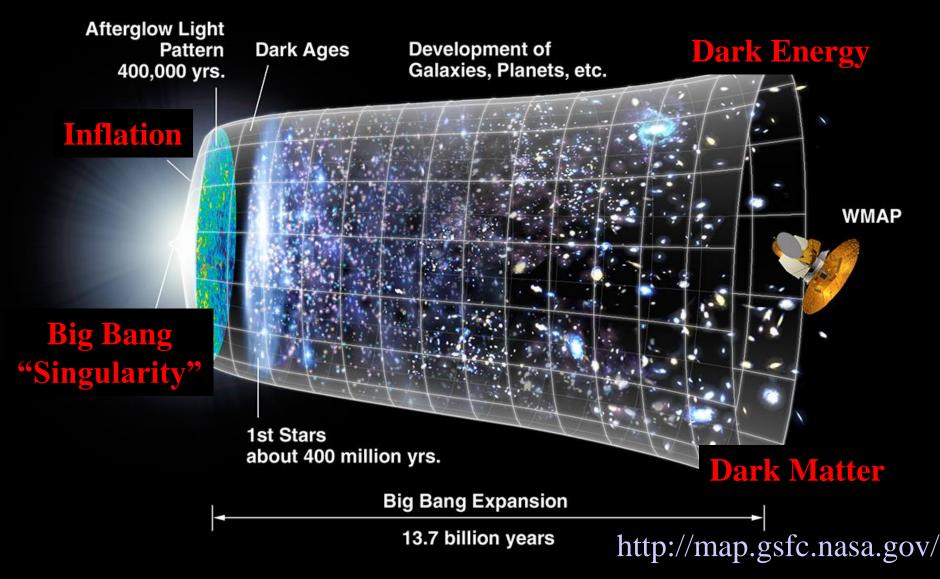
Massive gravity and cosmology

Shinji Mukohyama Yukawa Institute for Theoretical Physics Kyoto University

Based on collaborations with

Katsuki Aoki, Antonio DeFelice, Garrett Goon, Emir Gumrukcuoglu, Lavinia Heisenberg, Kurt Hinterbichler, Sachiko Kuroyanagi, David Langlois, Chunshan Lin, Charles Mazuet, Ryo Namba, Atsushi Naruko, Michele Oliosi, Takahiro Tanaka, Norihiro Tanahashi, Mark Trodden, Jean-Philippe Uzan, Mikhail Volkov

Why alternative gravity theories?



Three conditions for good alternative theories of gravity (my personal viewpoint)

- 1. Theoretically consistent e.g. no ghost instability
- 2. Experimentally viable solar system / table top experiments
- 3. Predictable e.g. protected by symmetry

Some examples

- I. Ghost condensation
 IR modification of gravity
 motivation: dark energy/matter
- II. Nonlinear massive gravityIR modification of gravitymotivation: "Can graviton have mass?"
- III. Horava-Lifshitz gravityUV modification of gravitymotivation: quantum gravity
- IV. Superstring theoryUV modification of gravitymotivation: quantum gravity, unified theory

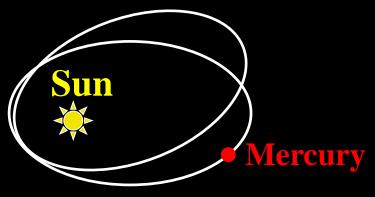
A motivation for IR modification

- Gravity at long distances
 Flattening galaxy rotation curves
 extra gravity

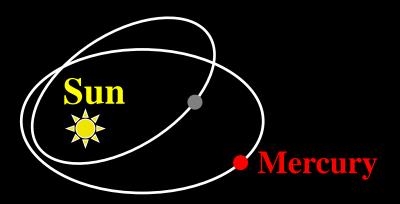
 Dimming supernovae
 accelerating universe
- Usual explanation: new forms of matter (DARK MATTER) and energy (DARK ENERGY).

Dark component in the solar system?

Precession of perihelion observed in 1800's...



which people tried to explain with a "dark planet", Vulcan,



But the right answer wasn't "dark planet", it was "change gravity" from Newton to GR. Can we change gravity in IR?

