

SLD 2001

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PHYSICAL REVIEW LETTERS

12 February 2001

Improved Direct Measurement of Leptonic Coupling Asymmetries with Polarized Z Bosons

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(The SLD Collaboration)

 $\sin^2\Theta_{\rm w} = 0.23098 \pm 0.00026$



POLARIZATION Highlights

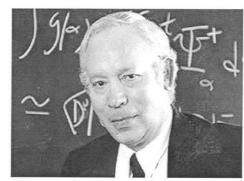
- The PRE-SLD Era
- The SLC
- Looking Ahead The Future Linear Collider



A Model of Leptons

Steven Weinberg - 1967

The first to unify the weak and electromagnetic forces, in 1967 Weinberg wrote down the most general form for an interaction for leptons, which included the concept of mixing, mass generation, and couplings to a heavy neutral gauge boson, the Z.



Steven Weinberg

The model assigned the electron and its neutrino to a left-handed doublet, while the right-handed electron was alone as a singlet. The neutral coupling that resulted was

Thus the left-handed and right-handed couplings were different.

This choice preserved the purely left-handed charged currents.



Polarized Electrons come to SLAC

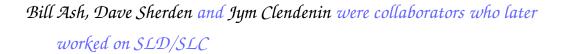
Fixed Target Experiment E80 proposed 1970

Vernon Hughes and collaborators propose a polarized beam/polarized target experiment to validate the quark model of the proton

"the SLAC - Yale experiment E80"

The polarized electrons were from a ⁶Li atomic beam ionized by a UV flash lamp

E80 was the first of a highly successful Spin Structure program at SLAC, CERN, DESY and elsewhere





Vernon W. Hughes



End Station A Experiments

In 1972 **E95** was proposed to look for parity violation in inelastic scattering – using the SLAC-YALE "PEGGY" source

The source was too low in intensity, and the proposal stated E95 was "insensitive to the weak interactions".

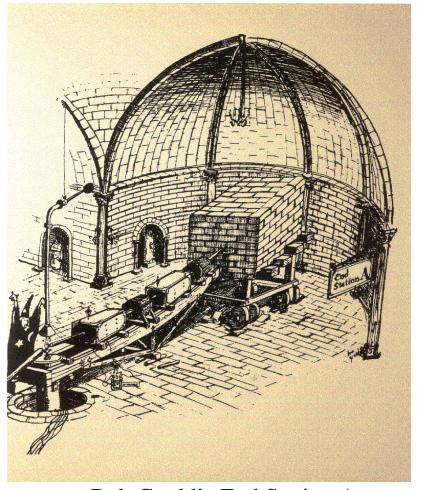
E95 published a limit in 1978

 A_{LB} < 2 x 10⁻³ at Q^2 = 1.2 GeV/ c^2

While E95 was underway, Charlie Sinclair and I were discussing ways to reach the weak level, as

defined in the Weinberg-Salam model. We needed an intense polarized electron source, and considered developing a laser-driven Fano source using a cesium vapor.

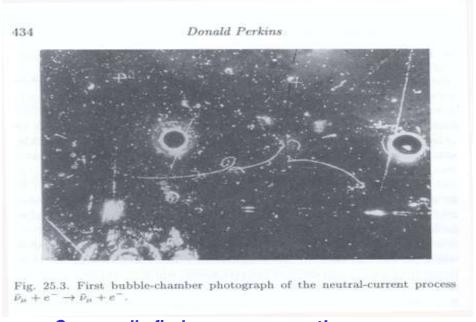
We proposed a new experiment, £122, which was approved, but we never built the Fano source.



Bob Gould's End Station A



Neutral Currents Discovered! Gargamelle CERN - 1973





(two more by 1976)

First Z⁰ seen in UA1 in 1983



Charlie Baltay

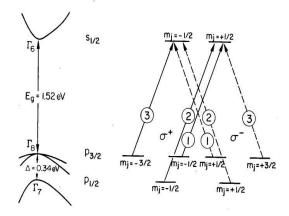


The Promise of Gallium Arsenide 1974

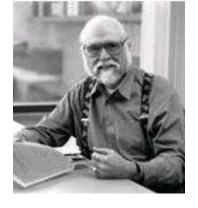
Gallium Arsenide was well known to have polarized internal electrons when optically pumped by circularly polarized light (Ekimov and Sakarov, JETP Letters 13, 495 (1971))

Bell and Spicer had shown that the conduction band electrons could be photoemitted by adding Cs-O monolayers to the surface.

