

UNIVERSITAT HAMBURG- KYOTO UNIVERSITY Symposium June 7th @ DESY, Hambrug U. Universität Hamburg FORSCHUNG I DER LEHRE I DER BILDUNG



Emergent Spacetime from Quantum Entanglement

Tadashi Takayanagi Yukawa Institute for Theoretical Physics (YITP) Kyoto U. **Γ**



HEORETICAL PHYSIC



Contents

- 1 Introduction
- 2 Quantum Entanglement and Holography
- 3 AdS from Optimization of Path-integrals

(4) Conclusions

1703.00456 + in preparation

Collaborators: Pawel Caputa (YITP) Nilay Kundu (YITP) Masamichi Miyaji (YITP) Kento Watanabe (YITP)



Why Quantum Entanglement ?

Quantum Entanglement = Measure of `Quantumness'

[`Quantumness'

 \Rightarrow We know the total system but not its part.]

= Structures of Quantum Matter

Holography (Gauge/Gravity duality, AdS/CFT)

= Structures of Spacetime

in Gravity (or String theory)

e.g. Surface Area = Entanglement Entropy (EE) Perturbative Einstein eq. = First law of EE

What is the quantum entanglement?

Consider the following states in **two spin systems**:

(i) A direct product state (unentangled state) $|\Psi\rangle = \frac{1}{2} \left[\uparrow \rangle_A + |\downarrow\rangle_A \right] \otimes \left[\uparrow \rangle_B + |\downarrow\rangle_B = |\psi_1\rangle \otimes |\psi_2\rangle.$ Independent

(ii) An entangled state (EPR pair)

$$|\Psi\rangle = \frac{1}{\sqrt{2}} \left[|\uparrow\rangle_{A} \otimes |\downarrow\rangle_{B} - |\downarrow\rangle_{A} \otimes |\uparrow\rangle_{B} \right] \neq |\psi_{1}\rangle \otimes |\psi_{2}\rangle.$$

One determines the other !

We know the state of A+B but not the state of A or B.

Entanglement entropy (EE)

Divide a quantum system into two parts A and B. The total Hilbert space becomes factorized:

$$H_{tot} = H_A \otimes H_B \quad .$$
Example1: Spin Chain
$$A \qquad B$$

Example2: QFT



Define the reduced density matrix P_A for A by

$$\rho_A = \mathrm{Tr}_B \rho_{tot}$$
,

Finally, the entanglement entropy (EE) S_A is defined by

$$S_A = -\mathrm{Tr}_A \ \rho_A \ \log \rho_A$$
 . (von-Neumann entropy)

Quantum Entanglement has recently been applied also to other topics in theoretical physics:

Condensed Matter Theory

Entanglement Entropy (EE), Entanglement Spectrum → Quantum Order Parameter

~ Required `Size' of numerical calculations

• Quantum Field Theories (QFTs)

(Renyi) Entanglement Entropy (EE)

 \rightarrow Universal quantities which characterize

the degrees of freedom of QFTs

 \rightarrow Proof of c-theorem, F-theorem etc.

``Geometrization of QFTs''

An example in Stat-Mech: Quantum Ising spin chain

Consider the Ising spin chain with a transverse magnetic field:



[Vidal-Latorre-Rico-Kitaev 02, Calabrese-Cardy 04]

[1] It is recently reported that (2nd Renyi) EE was measured even experimentally in a cold atom system.