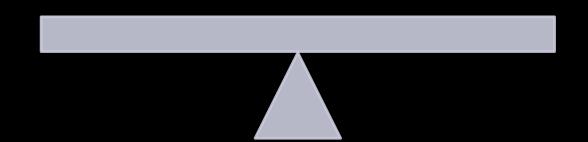
Change Theory? Massive gravity Fierz-Pauli 1939 DGP model Dvali-Gabadadze-Porrati 2000

Change State? Higgs phase of gravity The simplest: Ghost condensation

Arkani-Hamed, Cheng, Luty and Mukohyama, JHEP 0405:074,2004.

Simple question: Can graviton have mass? May lead to acceleration without dark energy





Simple question: Can graviton have mass? May lead to acceleration without dark energy



Fierz-Pauli theory (1939) Unique linear theory without instabilities (ghosts)

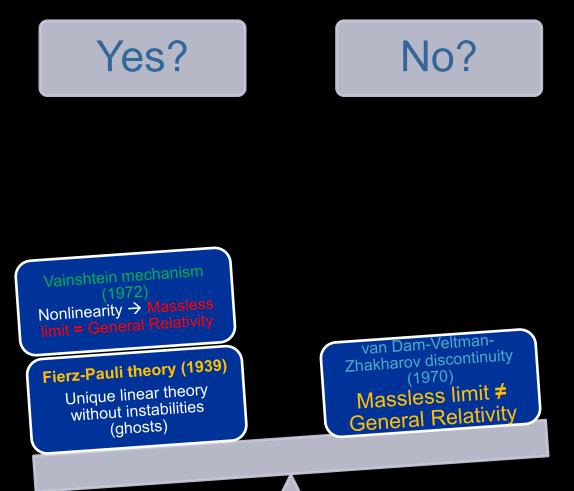
Simple question: Can graviton have mass? May lead to acceleration without dark energy



Fierz-Pauli theory (1939)

Unique linear theory without instabilities (ghosts) van Dam-Veltman-Zhakharov discontinuity (1970) Massless limit ≠ General Relativity

Simple question: Can graviton have mass? May lead to acceleration without dark energy



Simple question: Can graviton have mass? May lead to acceleration without dark energy



Vainshtein mechanism (1972) Nonlinearity → Massless limit = General Relativity

Fierz-Pauli theory (1939) Unique linear theory

without instabilities (ghosts) Boulware-Deser ghost (1972) 6th d.o.f.@Nonlinear level → Instability (ghost)

van Dam-Veltman-Zhakharov discontinuity (1970) Massless limit ≠ General Relativity

Nonlinear massive gravity

de Rham, Gabadadze 2010 de Rham, Gabadadze & Tolley 2010

- First example of fully nonlinear massive gravity without BD ghost since 1972!
- Purely classical (but technically natural)
- Properties of 5 d.o.f. depend on background
- 4 scalar fields φ^a (a=0,1,2,3)
- Poincare symmetry in the field space: $\phi^a \rightarrow \phi^a + c^a, \quad \phi^a \rightarrow \Lambda^a_b \phi^b$