Ed Garwin knew of these works and the need for a source at SLAC.







Ed Garwin

Bill Spicer



Gallium Arsenide proposed

Garwin, Pierce, and Siegmann
1974

Ed Garwin visited ETH Zurich in 1974, and while there proposed to develop a polarized electron source using gallium arsenide. The first source was built and demonstrated by Dan Pierce at ETH Zurich (now at NIST).

The density of electrons in GaAs is high, promising large available currents. GaAs as a source of polarized electrons appeared ideal for SLAC, but first, the principles had to be demonstrated.



H. C. Siegmann



Dan Pierce



E. L. Garwin



Parity Violation

1974-1978

The prospect of a GaAs photoemission source for high beam currents triggered a new proposal....which could test the Weinberg-Salam model in the End Station. Charlie Sinclair and I proposed such a test to the SLAC EPAC in 1974.

This experiment was E122.

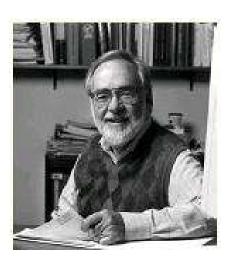
The proposal received conditional approval and we went to work on the source, with Ed Garwin and Roger Miller. That occupied us for 4 years.



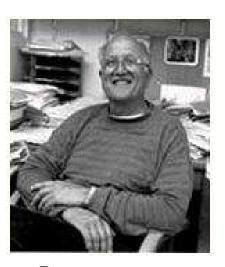
Charlie Sinclair



Ed Garwin



Charlie Prescott



Roger Miller



1977 – The Drama Intensifies Atomic Parity Violation lays an EGG

Two competing atomic physics groups eagerly pursue parity violation in bismuth vapor – Washington and Oxford

They hold noisy debates in conferences and reports – among themselves on the one hand, and with Weinberg and Salam on the other. They argued for the "hybrid" model which predicted no parity violation. They published back_to-back_null results in PRL.

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Zh. Eksp. Teor. Fiz. 71, 1665 (1976) [Sov. Phys. JETP (to be published)].

⁹I. P. Grant, N. C. Pyper, and P. G. H. Sandors, to be published.

³S. Weinberg, Phys. Rev. Lett. 19, 1264 (1967).
⁸A. Salam, in Proceedings of the Eighth Notel Symposium, edited by Svartholm (Almkvist and Wiksell, Stockholm, 1968).

³We use the optical convention that a positive rotation appears clockwise when looking toward the source, ¹⁵M. A. Bouchiat and C. C. Bouchiat, Phys. Lett. <u>48B</u>,

111 (1974).

11 Collisional broadening becomes noticeable for He

buffer gas pressures above 100 Torr, but no observable collisional enhancement of the integrated absorption of this M. line occurs.

12A convenient parameter is the mean number of ab-

A convenient parameter is the mean number of absorption lengths of the his components at their peaks, 15The central dip associated with the Faraday effect disappears in the average over the transmitted laser

light for conditions of strong absorption as in Fig. 2(b). The average over the laser profile of any λ -dependent background rotation will change when the absorp-

tion line alters the transmitted laser profile.

19 P. E. G. Baird et al., follwing Letter |Phys. Rev. Lett. 39, 798 (1977)].

Search for Parity-Nonconserving Optical Rotation in Atomic Bismuth

P. E. G. Baird, M. W. S. M. Brimicombe, R. G. Hunt, G. J. Roberts, P. G. H. Sandars, and D. N. Stacey Clarendon Laboratory, University of Oxford, Oxford, England (Received 7 July 1977)

We report the results of a laser experiment to search for the parity-nonconserving optical rotation in anomic bismuth. We work at wavelengths close to the 89^{4} -may = 8/2-M, transition from the ground state. We find $R = \ln(R/M) = (+2.1 - 4.1 \times 10^{-3})$ in disagreement with the theoretical value $R = -30 \times 10^{-4}$ predicted for this transition on the basis of the Weitherg-Salam model of the weak interactions combined with relativistic central-field atomic theory.