[Greiner et.al. 1509.01160]

Entanglement Entropy for 2 copies of 4-site systems



Markus Greiner (Harvard University) 05

1/1 ページ

Return <first] <pre>orev] [Markus Greiner (Harvard University) 05] [NEXT]> [Inst



<first] <pre>style="text-align: center;">style="text-align: center;"/style="text-align: center;"/>style="text-align: center;"/>styl

[2] Recent Experimental realization of `holographic EE' arXiv:1705.00365

Measuring Holographic Entanglement Entropy on a Quantum Simulator

Keren Li,^{1,2,*} Muxin Han,^{3,4,*} Guilu Long,¹ Yidun Wan,^{5,6,†} Dawei Lu,^{2,7,‡} Bei Zeng,^{2,7,8,9,§} and Raymond Laflamme^{2,9,10}

¹State Key Laboratory of Low-Dimensional Quantum Physics and Department of Physics, Tsinghua University, Beijing 100084, China ²Institute for Quantum Computing and Department of Physics and Astronomy, University of Waterloo, Naterloo N2L 3G1, Ontario, Canada ³Department of Physics, Florida Atantic University, 777 Glades Road, Boca Raton, FL 33431, USA ⁴Institut für Quantengravitation, Universität Erlangen-Nirmberg, Staudtstr. 7B2, 91058 Erlangen, Germany ⁵Department of Physics and Center for Field Theory and Particle Physics, Fudan University, Shanghai 200433, China ⁶Collaborative Innovation Center of Advanced Microstructures, Nanjing, 210093, China ⁷Department of Physics, Southern University of Science and Technology, Shenzhen 518055, China ⁸Department of Mathematics and Statistics, University of Guelph, Guelph NIG 2W1, Ontario, Canada ⁹Perimeter Institute for Advanced Research, Toronto MSG 128, Ontario, Canada ¹⁰Perimeter Institute for Theoretical Physics, Waterloo N2L 2Y5, Ontario, Canada

Anti-de Sitter/conformal field theory (AdS/CFT) correspondence is one of the most promising realizations of holographic principle towards quantum gravity. The recent development of a discrete version of AdS/CFT correspondence in terms of tensor networks motivates one to simulate and demonstrate AdS/CFT correspondence



PT—a build block of a complex delity, which is already state-of some non-negligible decoherence fore, our results successfully tes coherence. Our results further s processor would be able to sim formulae and hence the AdS/CFT of the two assumptions: the dec ments can be tolerated, or the or be *pure*. The latter has been in detail in [28].

FIG. 3: Entanglement entropy $S^{(k)}$ of the k-qubit subsystem of the mult 6 pt. In theory $S^{(k)} = \min\{k, 6, k\}$ as shown by the grange

2 Quantum Entanglement and Holography

(2-1) AdS/CFT



Basic Principle

(Bulk-Boundary relation):

 $|Z_{Gravity} = Z_{CFT}|$

(2-2) Holographic Entanglement Entropy (HEE)

[Ryu-TT 06; derived by Lewkowycz-Maldacena 13]

$$S_{A} = \underset{\substack{\partial \gamma_{A} = \partial A \\ \gamma_{A} \approx A}}{\operatorname{Min}} \left[\frac{\operatorname{Area}(\gamma_{A})}{4G_{N}} \right]$$

 γ_A is the minimal area surface (codim.=2) such that $\partial A = \partial \gamma_A$ and $A \sim \gamma_A$. homologous

Note: In time-dependent spacetimes, we need to take extremal surfaces. [Hubeny-Rangamani-TT 07]



Holographic Proof of Strong Subadditivity [Headrick-TT 07]



Algebraic relations in Quantum Information Theory ⇔ Geometric properties in Gravity

(2-3) Einstein Equation from Entanglement

A is a round ball (radius l) in 3d CFT. Its center is at (t, \vec{x}) . The perturbative Einstein equation is rewritten as follows:

$$R_{\mu\nu} - \frac{1}{2} Rg_{\mu\nu} + \Lambda g_{\mu\nu} = T_{\mu\nu}$$
Kinetic term
$$\int C.C.$$
Matter field
contributions
$$\left(\partial_{l}^{2} - \partial_{l} - \partial_{\vec{x}}^{2} - \frac{3}{l^{2}}\right) \Delta S_{A}(t, \vec{x}, l) = \langle O \rangle \langle O \rangle$$
[Nozaki-Numasawa-Prudenziati-TT 13, Bhattacharya-TT 13]

The perturbative Einstein eq. turns out to be equivalent to the first law of entanglement entropy. [Lashkari-McDermott-Raamsdonk 13, Faulkner-Guica-Hartman-Myers-Raamsdonk 13]

1st law of EE: $S(\rho_A + \Delta \rho_A \parallel \rho_A) = \Delta S_A - \Delta H_A \approx 0$, $(\rho_A \equiv e^{-H_A})$.