$$f_{\mu\nu} \equiv \eta_{ab} \partial_{\mu} \phi^a \partial_{\nu} \phi^b$$

fiducial metric

Pullback of Minkowski metric in field space to spacetime

Systematic resummation

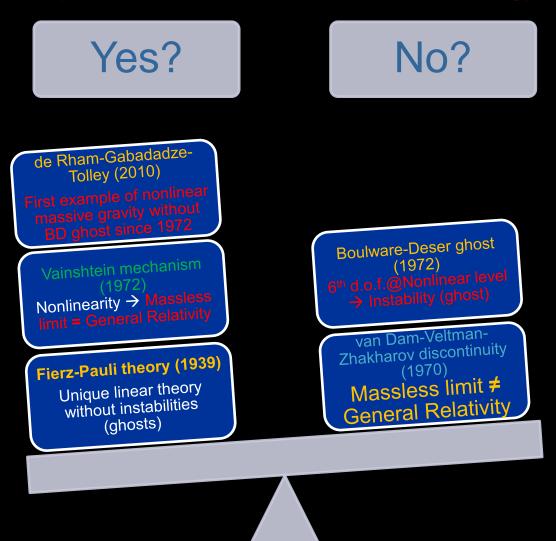
de Rham, Gabadadze & Tolley 2010

$$\begin{split} I_{mass}[g_{\mu\nu}, f_{\mu\nu}] &= M_P^2 m_g^2 \int d^4 x \sqrt{-g} \left(\mathcal{L}_2 + \alpha_3 \mathcal{L}_3 + \alpha_4 \mathcal{L}_4 \right) \\ f_{\mu\nu} &\equiv \eta_{ab} \partial_\mu \phi^a \partial_\nu \phi^b \\ \mathcal{L}_2 &= \frac{1}{2} \left([\mathcal{K}]^2 - [\mathcal{K}^2] \right) \\ \mathcal{L}_3 &= \frac{1}{6} \left([\mathcal{K}]^3 - 3 \left[\mathcal{K} \right] \left[\mathcal{K}^2 \right] + 2 \left[\mathcal{K}^3 \right] \right) \\ \mathcal{L}_4 &= \frac{1}{24} \left([\mathcal{K}]^4 - 6 \left[\mathcal{K} \right]^2 \left[\mathcal{K}^2 \right] + 3 \left[\mathcal{K}^2 \right]^2 + 8 \left[\mathcal{K} \right] \left[\mathcal{K}^3 \right] - 6 \left[\mathcal{K}^4 \right] \end{split}$$

No helicity-0 ghost, i.e. no BD ghost, in decoupling limit $\mathcal{K}_{\mu\nu} = \partial_{\mu}\partial_{\nu}\pi \implies \mathcal{L}_{2,3,4} = (\text{total derivative})$

No BD ghost away from decoupling limit (Hassan&Rosen)

Simple question: Can graviton have mass? May lead to acceleration without dark energy



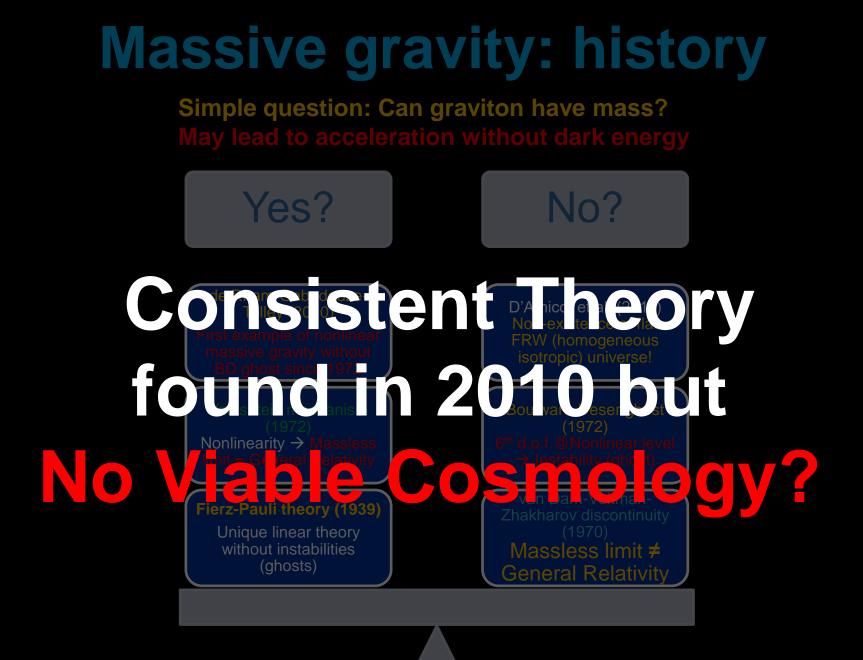
No FLRW universe?

D'Amico, de Rham, Dubovsky, Gabadadze, Pirtshalava, Tolley (2011)

- Flat FLRW ansatz in "Unitary gauge"
 - $g_{\mu\nu}dx^{\mu}dx^{\nu} = -N^{2}(t)dt^{2} + a^{2}(t)(dx^{2}+dy^{2}+dz^{2})$ $\phi^{a} = x^{a} \qquad \longrightarrow \qquad f_{\mu\nu} = \eta_{\mu\nu}$
- Bianchi "identity" \rightarrow a(t) = const. c.f. $\nabla^{\mu} \left(\frac{2}{\sqrt{-g}} \frac{\delta I}{\delta g^{\mu\nu}} \right) = \frac{1}{\sqrt{-g}} \frac{\delta I_g}{\delta \phi^a} \partial_{\nu} \phi^a$ \rightarrow poportrivial flat FLRW/ cospole

ightarrow no non-trivial flat FLRW cosmology

- "Our conclusions on the absence of the homogeneous and isotropic solutions do not change if we allow for a more general maximally symmetric 3-space"
- c.f. Inhomogeneous solutions do exist. [Koyama, Niz, Tasinato 2011; Chamseddine & Volkov 2011]



Cosmological solutions in nonlinear massive gravity

Good?