We report the results of an experiment to search for the parity-nonconserving (PNC) optical rotation¹⁻⁹, in atomic bismuth which has been predicted⁸⁻⁷ on the basis of the Weinberg-Salam⁸.

ing the different approaches employed.

Our apparatus is illustrated schematically in

Fig. 1. The Spectra-Physics 580A jet-stream dye
laser produces approximately 2 mW of light in a

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⁵R. E. Tribble, J. D. Cossairt, and R. A. Kenefick, Phys. Rev. C <u>15</u>, 2028 (1977).

⁶J. C. Harcy, J. E. Esterl, R. G. Sextro, and J. Cerny, Phys. Rev. C 3, 700 (1971). ⁵P. M. Endt and C. van der Leun, Nucl. Phys. <u>A214</u>,

1 (1973).

R. L. McGrath, J. Cerny, J. C. Hardy, G. Goth,

*R. L. McGrath, J. Cerny, J. C. Hardy, G. Goth, and A. Arima, Phys. Rev. C J., 184 (1870).

*J. C. Hardy and I. S. Towner, Nucl. Phys. A<u>254</u>, 221

notte, C.R. Acad, Sci. <u>271B</u>, 970 (1870).

(1975).

⁰A. H. Wapstra and K. Bos, At. Data Nucl. Data Tables <u>19</u>, 175 (1977).

¹E. G. Adelberger and D. P. Balamuth, Phys. Rev. Lett. <u>27</u>, 1597 (1971).

Fortier, H. Laurent, J. M. Maison, J. P. Schapira, and J. Vernotte, Phys. Rev. C 6, 378 (1972).
 S. Galès, M. Langevin, J. M. Maison, and J. Vernotte, C.R. Acad. Sci. 271B, 970 (1970).

Upper Limit on Parity-Nonconserving Optical Rotation in Atomic Bismuth

L. L. Lewis, J. H. Hollister, D. C. Soreide, E. G. Lindahl, and E. N. Fortson Department of Fhysics, University of Washington, Scattle, Washington 98195 (Received 7 July 1977)

We have searched for optical rotation near the 8737- $\tilde{\Lambda}$ magnetic-dipole absorption line in atomic bitsmuth vapor. The experiment is sensitive to parity nonconservation in the weak neutral-current interaction between electrons and nucleons in atoms. We find $R = \operatorname{Im}(E_q/M_s) = (-0.7 + 3.2 \times 10^{-4})$, which is considerably smaller than the value $R = -2.5 \times 10^{-7}$ obtained by central-field calculations for this bismuth line using the Weinberg-Salam theory of neutral currents.

We present here results of an experiment in which we search for parity-nonconserving (PNC) optical rotation in atomic bismuth vapor. 142 We

selected the $J=\frac{3}{2}+J=\frac{3}{2}$ absorption line at 8757 Å where there is no competing background absorption from Bi, molecular bands to limit the usable

At SLAC, the laser-driven GaAs source works; Polarized electrons are accelerated in December 1977.



E122 Announces Parity Violation June 1978

In June 1978, in the SLAC Auditorium, E122 announced the evidence for parity violation in inelastic ep and ed scattering. The statistical significance exceeded 10 sigma. Consistency checks and null texts were fully satisfied.

In the Fall, E122 again ran, and the combined data agreed with the W-S Model and gave

