu \mathbb{1}_{mn} + \frac{1}{2} \sigma_{mn}^{\beta\gamma} \nabla_\beta^\nu \left( \frac{\partial}{\partial x^\mu} \nabla_\nu \right) \right)$$

Using this we finally write down the Lagrangian for a spinor

$$\mathcal{L} \ni \bar{\psi} (\not{\partial} - m) \psi$$

↳

This can now be used to determine:

- EoM of the spinor
- a contribution to the energy-momentum in the Einstein equation

## Penrose diagrams:

Penrose diagrams are digrams that display the causal structure of a space. They make it easy to track the paths followed by light.

For example: Minkowski space:

$$ds^2 = - dt^2 + dx^2 \quad (1D)$$

$$= - du dv$$

$$u = t - x$$

$$v = t + x$$

$$(u, v) \in (-\infty, \infty)$$

$$p = \arctan v$$

$$q = \arctan u$$

$$(p, q) \in \left(-\frac{\pi}{2}, \frac{\pi}{2}\right)$$

$$= - \frac{dp dq}{\cos^2 p \cos^2 q}$$

$$= - \frac{dT^2 + dX^2}{\cos^2\left(\frac{T+X}{2}\right) \cos^2\left(\frac{T-X}{2}\right)}$$

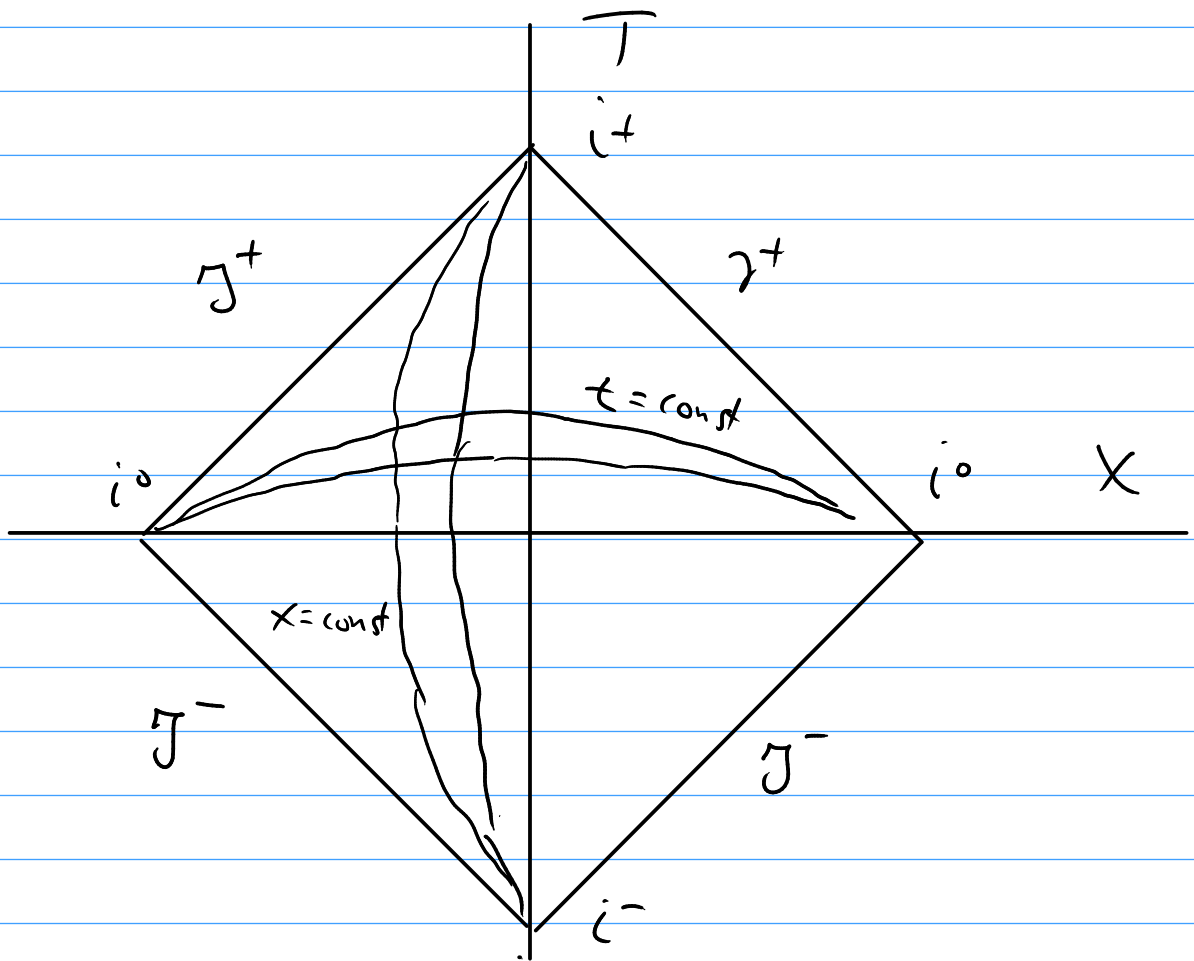
$$\equiv \Omega(x, \tau) (-d\tau^2 + dx^2)$$

$$T = p + q$$

$$X = p - q$$

$$T+X \in [-\pi, \pi]$$

$$T-X \in [-\pi, \pi]$$



Light follows:

$$ds^2 = 0$$

$$dT = dx$$

↳ diagonal lines

- $i^o$  space-like infinity
- $i^+$  time-like + infinity
- $i^-$  time-like - infinity
- $J^+$  light rays end here
- $J^-$  light rays originate here