D'Amico, et.al. (2011) Non-existence of flat FLRW (homogeneous isotropic) universe!

Open FLRW solutions

Gumrukcuoglu, Lin, Mukohyama, arXiv: 1109.3845 [hep-th]

- $f_{\mu\nu}$ spontaneously breaks diffeo.
- Both $g_{\mu\nu}$ and $f_{\mu\nu}$ must respect FLRW symmetry
- Need FLRW coordinates of Minkowski $f_{\mu\nu}$
- No closed FLRW chart
- Open FLRW ansatz

$$\begin{split} \phi^{0} &= f(t)\sqrt{1+|K|(x^{2}+y^{2}+z^{2})}, \\ \phi^{1} &= \sqrt{|K|}f(t)x, \\ \phi^{2} &= \sqrt{|K|}f(t)y, \\ \phi^{3} &= \sqrt{|K|}f(t)z. \end{split}$$

 $f_{\mu\nu}dx^{\mu}dx^{\nu} = -(\dot{f}(t))^2 dt^2 + |K| (f(t))^2 \Omega_{ij}(x^k) dx^i dx^j$

$$g_{\mu\nu}dx^{\mu}dx^{\nu} = -N(t)^{2}dt^{2} + a(t)^{2}\Omega_{ij}dx^{i}dx^{j},$$

$$\Omega_{ij}dx^{i}dx^{j} = dx^{2} + dy^{2} + dz^{2} - \frac{|K|(xdx + ydy + zdz)^{2}}{1 + |K|(x^{2} + y^{2} + z^{2})},$$

Open FLRW solutions

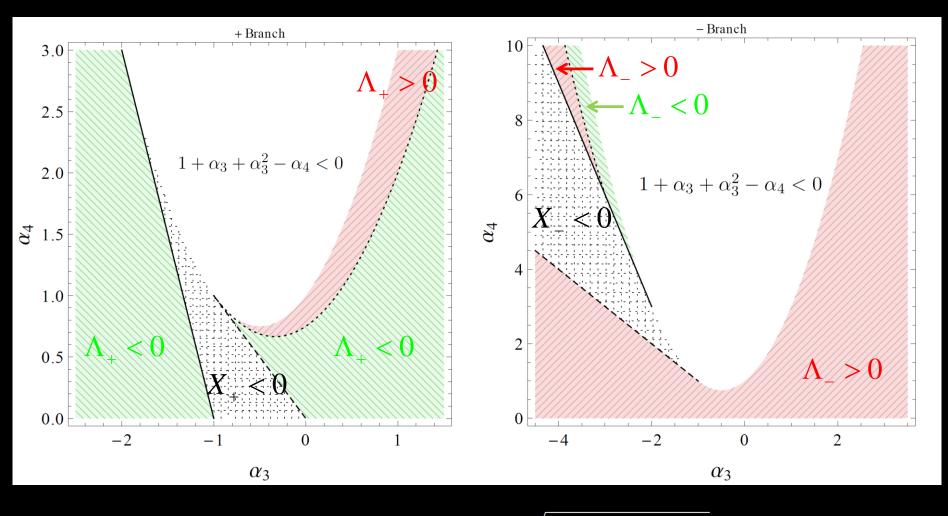
Gumrukcuoglu, Lin, Mukohyama, arXiv: 1109.3845 [hep-th]

• EOM for ϕ^{a} (a=0,1,2,3) $(\dot{a} - \sqrt{|K|}N) \left[\left(3 - \frac{2\sqrt{|K|}f}{a} \right) + \alpha_{3} \left(3 - \frac{\sqrt{|K|}f}{a} \right) \left(1 - \frac{\sqrt{|K|}f}{a} \right) + \alpha_{4} \left(1 - \frac{\sqrt{|K|}f}{a} \right)^{2} \right] = 0$ • The first sol $\dot{a} = \sqrt{|K|}N$ implies $g_{\mu\nu}$ is Minkowski \rightarrow we consider other solutions

$$f = \frac{a}{\sqrt{|K|}} X_{\pm}, \quad X_{\pm} \equiv \frac{1 + 2\alpha_3 + \alpha_4 \pm \sqrt{1 + \alpha_3 + \alpha_3^2 - \alpha_4}}{\alpha_3 + \alpha_4}$$