(2-4) Emergent Spacetime from Quantum Entanglement

The HEE suggests the following novel interpretation: A spacetime in gravity = Bits of quantum entanglement



⇒ Manifestly realized in the recently argued connection between AdS/CFT and Tensor networks !

[Swingle 09, cf. Raamsdonk 09, Susskind-Maldacena 13,]

(2-5) Tensor Networks and AdS/CFT

Tensor Network[See e.g. Cirac-Verstraete 09(review)]

Tensor network = Geometrical description of wave function

- ⇒ Efficient variational ansatz for the ground state wave functions in quantum many-body systems.
- ⇒ An ansatz should respect the correct <u>quantum entanglement</u> of ground state.

~Geometry of Tensor Network

<u>MERA</u> (Multiscale Entanglement Renormalization Ansatz): ⇒ An efficient variational ansatz for CFT vacuum. [Vidal 05] To increase entanglement in a CFT, we add (dis)entanglers.



The original idea: Tensor Network of MERA (\exists scale inv.) = a time slice of AdS space

However, recent detailed studies show several problems:(1) MERA does not have any isometry other than scale inv.(2) Why EE bound is saturated ?(3) Sub AdS scale locality

Some of these problems may be due to lattice artifacts. We should consider CFTs in genuine continuum limit.

In our latest work, we gave a new approach which realizes continuous description **using path-integrals**.

③ Optimization of Path-Integral and AdS/CFT(3-1) Motivation

Remember that the MERA can be obtained from the `optimization' of tensor networks

⇒ Tensor network renormalization [Evenbly-Vidal 14, 15]



Optimization of Path-Integral [Miyaji-Watanabe-TT 2016]



(3-2) Formulating Optimization of Path-integral

[Caputa-Kundu-Miyaji-Watanabe-TT 2017]

<u>A Basic Rule</u>: Simplify a path-integral such that it produces the correct UV wave functional $\Psi_{UV}^{\text{Flat}}[\Phi(x)]$.

Below we focus on 2d CFTs for simplicity.

Deformation of discretizations in path-integral = Curved metric such that one cell (bit) = unit length.

$$ds^{2} = e^{2\phi(x,z)}(dx^{2} + dz^{2}).$$

Note: The original flat metric is given as follows (ε is the UV cutoff): $ds^2 = \varepsilon^{-2} \cdot (dx^2 + dz^2).$

Ground state UV wave function in curved space

$$\Psi_{UV}^{g}\left[\Phi(x)\right] = \int \prod_{\substack{0 < z < \infty \\ -\infty < x < \infty}} D\Phi(x, z) \, e^{-S_{CFT}(\Phi)} \cdot \delta\left(\Phi(x) - \Phi(x, z = 0)\right)$$

In CFTs, owing to the Weyl invariance, we have

$$\Psi_{UV}^{g_{ab}=e^{2\phi}\delta_{ab}}\left[\Phi(x)\right] = N[\phi(x,z)] \cdot \Psi_{UV}^{\text{Flat}}\left[\Phi(x)\right].$$

<u>**Our Proposal**</u> (Optimization of Path-integral for CFTs): Minimize $N[\phi(x,z)]$ w.r.t $\phi(x,z)$ with the boundary condition $e^{2\phi}|_{z=\varepsilon} = \varepsilon^{-2}$.