$$sin^2 \Theta_w = 0.224 \pm .020$$

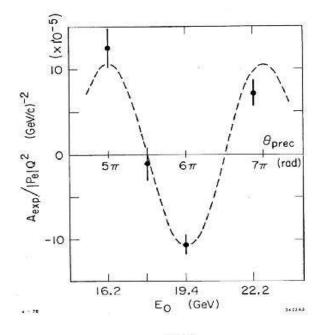


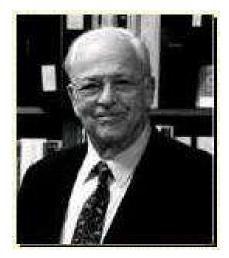
Fig. 3

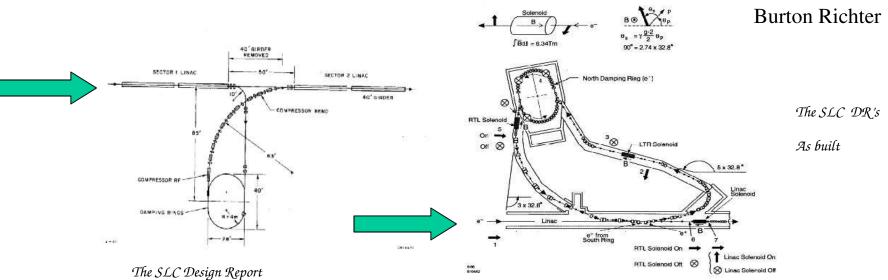


The SLC Era

By 1979, the ESA Parity Violation work was winding down. Richard Taylor and I were discussing the future program, and I argued for joining Group C's plans led by Burton. He was planning the SLC. I wanted to take polarized electrons to the SLC. Dick called a meeting in early 1980, to talk about the idea. Burton listened.

Burton also passed along to me an invitation to speak about this in Europe that summer.





June 1980



HEP Spin Physics Conference, Lausanne 1980 13th Erice Summer School, 1980

Burton received invitations to speak in 1980. He had two sent down to me, the Erice Summer School and the Lausanne HEP Spin Conference. I spoke on the physics of polarized electrons at the Z-pole.

Brian Montague, an accelerator theorist at CERN, also spoke at the Lausanne Conference. He worked in the LEP Machine group, and had a long standing interest in polarized electrons. He summarized his talk with the cartoon at the right.

Polarization at LEP was a difficult technical challenge. They eventually abandoned all serious efforts, leaving that physics to the SLC. where $\Gamma_{\rm f}$ is the partial width into final state f. Figure 7 shows the ratio of

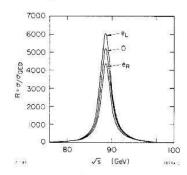


Fig. 7. $R = \sigma/\sigma_{QED}$ versus center-of-mass energy near the Z^O pole for electron beams of positive helicity (e_R) , negative helicity (e_L) , or unpolarized beams (0).

Figure / shows the ratio of σ/σ_{QED} versus \sqrt{s} , the center-of-mass energy, near the Z^0 mass. The cross section σ_{QED} is defined as $\sigma(e^-e^+ + \mu^- \mu^+)$ from single γ exchange only. In Fig. 7, single γ exchange has been included, but is negligible near the Z^0 mass. The curves are marked e_L , 0, e_R for $P_e = -1, 0 + 1$, respectively. The experimental asymmetry

$$A_{\text{exp}} = (\sigma_{\text{R}} - \sigma_{\text{L}})/(\sigma_{\text{R}} + \sigma_{\text{L}})$$
 (14)

where $\sigma_R(\sigma_L)$ is the total cross section for right-handed (left-handed) electrons, has the value -.16 at the Z^0 peak, for $\sin^2 \theta_w = .23$. In terms of the couplings

$$A_{\text{exp}} = \frac{2g_{\text{V}}^2 g_{\text{A}}^2}{g_{\text{V}}^2 + g_{\text{A}}^2} . \tag{15}$$

Measurement of the asymmetry, Eq. (14), and









Energy first, then Luminosity, then Polarization

The Heroes
Nan Phinney
Marc Ross
Pantaleo
Tracy Usher
Morris Swartz
Mike Woods
Takashi Maruyama
Paul Emma
Bob Kirby
Tim Barklow















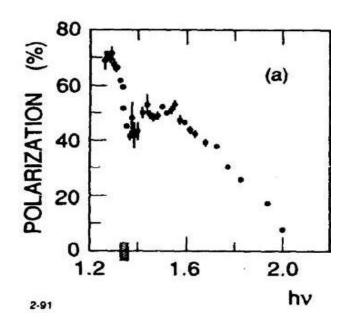


SLD Collaboration Meeting St. Francis Yacht Club October 5, 2001



Strained GaAs 1991

A Collaboration between SLAC-Wisconsin-Berkeley achieved a major breakthrough in 1991. Through the application of MBE techniques, they demonstrated the first strained Gallium Arsenide cathode, which yielded 70% polarization at the bandgap edge.