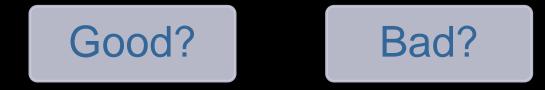
- Latter solutions do not exist if K=0
- Metric EOM \rightarrow self-acceleration $3 H^2 + \frac{3 K}{a^2} = \Lambda_{\pm} + \frac{1}{M_{Pl}^2} \rho$ $\Lambda_{\pm} \equiv -\frac{m_g^2}{(\alpha_3 + \alpha_4)^2} \left[(1 + \alpha_3) \left(2 + \alpha_3 + 2 \alpha_3^2 - 3 \alpha_4 \right) \pm 2 \left(1 + \alpha_3 + \alpha_3^2 - \alpha_4 \right)^{3/2} \right]$

Self-acceleration



$$f = \frac{a}{\sqrt{|K|}} X_{\pm}, \quad X_{\pm} \equiv \frac{1 + 2\alpha_3 + \alpha_4 \pm \sqrt{1 + \alpha_3 + \alpha_3^2 - \alpha_4}}{\alpha_3 + \alpha_4}$$

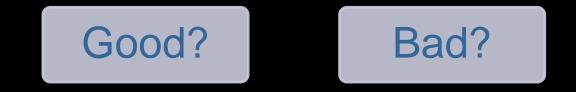
Cosmological solutions in nonlinear massive gravity



Open universes with selfacceleration GLM (2011a) D'Amico, et.al. (2011) Non-existence of flat FLRW (homogeneous isotropic) universe!

GLM = Gumrukcuoglu-Lin-Mukohyama

Cosmological solutions in nonlinear massive gravity



More general fiducial metric f_{μυ} closed/flat/open FLRW universes allowed GLM (2011b)

Open universes with selfacceleration GLM (2011a) D'Amico, et.al. (2011) Non-existence of flat FLRW (homogeneous isotropic) universe!

GLM = Gumrukcuoglu-Lin-Mukohyama

Summary of Introduction + α

- Nonlinear massive gravity free from BD ghost
- FLRW background No closed/flat universe
 Open universes with self-acceleration!
- More general fiducial metric $f_{\mu\nu}$ closed/flat/open FLRW universes allowed Friedmann eq does not depend on $f_{\mu\nu}$
- Cosmological linear perturbations Scalar/vector sectors → same as in GR Tensor sector → time-dependent mass

Nonlinear instability

DeFelice, Gumrukcuoglu, Mukohyama, arXiv: 1206.2080 [hep-th]

- de Sitter or FLRW fiducial metric
- Pure gravity + bare $cc \rightarrow$ FLRW sol = de Sitter
- Bianchi I universe with axisymmetry + linear perturbation (without decoupling limit)
- Small anisotropy expansion of Bianchi I + linear perturbation
 - \rightarrow nonlinear perturbation around flat FLRW

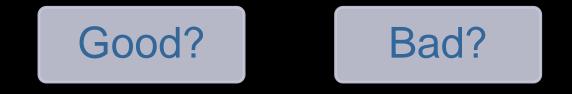
Odd-sector:

1 healthy mode + 1 healthy or ghosty mode

Even-sector: 2 healthy modes + 1 ghosty mode

• This is not BD ghost nor Higuchi ghost.

Cosmological solutions in nonlinear massive gravity



More general fiducial metric f_{μυ} closed/flat/open FLRW universes allowed GLM (2011b)

Open universes with self acceleration GLM (2011a) NEW Nonlinear instability of FLRW solutions DGM (2012)

D'Amico, et.al. (2011) Non-existence of flat FLRW (homogeneous isotropic) universe!