(3-3) AdS3/CFT2 (Vaccum State)

[Caputa-Kundu-Miyaji-Watanabe-TT 2017]

Metric of Discretized Lattice : $ds^2 = e^{2\phi}(dx^2 + dz^2)$

$$\Rightarrow \frac{\Psi_{g_{ab}=e^{2\phi}\cdot\delta_{ab}}}{\Psi_{g_{ab}=\delta_{ab}}} \propto e^{\frac{c}{24\pi}S_L},$$

$$S_L = \int dx dz \Big[(\partial_x \phi)^2 + (\partial_z \phi)^2 + e^{2\phi} \Big] \quad \text{[Liouville Theory]}$$

$$= \int dx dz \Big[(\partial_x \phi)^2 + (\partial_z \phi + e^{\phi})^2 \Big]$$

$$\Rightarrow \text{Minimum}: \quad e^{2\phi} = \frac{1}{z^2}. \quad \Rightarrow 2 \text{dim. Hyperbolic space}$$

$$= a \text{ time slice of AdS3}$$

(3-4) Analysis of Excited States

Consider the vacuum state of 2d CFT on a circle dual to a global AdS3. ⇒ Optimize the path-integral on a disk with the unit radius.

The solution is simply given by (in the general expressions)

$$A(w) = w, \quad B(\overline{w}) = \overline{w}.$$

$$ds^{2} = \frac{4dwd\overline{w}}{(1-|w|^{2})^{2}}.$$
 Hyperbolic Disk H2

Now we insert an operator O(x) in the center of the disk x=x0. O(x): conformal dim. hL=hR=h



 $\Rightarrow O(x) \sim e^{-2h \cdot \varphi}.$

Thus we minimize

$$\frac{\Psi_{g=e^{2\varphi}}}{\Psi_{\text{Flat}}} \propto e^{\frac{c}{24\pi}S_L} \cdot e^{-2h\varphi(x_0)}.$$

$$\partial_{w}\partial_{\overline{w}}\varphi - \frac{1}{4}e^{2\varphi} + \frac{6\pi h}{c}\delta^{2}(w) = 0.$$

Solution:
$$A(w) = w^{a}$$
, $B(\overline{w}) = \overline{w}^{a}$, $(a \equiv 1 - 12h/c)$.
Metric: $ds^{2} = \frac{4d\zeta d\overline{\zeta}}{(1 - |\zeta|^{2})^{2}}$, $\zeta \equiv w^{a} = re^{i\theta}$

 \Rightarrow Deficit angle geometry $\theta \sim \theta + 2\pi a$.

This agrees with the expected gravity dual if h/c<<1.

Note: the AdS/CFT predicts $a = \sqrt{1 - 24h/c}$. Interestingly, if we consider the quantum Liouville CFT, then $h = \frac{\gamma \alpha}{4}(Q - \alpha \gamma/2.), \quad c = 1 + 3Q^2, \quad (Q \equiv 2/\gamma + \gamma).$ \Rightarrow We get $a = \sqrt{1 - 24h/c}$. If we insert a local operator in the path-integral:



This agrees with general relativity !



④ Conclusions

- Holography tells us that gravitational dynamics is dual to that of quantum entanglement in QFTs.
 - ⇒ Emergent spacetime from quantum entanglement !
- We argue that the optimization of Euclidean path-integral is an interesting candidate of emergent spacetime in AdS/CFT. This works even in continuum limit.
 - ⇒Many future problems: Higher dim. version ?
 - Time-dependent b.g. ?
 - dS/CFT version ?

Motivation of Our Proposal

The normalization N estimates repetitions of same operations of path-integration. \rightarrow Minimize this !

⇒ Our conjecture: $N[\varphi(x,z)] \approx \exp[C[\varphi]]$ $C[\phi] \equiv \text{computational complexity of TN state}$ [Relevance of Complexity in ~ # of Tensors in TN holography: Susskind et.al. See also Czech 1706.00965] **Our Hope Optimization** of Tensor Network (~ Path-integral) **Minimizing** Computational Complexity **Solving** Einstein equation

Thank you !

Your visit to Kyoto U is very welcome !







Strings 2018 in OIST, Okinawa: June 25-29



Post Strings WS in YITP, Kyoto: July 2- Aug 3