Takashi Maruyama



Ed Garwin



Dick Prepost



Cathode Charge Limit

SLAC-PUB-6268 April 1993 (A)

MEASUREMENT OF CHARGE LIMIT IN A STRAINED LATTICE GRAS PHOTOCATHODE*

P. Sáez, R. Alley, H. Aoyagi, J. Clendenin, J. Frisch, C. Garden, E. Hoyt, R. Kirby, L. Klaisner, A. Kulikov, C. Prescott, D. Schultz, H. Tang, J. Turner, K. Witte, M. Woods and M. Zolotorev

Stanford Linear Accelerator Center, Stanford University, Stanford, California 94309

ABSTRACT

The SLAC Linear Collider (SLC) Polarized Electron Source (PES) photocathodes have shown a charge saturation when illuminated with a high intensity laser pulse. This charge limit in the cesium-activated GaAs crystal seems to be strongly dependent on its surface condition and on the incident light wavelength. Charge limit studies with highly polarized strained lattice GaAs materials are presented.

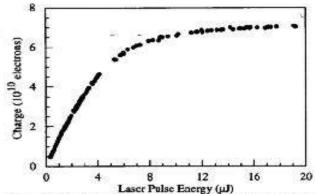


Figure 3. Typical photoemitted charge in the charge limit regime as a function of Ti:sapphire (850 nm) laser pulse energy. The photocathode QE was 1.51% and 0.57% at 750 nm and 833 nm, respectively. The cathode voltage was V = 120 kV.



Mike Woods



February 1992 LEP turns on, and the SLC is in trouble

The polarized electron guns are now in deep trouble. They don't hold voltage, and no one knows why. Burton is under tremendous pressure to get the SLC running, so turns to Sid Drell to help figure out what to do for polarization. Sid talks to many SLACers. He decides to set up a task force and asks me to join. At first I refuse, but then agree. He asks Lowell Klaisner to join, and Bob Kirby. My job was the gun test lab... to figure out what is wrong. Lowell's job was to coordinate the Tech division support. Bob's was to build the

load-lock system.

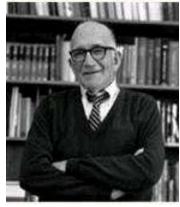
We barely survived the 1992 turn on in May.



Charles Prescott



Bob Kirby



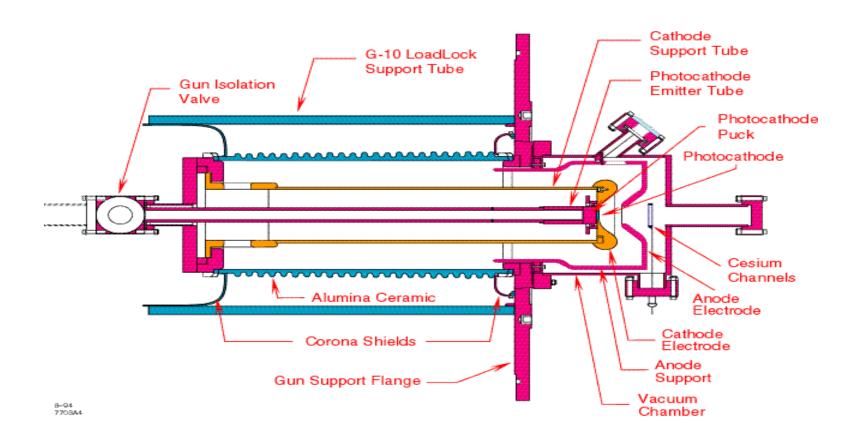
Sid Drell



Lowell Klaisner



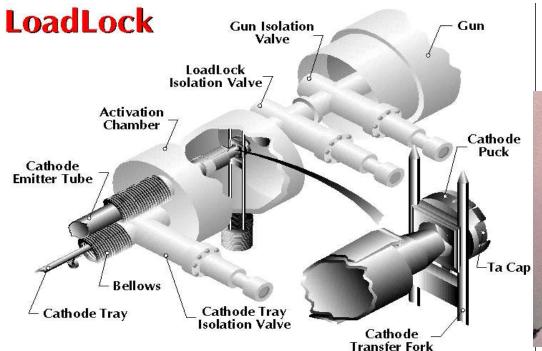
SLC Polarized Gun





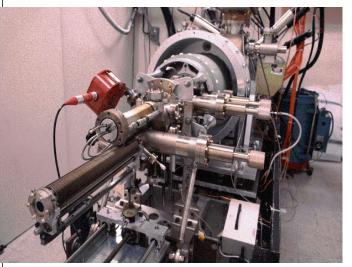
The Need for Load-Lock

The Gun structure could not hold voltage.....unless "processed", which contaminated the clean cathode surface. The fix was to install the cleaned cathode through a valve system into the processed gun... a "load-lock" device, which had to be invented. Call in Bob Kirby to the rescue.



