

Irreducible representations of the rotation group

We are interested in the representations of the rotation group with the algebra

$$[\partial_i, \partial_j] = i \epsilon_{ijk} \partial_k \quad i, j, k \in \{1, 2, 3\}$$

Since the operators do not commute, we can diagonalize at most one of the them.

However, consider

$$\partial^2 = \sum_i \partial_i^2$$

Then

$$\begin{aligned} [\partial^2, \partial_i] &= \sum_k [\partial_k^2, \partial_i] \\ &= \sum_k \partial_k [\partial_k, \partial_i] + [\partial_k, \partial_i] \partial_k \\ &= \sum_k i \epsilon_{ijk} \{ \partial_k \partial_j + \partial_j \partial_k \} \\ &= 0 \end{aligned}$$

Hence we can diagonalize

$$\begin{aligned} \mathcal{J}^2 \text{ and } \mathcal{J}_z \\ \mathcal{J}^2 |\alpha, \beta\rangle &= \alpha |\alpha, \beta\rangle \\ \mathcal{J}_z |\alpha, \beta\rangle &= \beta |\alpha, \beta\rangle \end{aligned}$$

Now consider the ladder operators

$$\begin{aligned} \mathcal{J}_+ &= \mathcal{J}_x + i\mathcal{J}_y \\ \mathcal{J}_- &= \mathcal{J}_x - i\mathcal{J}_y \end{aligned}$$

Then

$$\begin{aligned} [\mathcal{J}_z, \mathcal{J}_+] &= [\mathcal{J}_x, \mathcal{J}_x] + i[\mathcal{J}_x, \mathcal{J}_y] \\ &= i\mathcal{J}_y + \mathcal{J}_x = \mathcal{J}_+ \end{aligned}$$

also: $[\mathcal{J}_x, \mathcal{J}_-] = -\mathcal{J}_-$

and

$$[\mathcal{J}_+, \mathcal{J}_-] = 2\mathcal{J}_z$$

and obviously

$$[\mathcal{J}_\pm, \mathcal{J}^2] = 0$$

What are the quantum numbers of

$$J_+ |\alpha, \beta\rangle \quad \text{and} \quad J_- |\alpha, \beta\rangle ?$$

$$\begin{aligned} J^2 J_{\pm} |\alpha, \beta\rangle &= J_{\pm} J^2 |\alpha, \beta\rangle \\ &= \alpha J_{\pm} |\alpha, \beta\rangle \end{aligned}$$

But

$$\begin{aligned} J_+ J_+ |\alpha, \beta\rangle &= J_+ + J_+ J_+ |\alpha, \beta\rangle \\ &= (\beta+1) J_+ |\alpha, \beta\rangle \end{aligned}$$

Hence

$$J_+ |\alpha, \beta\rangle = N |\alpha, \beta+1\rangle$$

Likewise

$$J_- |\alpha, \beta\rangle = N |\alpha, \beta-1\rangle$$

Could $N=0$ at some point and terminate the series?

Now consider the norm of $\mathcal{J}_+ |\alpha, \beta\rangle$,

The hermitian conjugate of \mathcal{J}_+

$$(\mathcal{J}_+)^{\dagger} = \mathcal{J}_-$$

yields

$$\langle \alpha, \beta | \mathcal{J}_- \mathcal{J}_+ | \alpha, \beta \rangle \geq 0$$

and

$$\langle \alpha, \beta | \mathcal{J}_+ \mathcal{J}_- | \alpha, \beta \rangle \geq 0$$

Hence also

$$\langle \alpha, \beta | \frac{1}{2} (\mathcal{J}_+ \mathcal{J}_- + \mathcal{J}_- \mathcal{J}_+) | \alpha, \beta \rangle \geq 0$$

$$\frac{1}{2} (\mathcal{J}_+ \mathcal{J}_- + \mathcal{J}_- \mathcal{J}_+) = \mathcal{J}_x^2 + \mathcal{J}_y^2 = \mathcal{J}^2 - \mathcal{J}_z^2$$

$$\hookrightarrow \alpha - \beta^2 \geq 0$$

Hence the chain of eigenstates has to terminate in both directions and there are minimal and maximal values of β .

Using

$$J^+ J^- = J^L - J^2 + J_z$$

and

$$J^- J^+ = J^L - J^2 - J_z$$

One obtains for the minimal β

$$0 = \langle \alpha, \beta_{\min} | J^+ J^- | \alpha, \beta_{\min} \rangle$$

$$\hookrightarrow \alpha - \beta_{\min}^2 + \beta_{\min} = 0$$

And likewise

$$\alpha - \beta_{\max}^2 - \beta_{\max} = 0$$

Accordingly one finds

$$-\beta_{\min} = +\beta_{\max} = \alpha$$

$$\hookrightarrow J^2 |l, m\rangle = l(l+1) |l, m\rangle$$

$$J_z |l, m\rangle = m |l, m\rangle$$

l has to be half-integer and positive and m in $[-l, l]$.

Examples:

$$l = \frac{1}{2} \quad \therefore \quad J_i = \frac{1}{2} \sigma_i$$

with the Pauli matrices

$$l = 1 :$$

$$J_z = \begin{pmatrix} 0 & -i \\ i & 0 \\ & & 0 \end{pmatrix}$$

$$J_x = \begin{pmatrix} 0 & & \\ & 0 & i \\ & -i & 0 \end{pmatrix}$$

$$J_y = \begin{pmatrix} 0 & & i \\ & 0 & \\ -i & & 0 \end{pmatrix}$$

As an example of a reducible representation, consider a rank 2 tensor.

This contains representations with $l = 2, 1, 0$.

Total number of degrees of freedom is

$$5 + 3 + 1 = 9$$