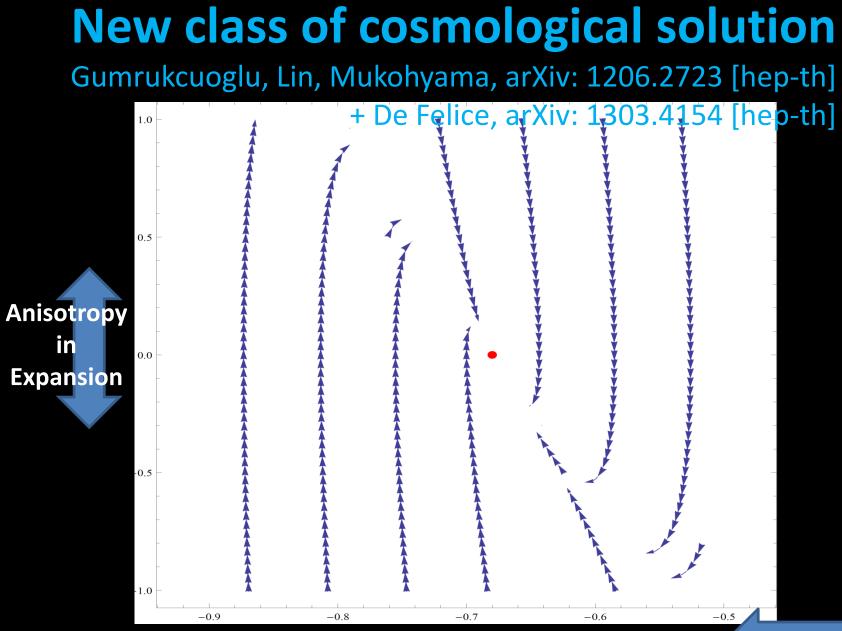
GLM = Gumrukcuoglu-Lin-Mukohyama DGM = DeFelice-Gumrukcuoglu-Mukohyama

New class of cosmological solution

Gumrukcuoglu, Lin, Mukohyama, arXiv: 1206.2723 [hep-th] + De Felice, arXiv: 1303.4154 [hep-th]

- Healthy regions with (relatively) large anisotropy
- Are there attractors in healthy region?
- Classification of fixed points
- Local stability analysis
- Global stability analysis

At attractors, physical metric is isotropic but fiducial metric is anisotropic.
 → Anisotropic FLRW universe! statistical anisotropy expected (suppressed by small m_g²)



Anisotropy in fiducial metric

Cosmological solutions in nonlinear massive gravity



GLM = Gumrukcuoglu-Lin-Mukohyama DGM = DeFelice-Gumrukcuoglu-Mukohyama

New backgrounds or Extended theories

- New nonlinear instability [DeFelice, Gumrukcuoglu, Mukohyama 2012]
 → (i) new backgrounds, or (ii) extended theories
- (i) Anisotropic FLRW (Gumrukcuoglu, Lin, Mukohyama 2012): physical metric is isotropic but fiducial metric is anisotropic
- (ii) Extended quasidilaton (De Felice&Mukohyama 2013), Bimetric theory (Hassan, Rosen 2011; DeFelice, Nakamura, Tanaka 2013; DeFelice, Gumrukcuoglu, Mukohyama, Tanahashi, Tanaka 2014), Rotation-invariant theory (Rubakov 2004; Dubovsky 2004; Blas, Comelli, Pilo 2009; Comelli, Nesti, Pilo 2012; Langlois, Mukohyama, Namba, Naruko 2014), Composite metric (de Rham, Heisenberg, Ribeiro 2014; Gumrukcuoglu, Heisenberg, Mukohyama 2014, 2015), New quasidilaton (Mukohyama 2014; De Felice, Gumrukcuoglu, Heisenberg, Mukohyama, Tanahashi 2016), Chameleonic bigravity (De Felice, Mukohyama, Uzan 2017), ...
- They provide stable cosmology.

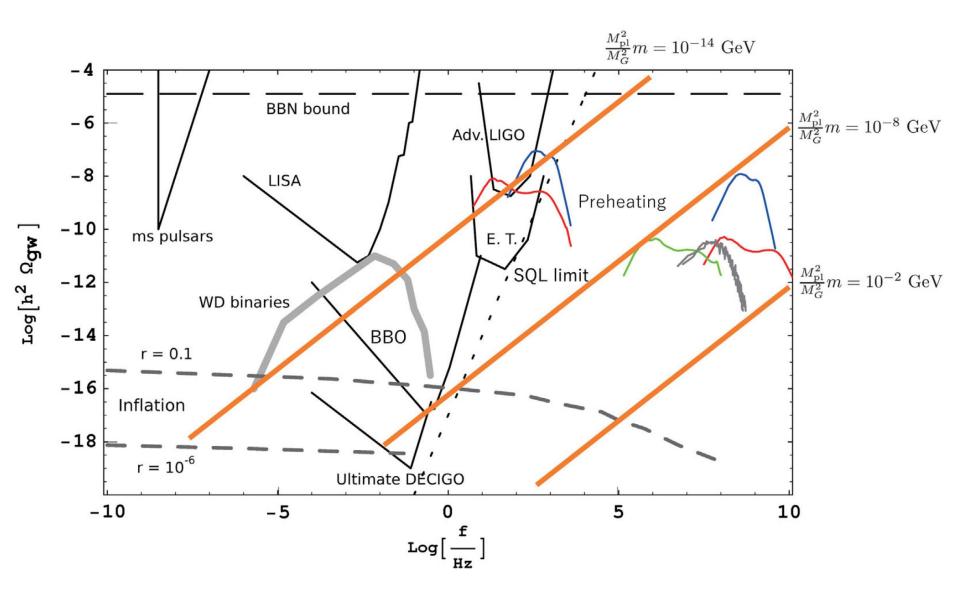
PHYSICAL REVIEW D 94, 024001 (2016)

Massive gravitons as dark matter and gravitational waves

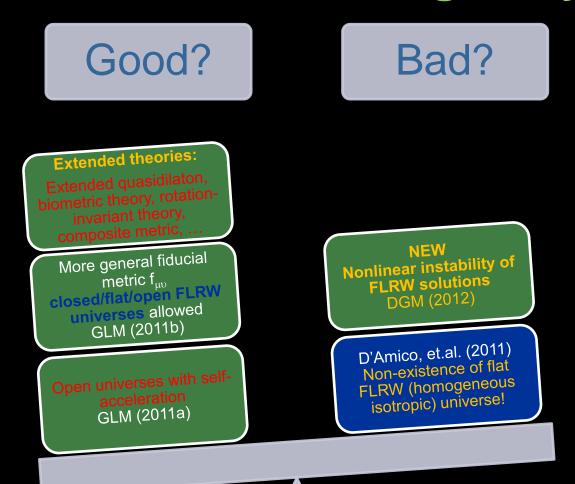
Katsuki Aoki^{1,*} and Shinji Mukohyama^{2,3,†}

¹Department of Physics, Waseda University, Shinjuku, Tokyo 169-8555, Japan ²Center for Gravitational Physics, Yukawa Institute for Theoretical Physics, Kyoto University, 606-8502 Kyoto, Japan ³Kavli Institute for the Physics and Mathematics of the Universe (WPI), UTIAS, The University of Tokyo, Kashiwa, Chiba 277-8583, Japan (Received 2 May 2016; published 1 July 2016)

We consider the possibility that the massive graviton is a viable candidate for dark matter in the context of bimetric gravity. We first derive the energy-momentum tensor of the massive graviton and show that it indeed behaves as that of dark matter fluid. We then discuss a production mechanism and the present abundance of massive gravitons as dark matter. Since the metric to which ordinary matter fields couple is a linear combination of the two mass eigenstates of bigravity, production of massive gravitons, i.e., the dark matter particles, is inevitably accompanied by generation of massless gravitons, i.e., the gravitational waves. Therefore, in this scenario some information about dark matter in our Universe is encoded in gravitational waves. For instance, if LIGO detects gravitational waves generated by the preheating after inflation, then the massive graviton with the mass of ~0.01 GeV is a candidate for dark matter.



Cosmological solutions in nonlinear massive gravity



GLM = Gumrukcuoglu-Lin-Mukohyama DGM = DeFelice-Gumrukcuoglu-Mukohyama

More recent development Minimal Theory of Massive Gravity

De Felice & Mukohyama, arXiv: 1506.01594 1512.04008

- 2 physical dof only = massive gravitational waves
- exactly same FLRW background as in dRGT
- no BD ghost, no Higuchi ghost, no nonlinear ghost

Three steps to the Minimal Theory

- 1. Fix local Lorentz to realize ADM vielbein in dRGT
- 2. Switch to Hamiltonian
- 3. Add 2 additional constraints

(It is easy to go back to Lagrangian after 3.)

Gravitational wave signal from massive gravity

A Emir Gümrükçüoğlu¹, Sachiko Kuroyanagi², Chunshan Lin¹, Shinji Mukohyama¹ and Norihiro Tanahashi³

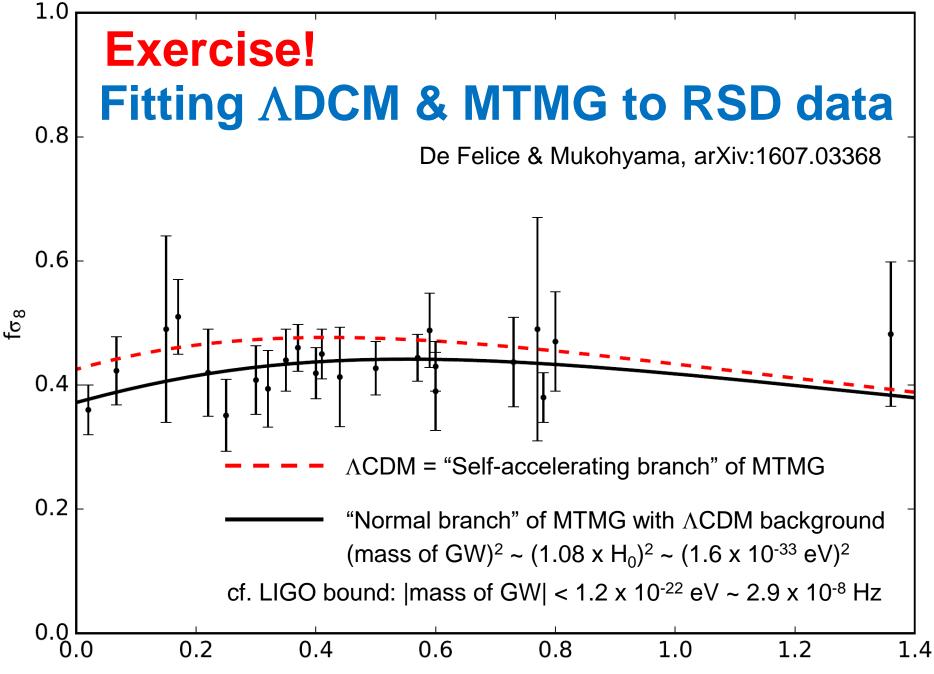
¹ Kavli Institute for the Physics and Mathematics of the Universe, Todai Institutes for Advanced Study, University of Tokyo, 5-1-5 Kashiwanoha, Kashiwa, Chiba 277-8583, Japan
 ² Research Center for the Early Universe (RESCEU), Graduate School of Science, The University of Tokyo, Tokyo 113-0033, Japan
 ³ Department of Physics, University of California, Davis, CA 95616, USA

E-mail: emir.gumrukcuoglu@ipmu.jp, skuro@resceu.s.u-tokyo.ac.jp, chunshan.lin@ipmu.jp, shinji.mukohyama@ipmu.jp and tanahashi@ms.physics.ucdavis.edu

Received 14 September 2012, in final form 14 October 2012 Published 14 November 2012 Online at stacks.iop.org/CQG/29/235026

Abstract

We discuss the detectability of gravitational waves with a time-dependent mass contribution, by means of the stochastic gravitational wave observations. Such a mass term typically arises in the cosmological solutions of massive gravity theories. We conduct the analysis based on a general quadratic action, and thus the results apply universally to any massive gravity theories in which the modification of general relativity appears primarily in the tensor modes. The primary manifestation of the modification in the gravitational wave spectrum is a sharp peak. The position and height of the peak carry information on the present value of the mass term, as well as the duration of the inflationary stage. We also discuss the detectability of such a gravitational wave signal using the future-planned gravitational wave observatories.



Cosmological solutions in nonlinear massive gravity



GLM = Gumrukcuoglu-Lin-Mukohyama DGM = DeFelice-Gumrukcuoglu-Mukohyama

DGHM = DeFelice-Gumrukcuoglu-Heisenberg-Mukohyama

Summary

- Nonlinear massive gravity free from BD ghost
- FLRW background No closed/flat universe
 Open universes with self-acceleration!
- More general fiducial metric $f_{\mu\nu}$ closed/flat/open FLRW universes allowed Friedmann eq does not depend on $f_{\mu\nu}$
- Cosmological linear perturbations
 Scalar/vector sectors → same as in GR
 Tensor sector → time-dependent mass
- All homogeneous and isotropic FLRW solutions in the original dRGT theory have infinitely strong coupling and ghost instability
- Stable cosmology realized in (i) new class of cosmological solution or (ii) extended theories
- Minimal theory of massive gravity with 2dof provides a nonlinear completion of dRGT self-accelerating cosmology

Why alternative gravity theories?