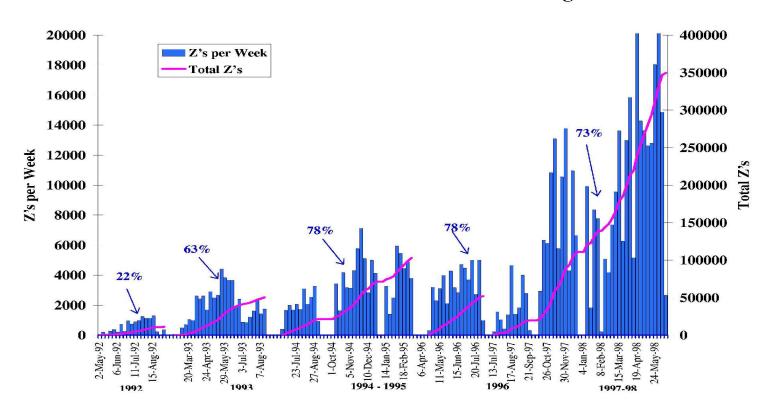


Bob Kirby





1992 - 1998 SLD Polarized Beam Running



Vanda 6/22/98



Polarization at a Future Linear Collider



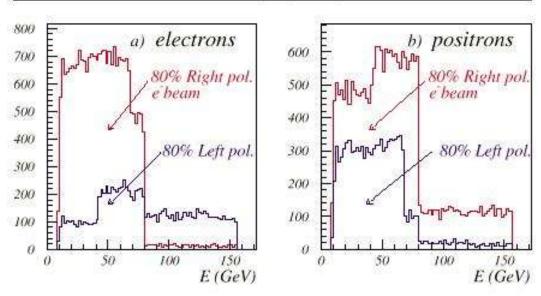


Figure 4.3: Electron and positron energy distributions for selectron pair production, with the indicated beam polarizations and integrated luminosity 50 fb⁻¹ [157].

The case for electron beam polarization is clear. What about the positrons?



The high price of polarized positrons

Photoemission sources give high yields of electrons. The quantum yield is in the range 0.1% to 1.0%, and the "cost" is 1.7 eV per electron. Polarizations as high as 80% are typical.

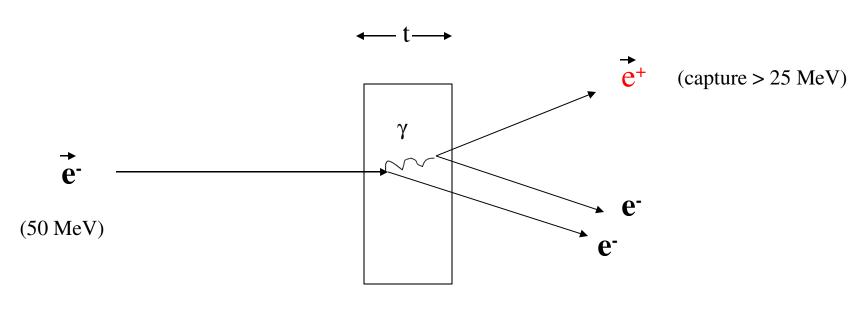
Positrons need 10-100 MeV photons each! Yields of .01% to 0.1 % are typical

- 1) Bremsstrahlung/pair production
- 2) Compton backscattered photons
- 3) Helical Undulators



Bremsstrahlung

A. P. Potylitsin, NIM A 398 395 (1997)



Radiator ~ $0.1 X_0$

Efficiency ~ .001

Polarization ~ 0.5

Beam Power ~ 1.5 MWatts (~ CEBAF)



Compton Source

T. Omori, KEK

Polarized Positron

Source

For Colliders

5-Mar-1997 Delac

We on Positron Source

Tsundiko OMORI (KEK)

by KST Collaboration

KEK

Sumitomo Heavy Industry

Tokyo Matropolitan Univ.

Our Choice 10.6 pm 0.117et 10.6 pm 0.117et 150 Hz

1×10 e/bunch 150 Hz

E beam 80 Met 150 Hz

Low Energy 6.7 GeV target target target (1×1011 e/bunch) t=3.5 mm

× (85 bunch/train)

× (150 Hz)

time structure is the same as the main lines.

SUMMARY

(1) -85 CO2 lasers

10 Jule, 150 Hz

6.7 Get e-linac, 150 Hz

1×10¹¹e/bunch, 85 bunch/train

2. positrons

0.7×10¹⁰et/bunch,

85 bunch/train, 150 Hz

Pol. = 50%

(2) Wall Plug Power = 20 MW

(3) Still Need R/D to get

Conclusion.

(i) design: Capture Section

(ii) Simulation:

CAIN -> Pair Creation on Target



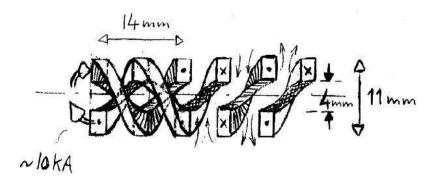
Undulator Source

Balakin and Mikhailichenko, Novosibirsk (1979)

Polarized Positron Sources

helical undulator Surget capture optics

V.E. Balakin, A.A. Mikhai lichenko 1979



	undulator with iron	undulator without iron
undulator period \(\lambda \)	10.0 mm	10.0 mm
inner radius r _i	2.0 mm	2.0 mm
coil width w	2.8 mm	3.3 mm
coil height h	5.5 mm	4.0° mm
yoke height y	5.0 mm	
on-axis field Br	1.3 T	0.62 T
required undulator length	100m	150m



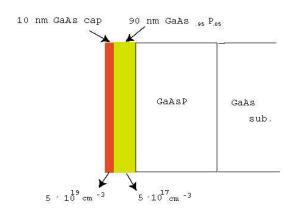
Cathode Charge Limit Solved SLAC 2001

Takashi has worked diligently on the cathode charge limit problem. He has a new design under test, that looks promising for the NLC.

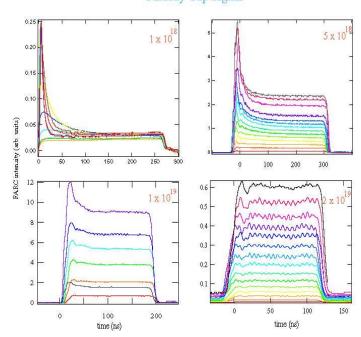


Takashi Maruyama

High Gradient Doped Strained GaAsP

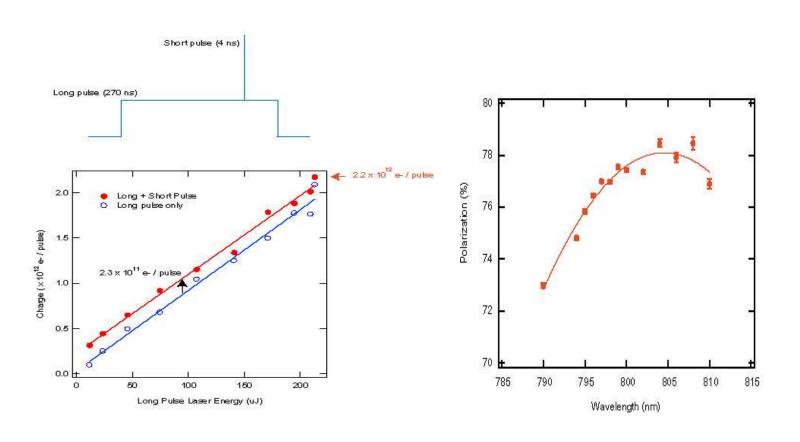


Faraday Cup Signal





The Cathode Charge Limit solved (continued)





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Improved Direct Measurement of Leptonic Coupling Asymmetries with Polarized Z Bosons

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(The SLD Collaboration)

 $\sin^2\Theta_{\rm w} = 0.23098 \pm 0.00026$