Another example: deSitter space  
(cosmology with a CC)

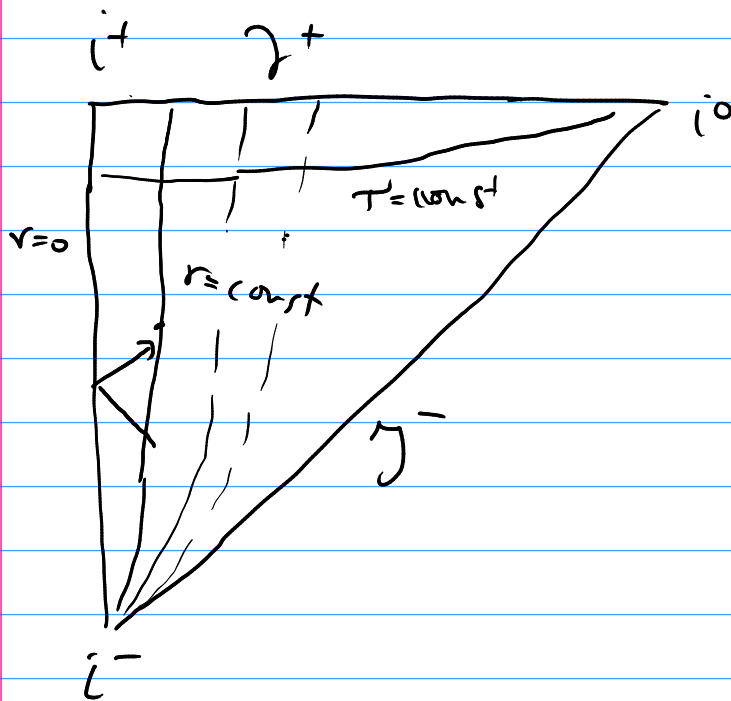
$$a(t) = e^{Ht} \quad ; \quad H = \text{const.}$$

The metric in 3D looks like

$$ds^2 = -dt^2 + a(t)^2 (dr^2 + r^2 d\Omega^2)$$

(let's neglect.)

$$= a(t)^2 (-dT^2 + dr^2)$$



conformal time:

$$a(t)dT = dt$$

$$dT = e^{-Ht} dt$$

$$T = -\frac{1}{a \cdot H}$$

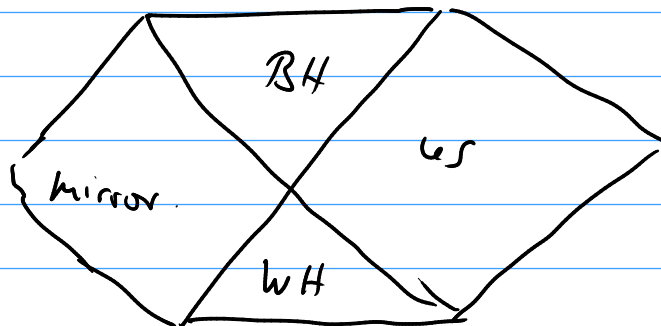
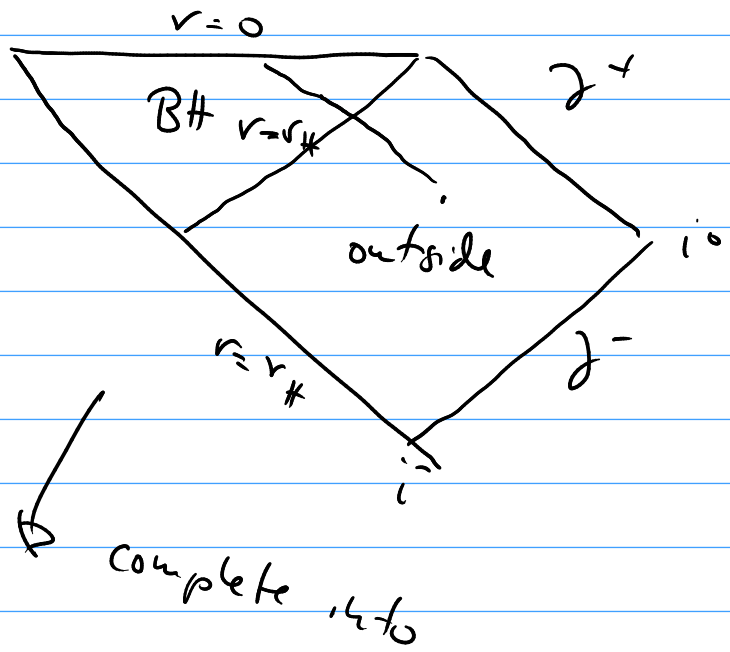
$$a \in [0, \infty]$$

$$r \in [-\infty, 0]$$

$$r \in [0, \infty]$$

Notice that there are horizons.

black holes:



collapsing star:

