

*Soft interaction processes at HERA  
(leading baryons, multiple interactions)*



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representing the H1 and ZEUS Collaborations

Outline:

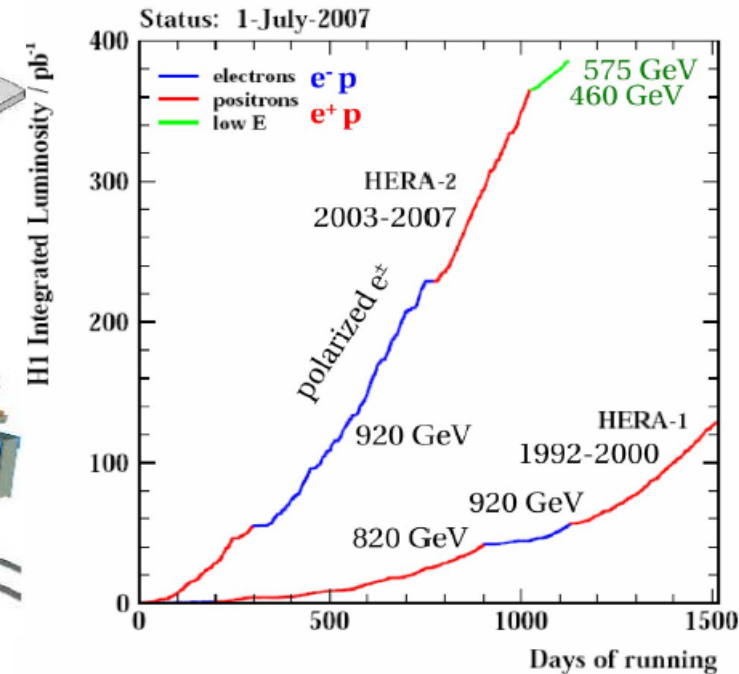
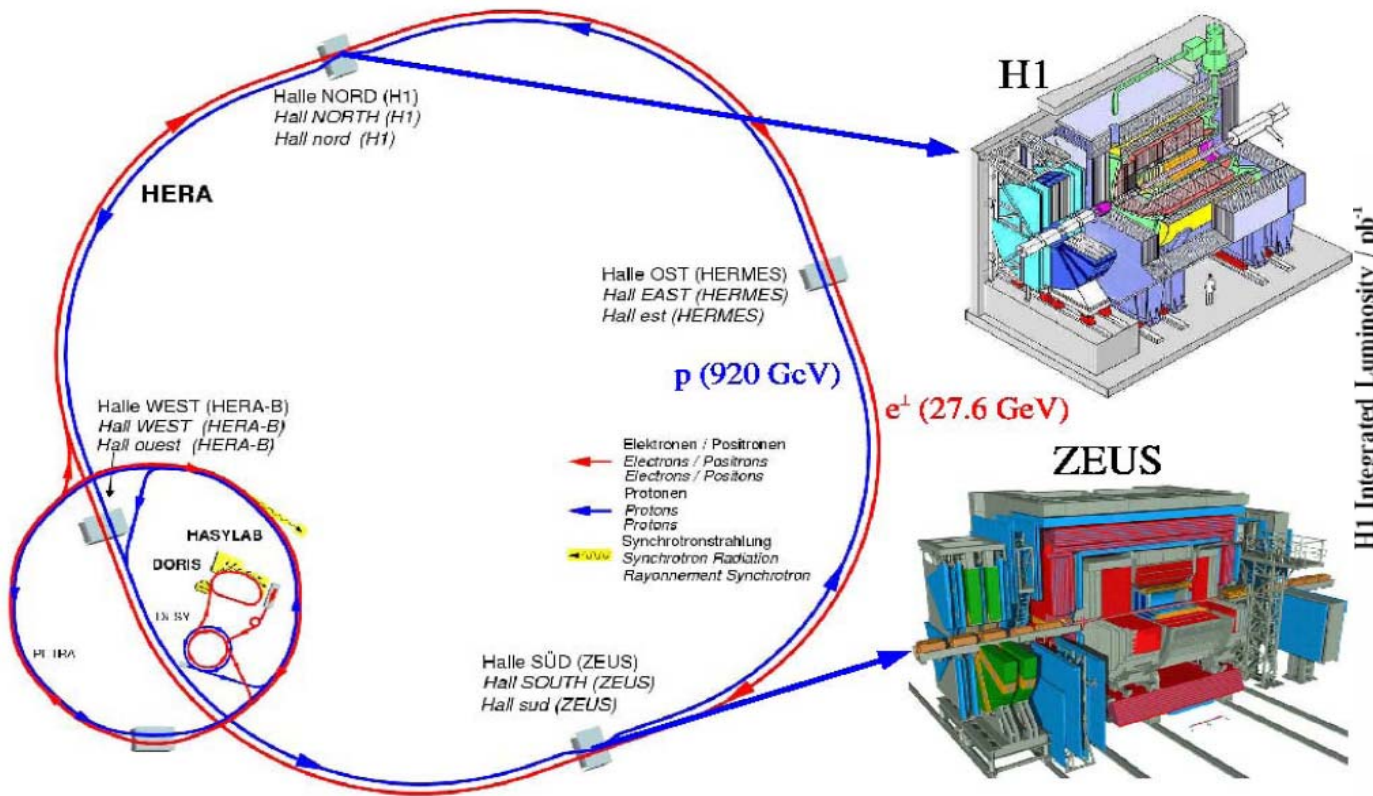
- Leading baryon production
  - comparison with models
  - estimation of pion structure function
  - absorption/rescattering effects
- Studies of multi-parton interactions in photoproduction

# HERA

The world's only electron/positron-proton collider at DESY, Hamburg

$E_e = 27.6 \text{ GeV}$   $E_p = 920 \text{ GeV}$  (also 820, 460 and 575 GeV)

(total centre-of-mass energy of collision up to  $\sqrt{s} \approx 320 \text{ GeV}$ )

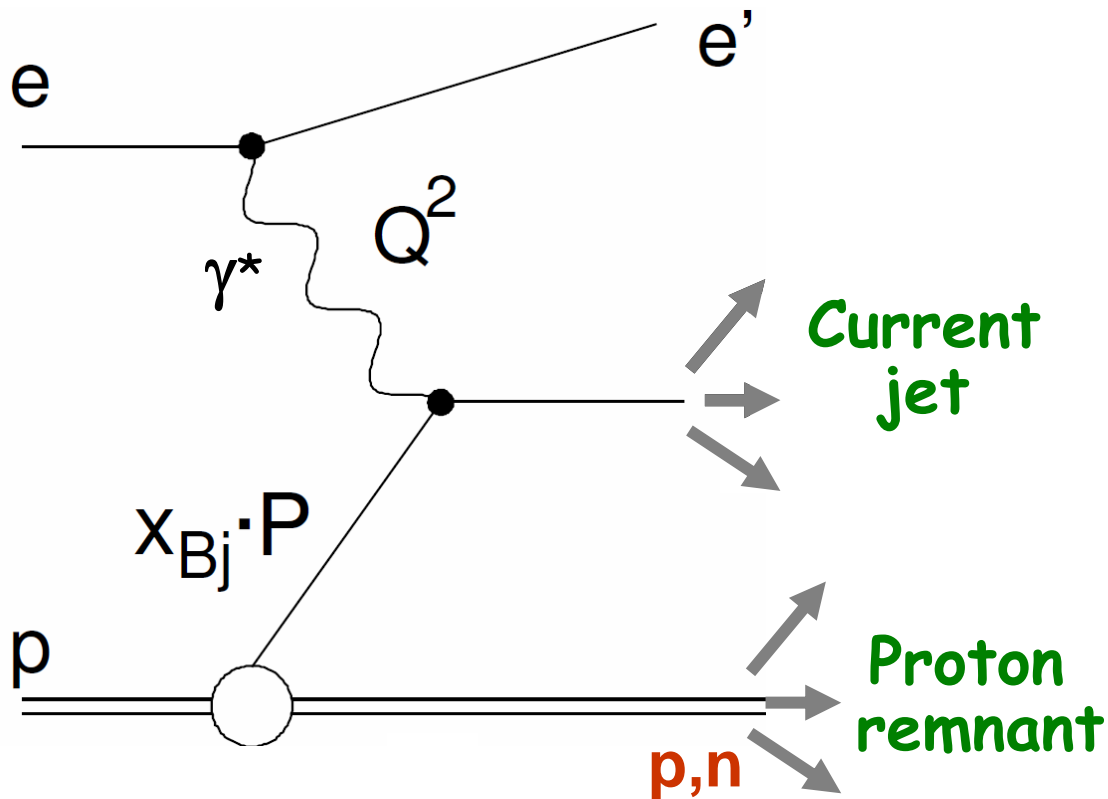


Two colliding experiments: H1 and ZEUS

HERA-1: 1992 - 2000

HERA-2: 2003 - 2007

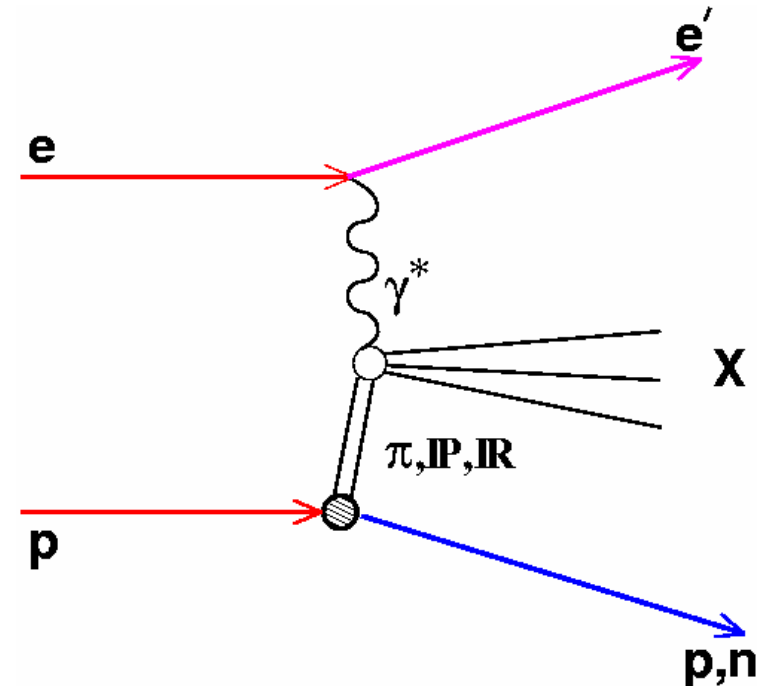
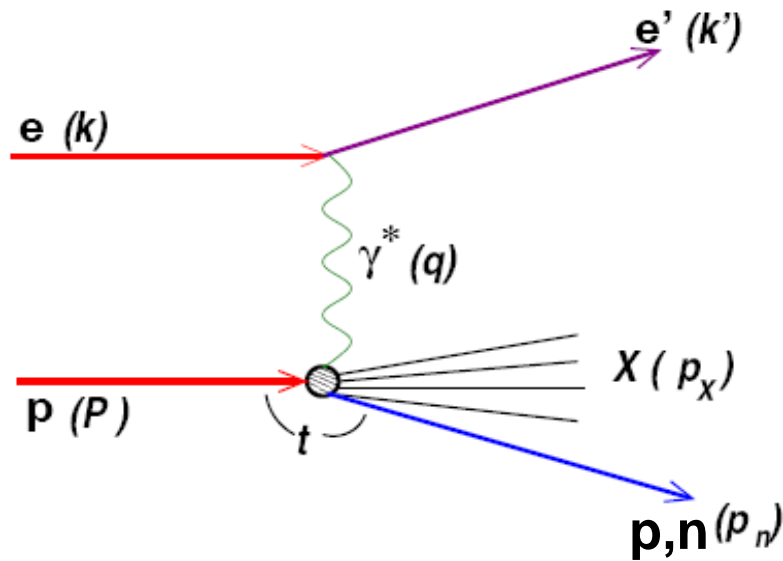
total lumi:  $0.5 \text{ fb}^{-1}$  per experiment



scale for secondary particle production decreases from  $Q^2$  in current region (or high  $P_T$  jets if  $Q^2 \sim 0$ ) to a soft hadronic scale (proton fragmentation region)

Significant fraction of  $ep$  scattering events contains in the final state a leading proton or neutron which carry a substantial portion of the energy of the incoming proton:  $e + p \rightarrow e' + n + X$  or  $e' + p + X$

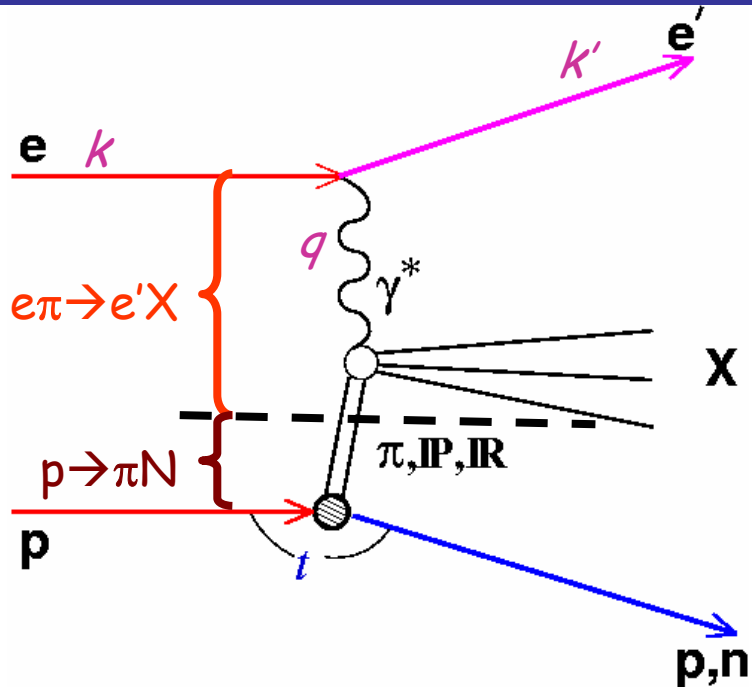
Production mechanism of leading baryons:



'conventional' fragmentation of proton remnant (e.g. Lund string)

exchange of virtual particle

- LP: neutral iso-scalar, iso-vector ( $\pi, IR, IP$ )
- LN: charged iso-vector ( $\pi^+, \rho^+, a_2 \dots$ )



$ep \rightarrow e'NX$

Lepton variables:

$$Q^2 = -(k - k')^2$$

$$x = Q^2 / (2p \cdot q)$$

Leading baryon variables:

$$x_L = E_{LB} / E_p$$

$$t = (p - p_{LB})^2 \quad (\text{or } p_{T,LB}^2)$$

In the exchange model the cross sections factorise, e.g. for one pion exchange

$$\sigma(ep \rightarrow e'NX) = f_{\pi/p}(x_L, t) \times \sigma(e\pi \rightarrow e'X)$$

$f_{\pi/p}(x_L, t)$  - **pion flux**:  
probability to emit pion from the photon with given  $x_L, t$

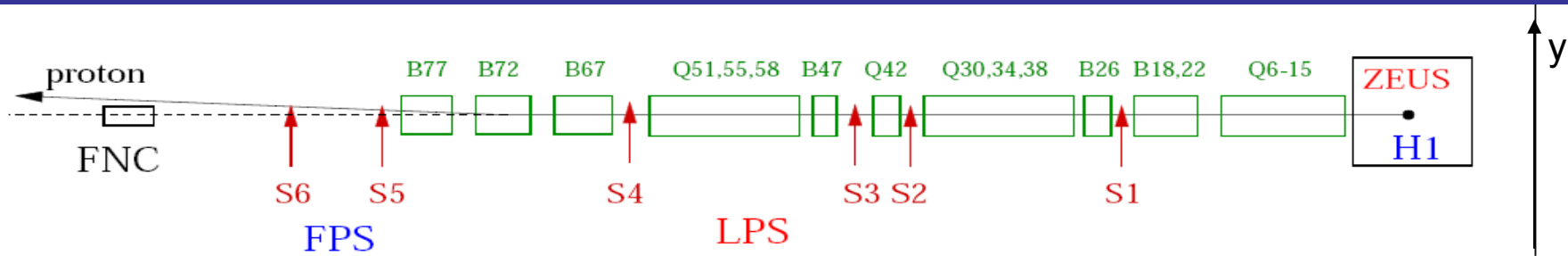
$\sigma(e\pi \rightarrow e'X)$  - **cross-section of  $e\pi$  scattering**

-LB production independent from photon vertex

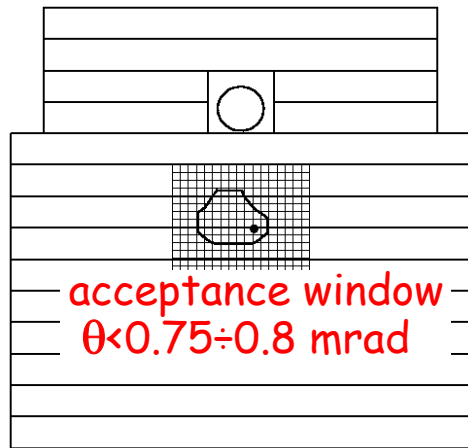
-**probe structure of exchanged particle**

-factorisation violation predicted- absorption/rescattering

# H1 and ZEUS detectors for leading baryons



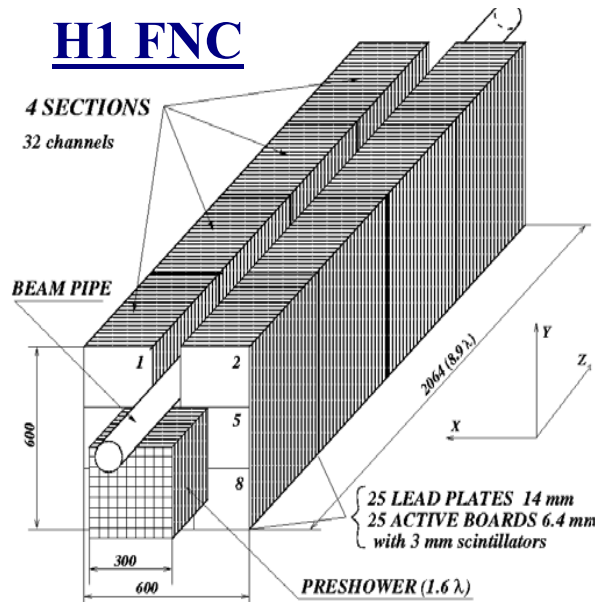
## ZEUS FNC+FNT



14 towers, 17x15 grid  
of the FNT hodoscopes,  
 $\sigma_E/E \approx 0.7/\sqrt{E}$

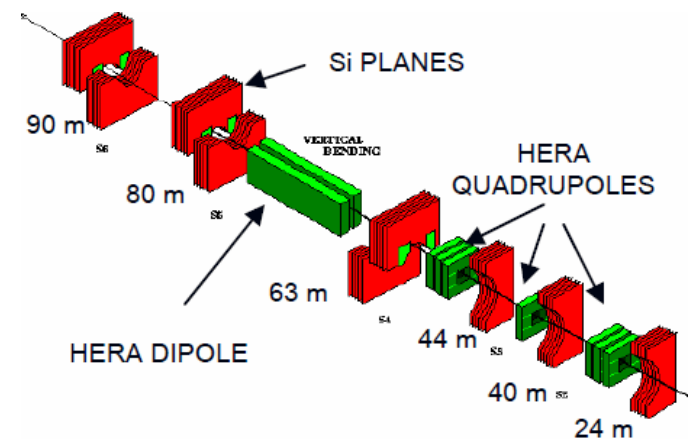
position resolution 2-3mm

## H1 FNC



$\sigma_E/E \approx 0.63/\sqrt{E} \oplus 2\%$

## ZEUS LPS

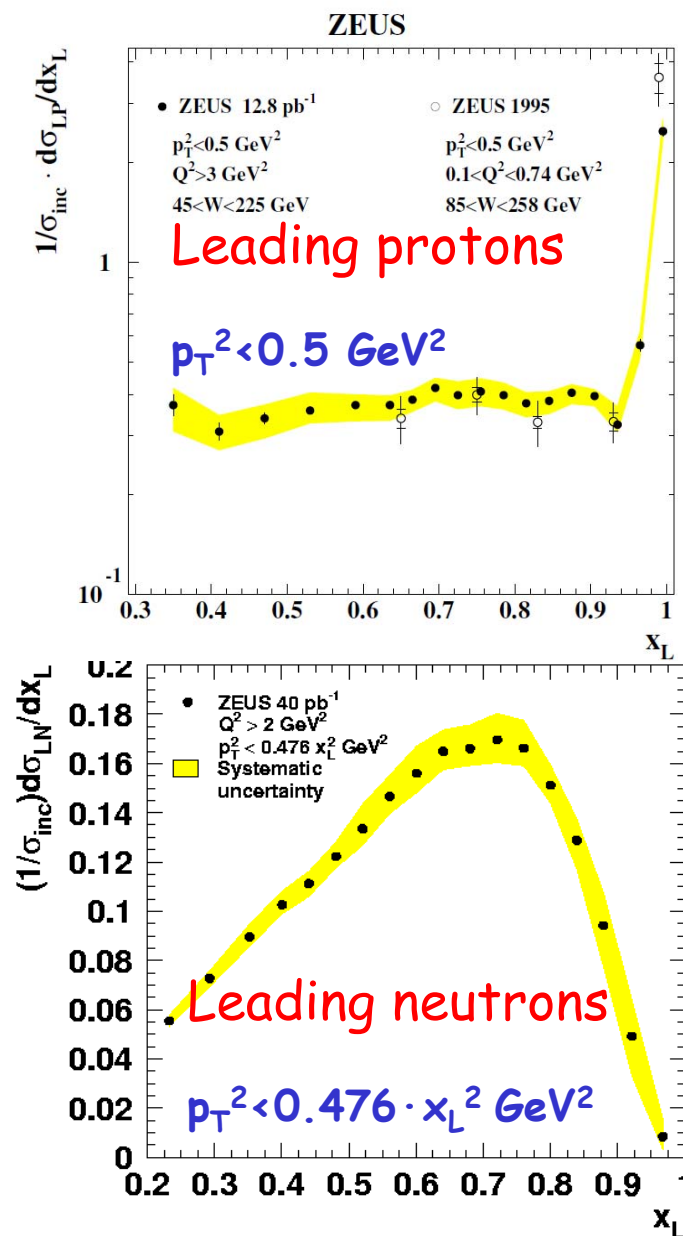


6 stations with  $\mu$ strip detectors  
hit position resolution  $\sim 30 \mu\text{m}$   
 $\sigma_{X_L} < 1\%$ ,  $\sigma_{P_T} \sim \text{few MeV}$

momentum accuracy  $< 1\%$

Acceptance limited by beam apertures and detector size  
 $p_T$  resolution is dominated by  $p_T$  spread of proton beam (50-100 MeV)

# Cross sections vs $x_L$ normalised to $\sigma_{DIS}$ ( $1/\sigma_{DIS} \cdot d\sigma/dx_L$ )



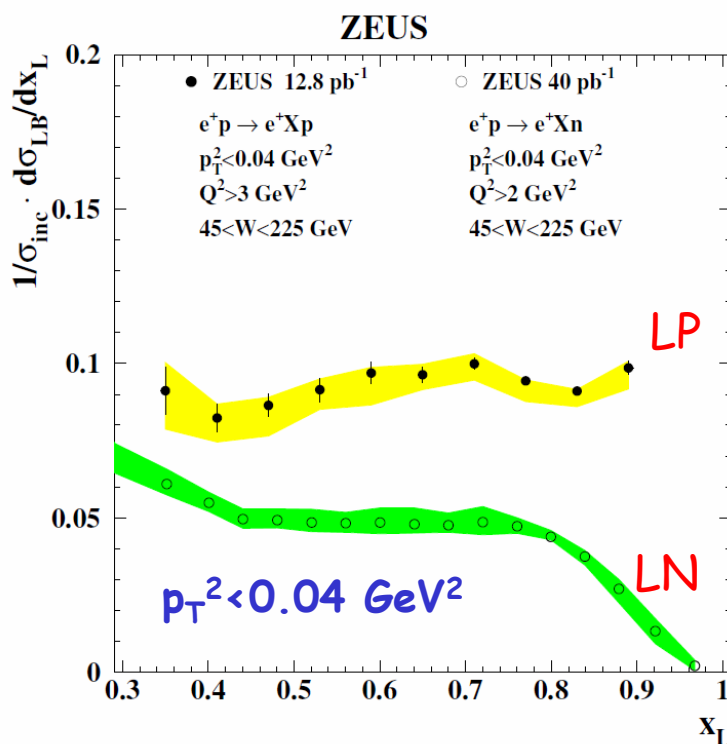
**Leading protons:** (DESY-08-176)

- diffractive peak at  $x_L = 1$  ; flat at  $x_L < 0.95$

**Leading neutrons:** (Nucl.Phys.B776(2007)1)

- yield  $\rightarrow 0$  as  $x_L \rightarrow 1$  ;
- drop at  $x_L < 0.7$  due to drop in acceptance

restrict to the same  $p_T^2$  range



- measurement:  
 $LP \sim 2 \cdot LN$

- for pure isovector particle exchange (e.g. pion) one expects  
 $LP = \frac{1}{2} \cdot LN$

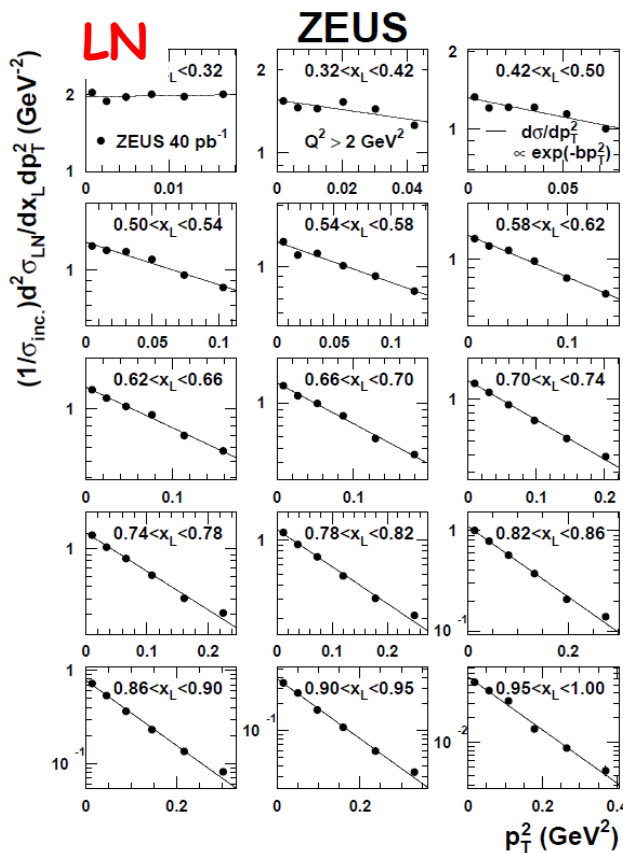
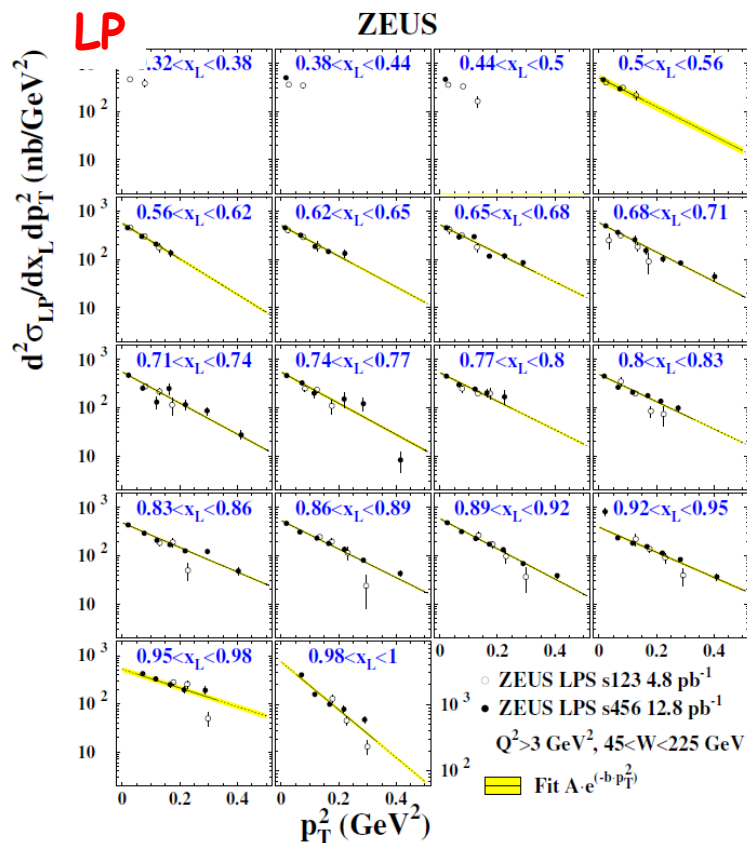
$\rightarrow$  more isoscalar exchanges contribute to the LP rates



# Double differential cross sections vs $p_T^2$ , $x_L$

DESY-08-176

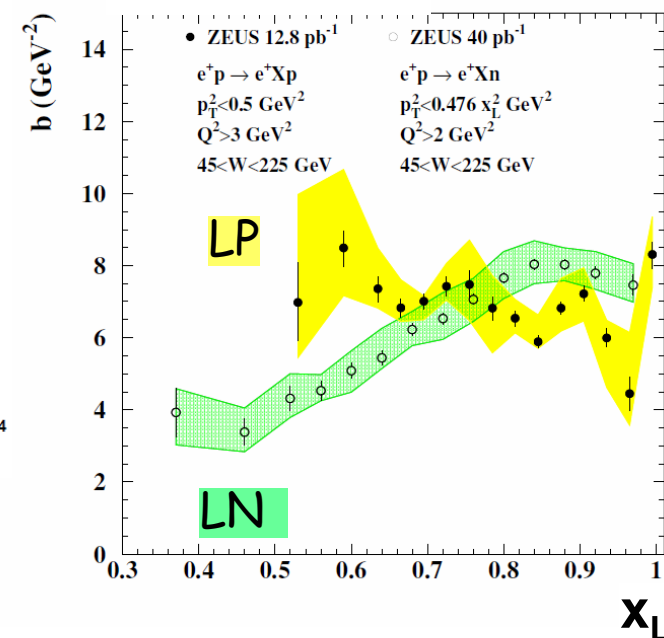
Nucl.Phys.B776(2007)1



Exponential behavior

$$\frac{d^2 \sigma}{dx_L dp_T^2} \sim a(x_L) e^{-b(x_L) p_T^2}$$

slope  $b(x_L)$

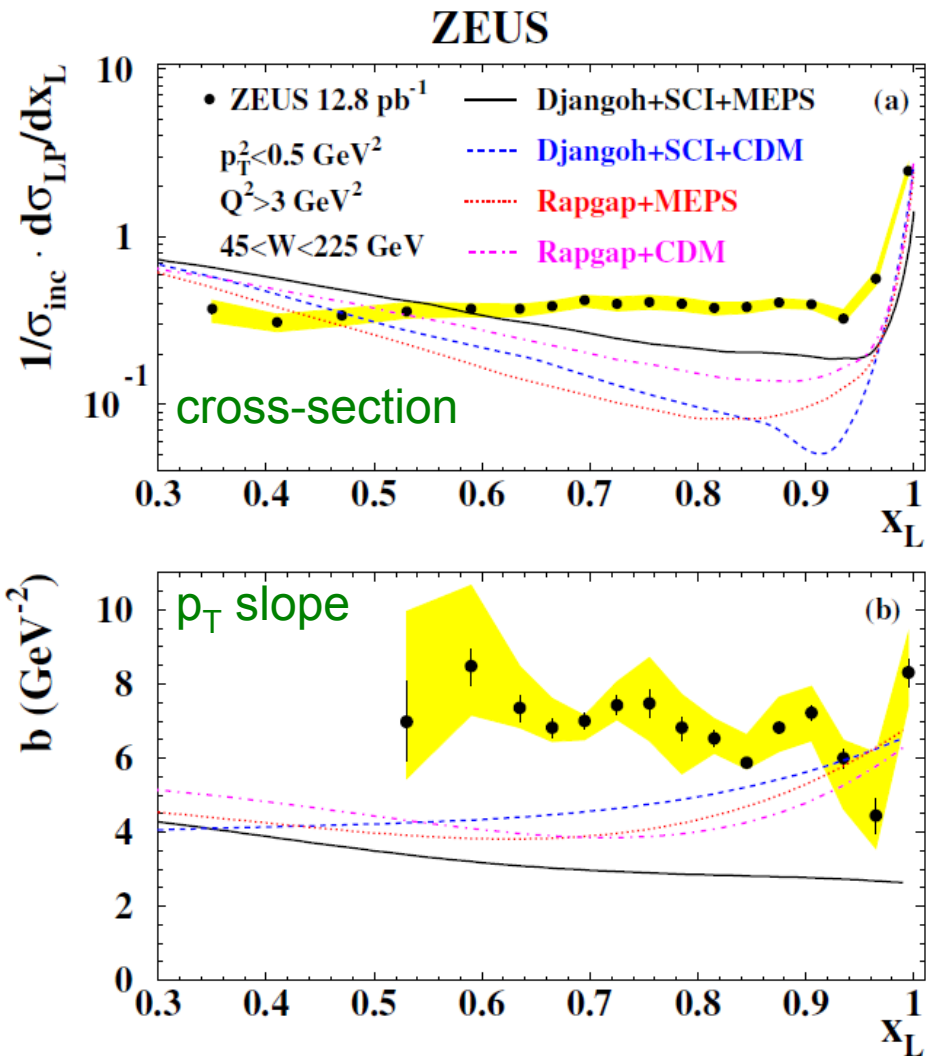


- different behavior for LP and LN
- similar around  $x_L \sim 0.7$

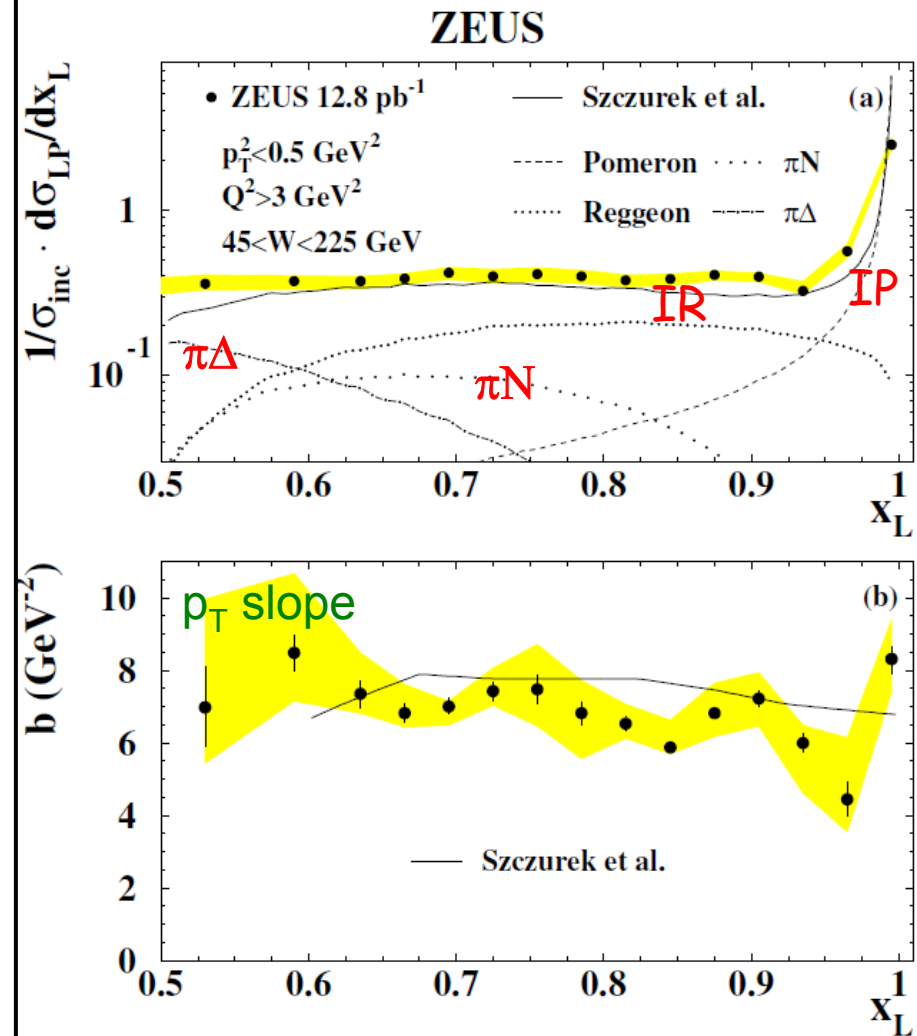


# Comparison with fragmentation and exchange models: Leading Protons in DIS

DESY-08-176



- standard fragmentation MC models don't describe the data out of the diffractive peak
- slopes too low at low  $x_L$

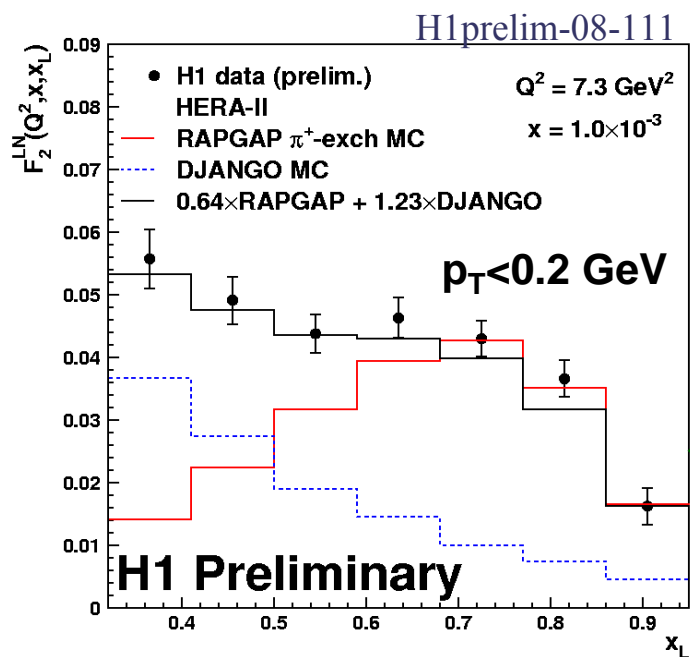
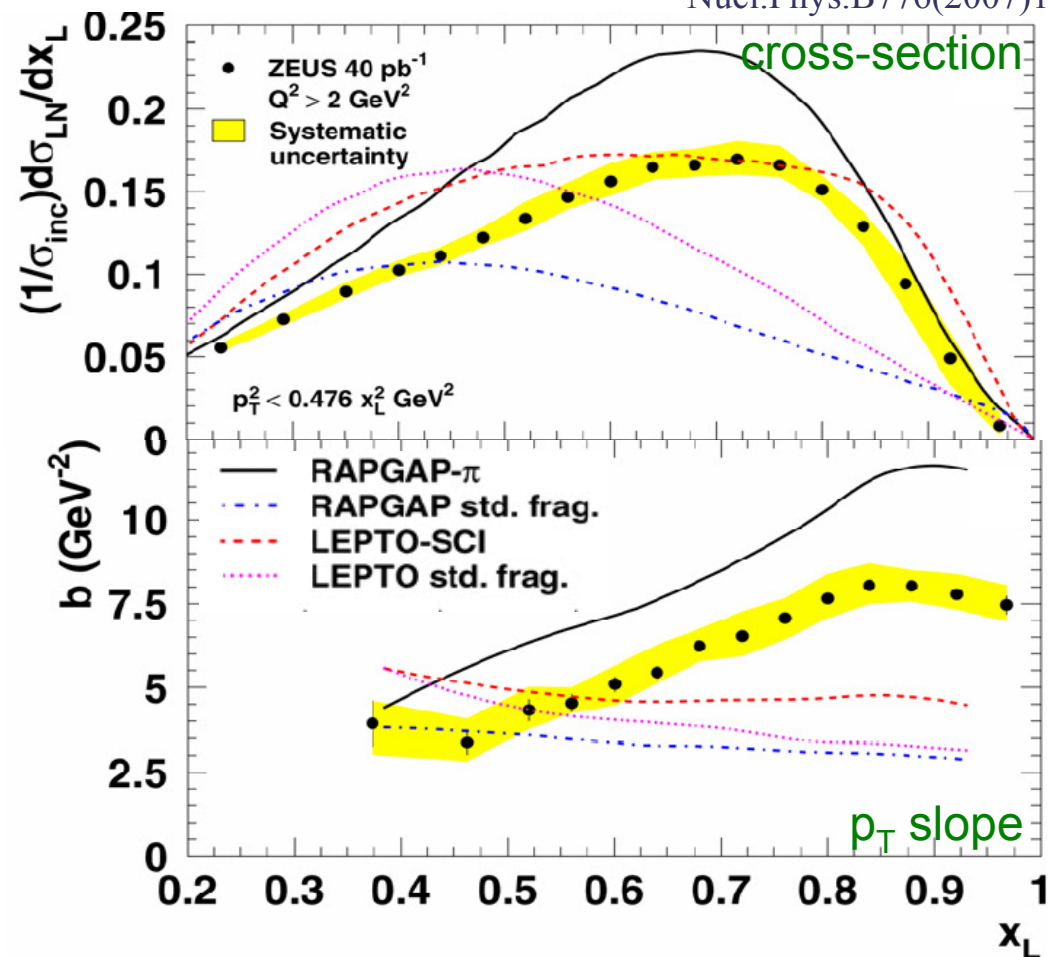


- good description by exchange models
- isoscalar reggeon dominant at intermediate  $x_L$

# Comparison with fragmentation and exchange models: Leading Neutrons in DIS

Nucl.Phys.B776(2007)1

- all standard fragmentation models underestimate the neutron yield at high  $x_L$
- LEPTO-SCI better for  $x_L$  shape, but not for the slope
- RAPGAP- $\pi$ -exchange describes data well for  $x_L > 0.6$ , underestimate data at lower  $x_L$
- Mixture of RAPGAP- $\pi$ -exchange and std. fragmentation (e.g. DJANGO-CDM) gives the best description of the data



←  $F_2^{LN}$  - reduced DIS cross section

$$\frac{d^3 \sigma(ep \rightarrow eNX)}{dQ^2 dx dx_L} = \frac{4 \pi \alpha^2}{x Q^4} \left[ 1 - y + \frac{y^2}{2} \right] \cdot F_2^{LN}(Q^2, x, x_L)$$

# Estimate the Pion structure function from $F_2^{LN}(Q^2, x, x_L)$

H1prelim-08-111

within  $\pi^+$ -exchange model we may try to estimate  $F_2^\pi$  from measured  $F_2^{LN}$ :

$$F_2^{LN(3)}(\beta, Q^2, x_L) = \Gamma_\pi(x_L) \cdot F_2^\pi(\beta, Q^2)$$

where

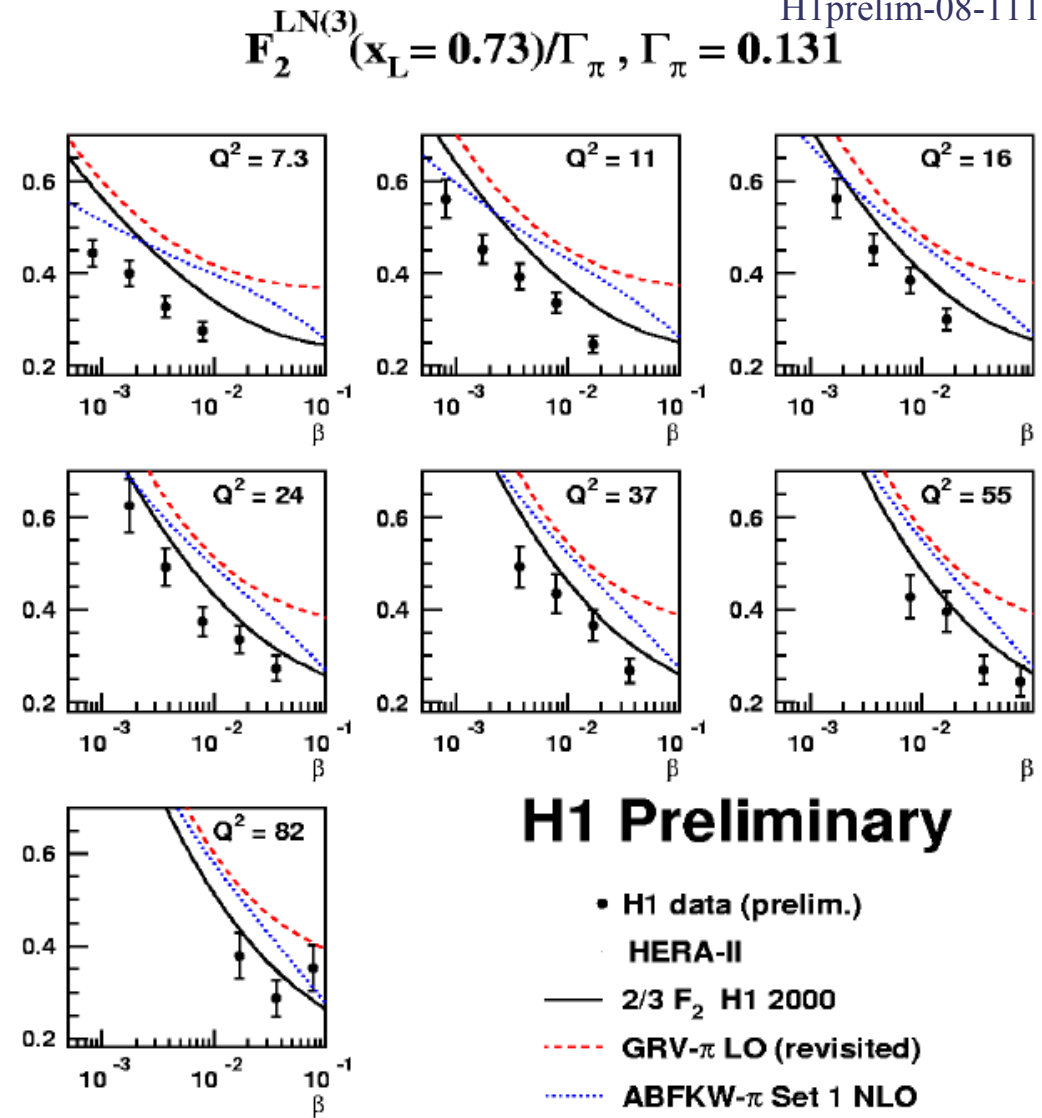
$\beta = x/(1-x_L)$  - fraction of pion momentum carried by struck quark (i.e.  $x_{Bj}$  for pion)

$\Gamma_\pi(x_L)$  is integrated over t pion flux

$$\Gamma_\pi = \int f_{\pi/p}(x_L = 0.73, t) dt$$

use pion flux parameterisation (Holtmann et al.):

$$f_{\pi^+/p} = \frac{1}{2\pi} \frac{g_{p\pi\pi}^2}{4\pi} (1-x_L) \frac{-t}{(m_\pi^2 - t)^2} \cdot \exp\left(-R_{\pi n}^2 \frac{m_\pi^2 - t}{1-x_L}\right)$$



Data are sensitive to the parameterisations of the pion structure function (constrained for  $x > 0.1$  from the fixed target experiments).

# Estimate the Pion structure function from $F_2^{LN}(Q^2, x, x_L)$

H1prelim-08-111

within  $\pi$ -exchange model we can estimate  $F_2^\pi$  from measured  $F_2^{LN}$ :

$$F_2^{LN(3)}(\beta, Q^2, x_L) = \Gamma_\pi(x_L) \cdot F_2^\pi(\beta, Q^2)$$

where

$$\beta = x/(1-x_L)$$

$\Gamma_\pi(x_L)$  is integrated over  $\dagger$  pion flux

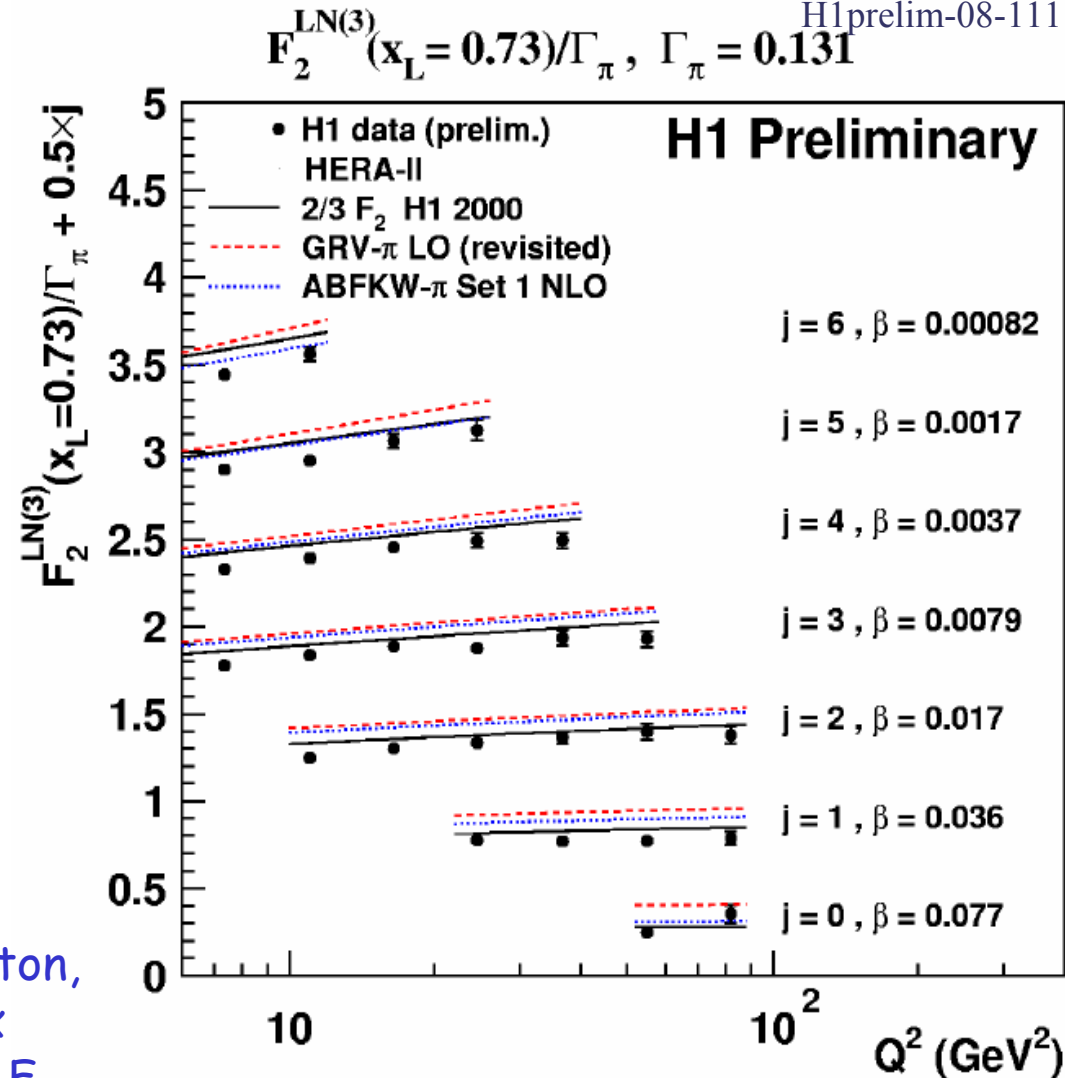
$$\Gamma_\pi = \int f_{\pi/p}(x_L = 0.73, t) dt$$

use pion flux expression (Holtmann et al.):

$$f_{\pi^+/p} = \frac{1}{2\pi} \frac{g_{p\pi\pi}^2}{4\pi} (1-x_L) \frac{-\dagger}{(m_\pi^2 - \dagger)^2} \cdot \exp\left(-R_{\pi n}^2 \frac{m_\pi^2 - \dagger}{1-x_L}\right)$$

- $F_2^{LN}$  dependence on  $x$  and  $Q^2$  similar to proton,  $\rightarrow$  universality of hadron structure at low  $x$
- in absolute values  $F_2^{LN}/\Gamma$  below the  $F_2^\pi$  and  $F_2$

However: large uncertainty of pion flux normalisation: choice of pion flux (formfactor), absorption/rescattering, background...



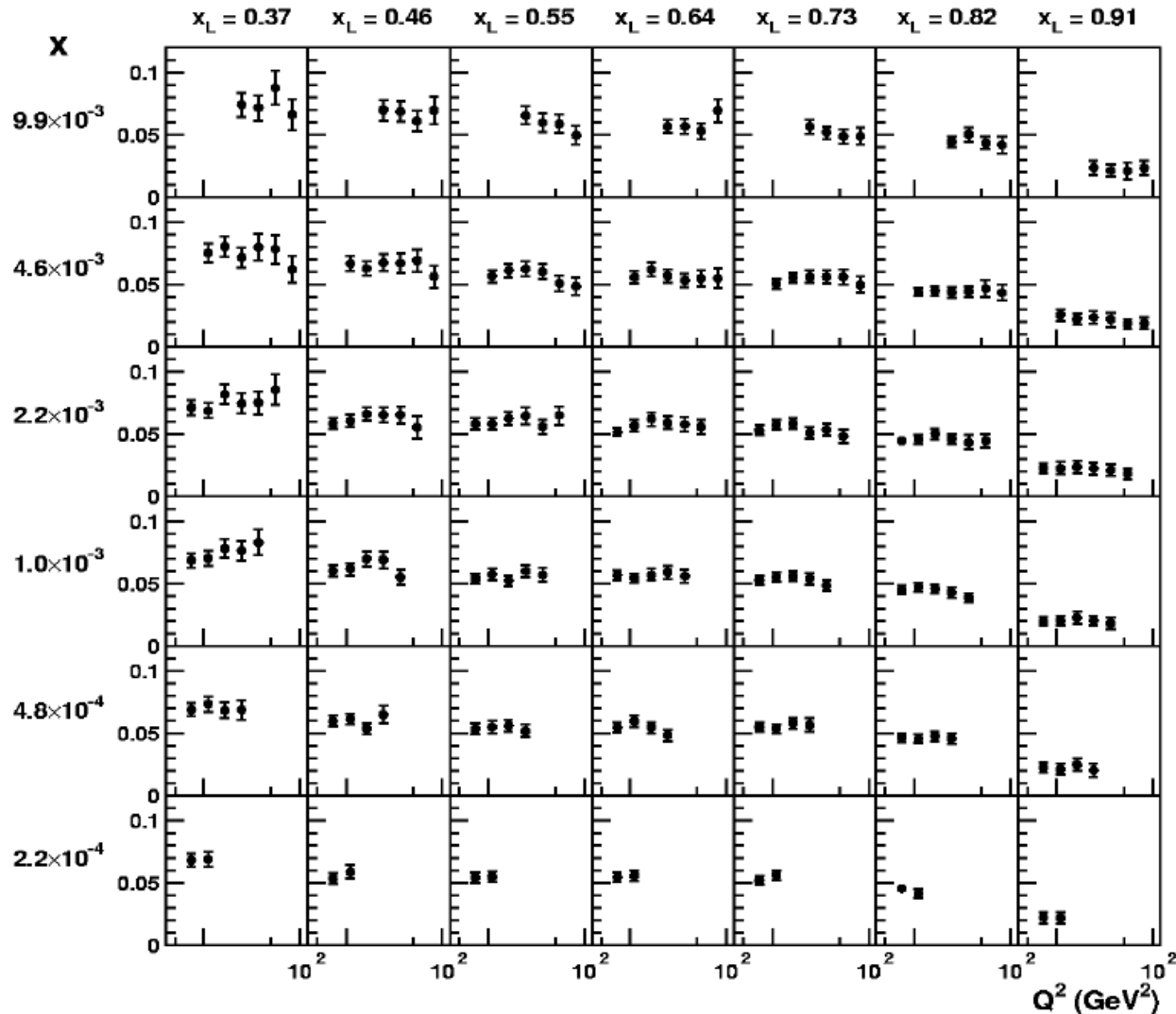
# $F_2^{LN}(Q^2, x, x_L)$ to $F_2(Q^2, x)$ ratio

$6 < Q^2 < 100 \text{ GeV}^2$ ,  $p_T < 0.2 \text{ GeV}$

$$\frac{d^3\sigma(ep \rightarrow eNX)}{dQ^2 dx dx_L} = \frac{4\pi\alpha^2}{xQ^4} \left[ 1 - y + \frac{y^2}{2} \right] F_2^{LN}(Q^2, x, x_L)$$

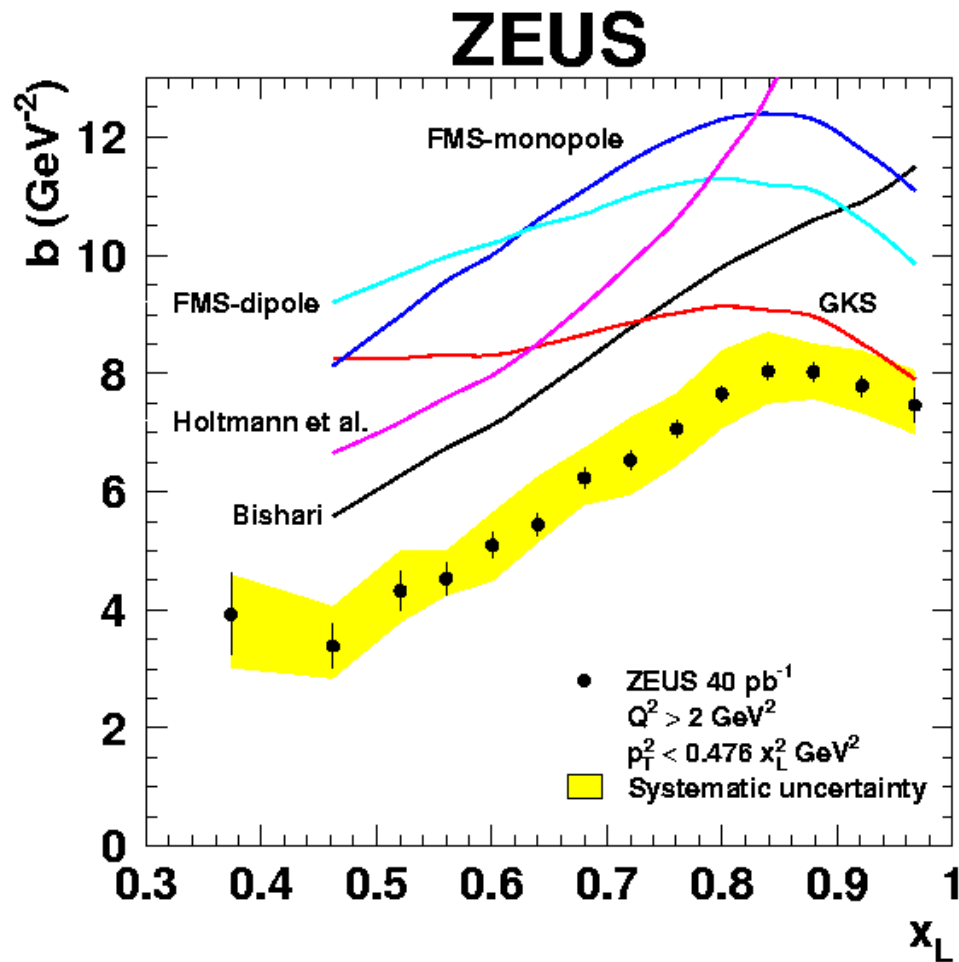
$F_2(Q^2, x)$  from the H1-2000-PDF parameterisation

$F_2^{LN}(Q^2, x, x_L)/F_2(Q^2, x)$  H1 Preliminary (HERA-II)



$F_2^{LN}(Q^2, x, x_L)/F_2(Q^2, x)$  is mostly flat in  $Q^2$  and  $x$

i.e. LN production rate, kinematics is approx. independent of  $(Q^2, x)$   
 → consistent with factorisation, limiting fragmentation (overall suppression of events is also possible)



in  $\pi$ -exchange

$$\sigma(ep \rightarrow e'nX) = f_{\pi/p}(x_L, t) \times \sigma(e\pi^+ \rightarrow e'X)$$

- $p_T^2$  (or  $t$ ) distribution is determined solely by pion flux

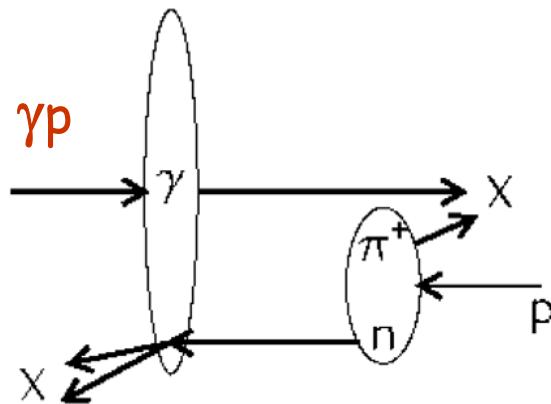
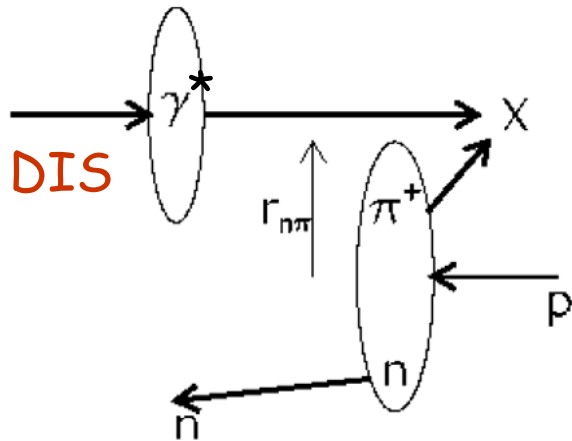
$$f_{\pi/p} = \frac{1}{2\pi} \frac{g_{p\pi\pi}^2}{4\pi} (1-x_L)^{1-2a(t)} \frac{-t}{(m_\pi^2 - t)^2} \cdot |F(x_L, t)|^2$$

- many parameterizations of pion flux  $f_{\pi/p}(x_L, t)$  in literature
- compare measured  $p_T$  slope  $b(x_L)$  with models (shown best agreeing models)
- reasonable agreement in shape but not in absolute values: all give too large  $b(x_L)$
- $\pi$ -exchange models alone don't describe  $p_T^2$ -distribution



# Exchange model refinement: absorptive corrections

Absorption: important ingredient to interpret the results in terms of particle exchange



Neutron absorption through rescattering:

enhanced when  $\pi$ -n system size  $r_{\pi n} \sim 1/p_T$  is small w.r.t. the  $\gamma$ -transverse size, e.g. at high  $p_T$ , low  $x_L$   
→ neutron breaks up or  
→ is kicked to lower  $x_L$ , higher  $p_T$  (migration) and/or escapes detector acceptance (absorption loss)  
(in other language: multi-Pomeron exchange)

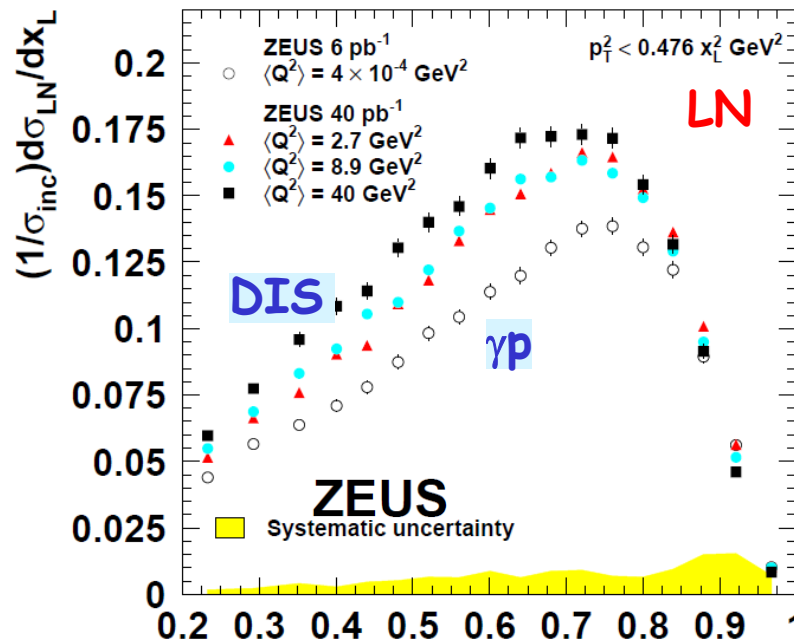
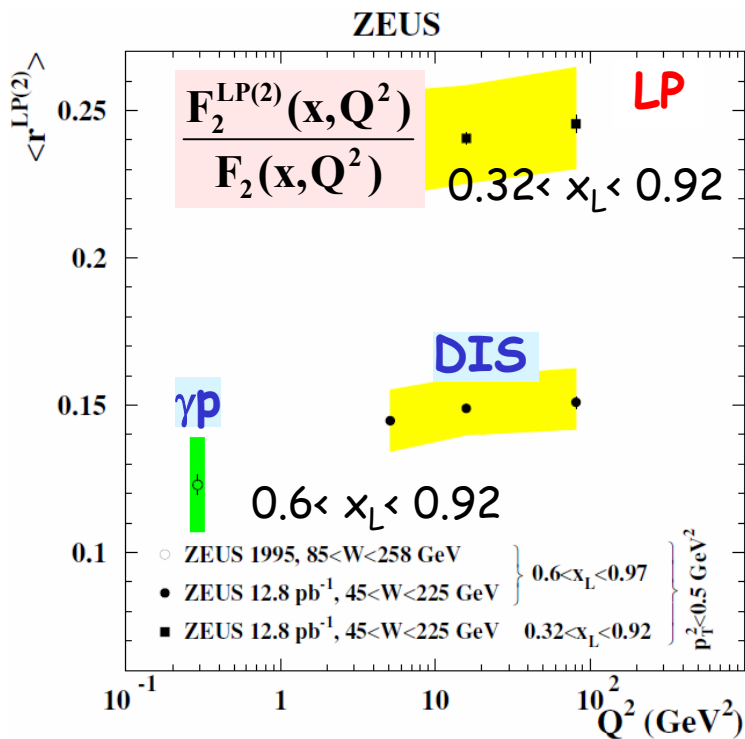
- Affects the relative rate of leading neutrons (depends on the scale  $Q$ )  
more absorption in photoproduction than in DIS, (real  $\gamma$  transverse size larger than at higher  $Q^2$ )  
→ The calculations/models made without absorption may overestimate the measurements

Effects of absorption and migration estimated:  
D'Alesio, Pirner; Nikolaev, Speth, Zakharov;  
Kaidalov, Khoze, Martin, Ryskin;  
Kopeliovich, Potashnikova, Schmidt, Soffer

Absorption- key ingredient in calculations of gap-survival probability in pp interactions at LHC, critical in interpreting hard diffractive processes, e.g. central exclusive Higgs prod.



# Comparison $\gamma p$ /DIS: $Q^2$ dependence



increase of LP and LN rate with  $Q^2$

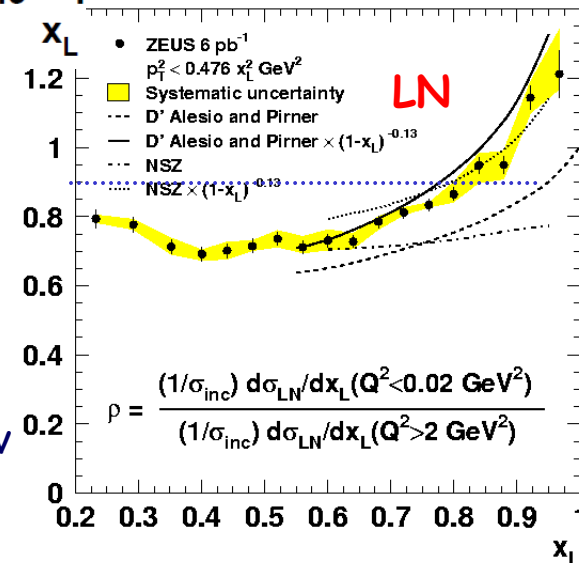
Suggest violation of vertex factorisation

From simple geometrical picture:  
smaller  $\gamma^*$  transverse size at higher  $Q^2$   
 $\rightarrow$  less absorption  $\rightarrow$  larger event yield

Ratio  $\gamma p$ /DIS:  
absorption models  
describe the data

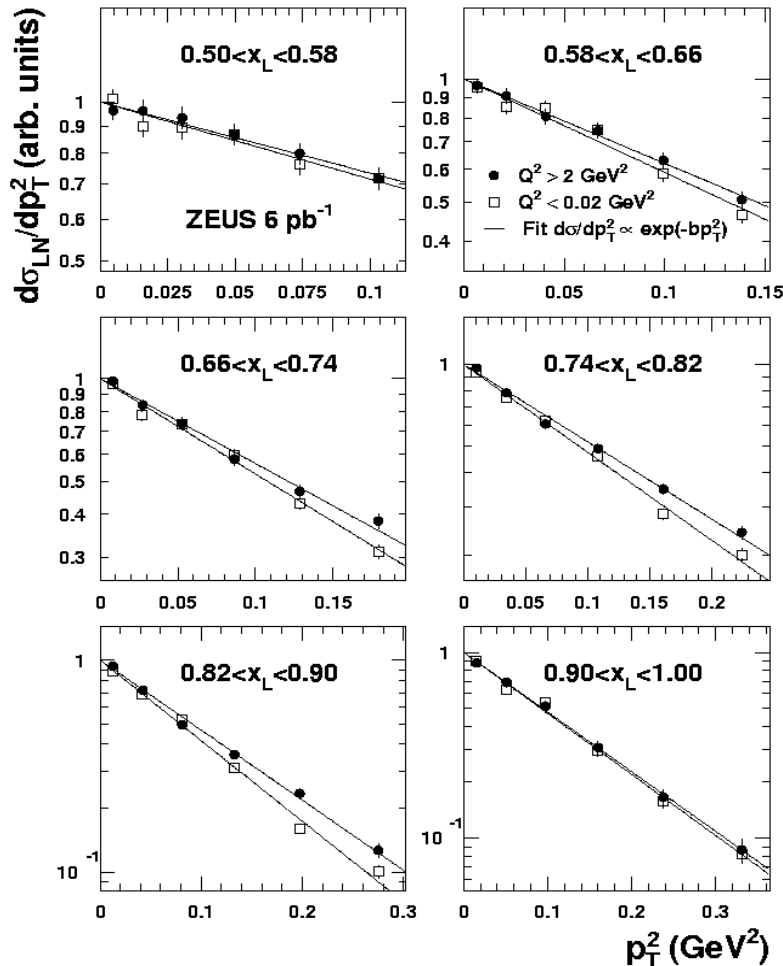
-d'Alesio, Pirner  
geometrical model

-Nikolaev, Speth, Zakharov  
Regge based model with  
multipomeron exchanges



# Comparison $\gamma p$ /DIS: $p_T^2$ distributions (LN)

ZEUS



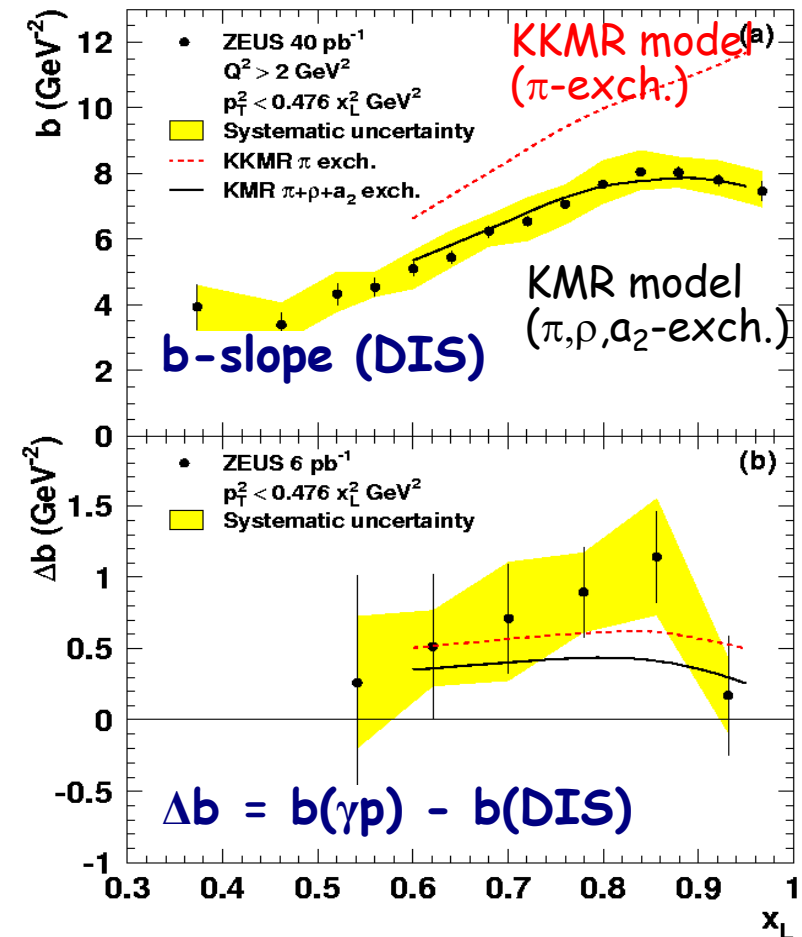
$p_T^2$  slopes at  $\gamma p$  steeper than at DIS

From simple geometrical picture:

Larger  $p_T \rightarrow$  smaller  $r_{\pi n} \rightarrow$  more absorption  
 $\rightarrow$  less neutrons at high  $p_T \rightarrow$  steeper slope

model of Kaidalov, Khoze, Martin, Ryskin

- rescattering on intermediate partons in central rapidity region; migration of LN in ( $x_L, p_T$ )
- ~50% absorption loss in  $\gamma p$
- addition of ( $\rho, a_2$ ) exchanges



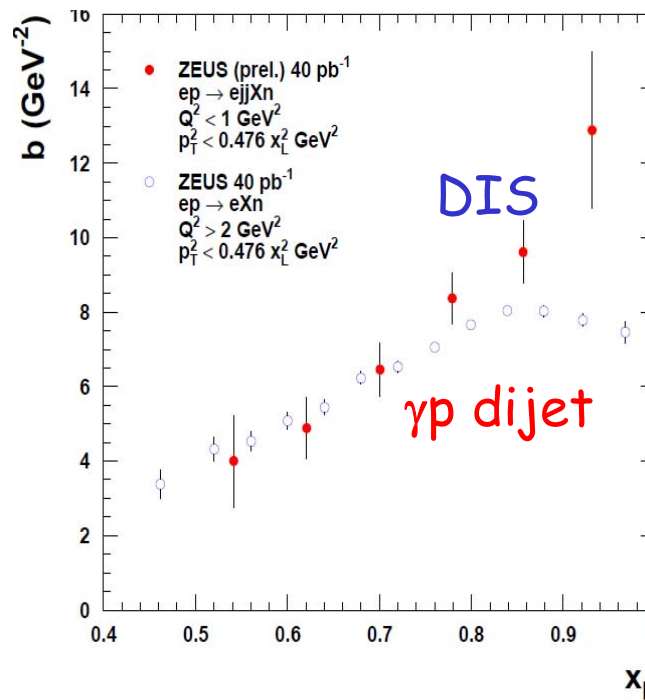
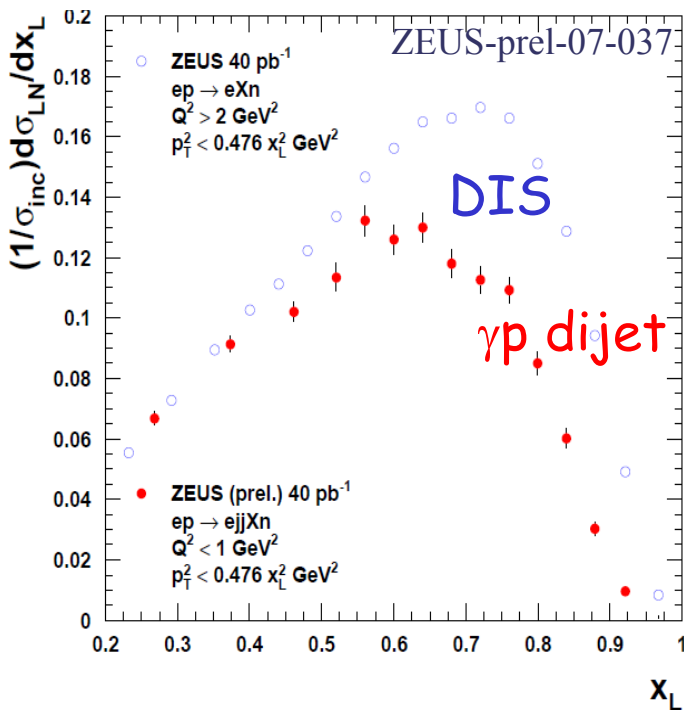
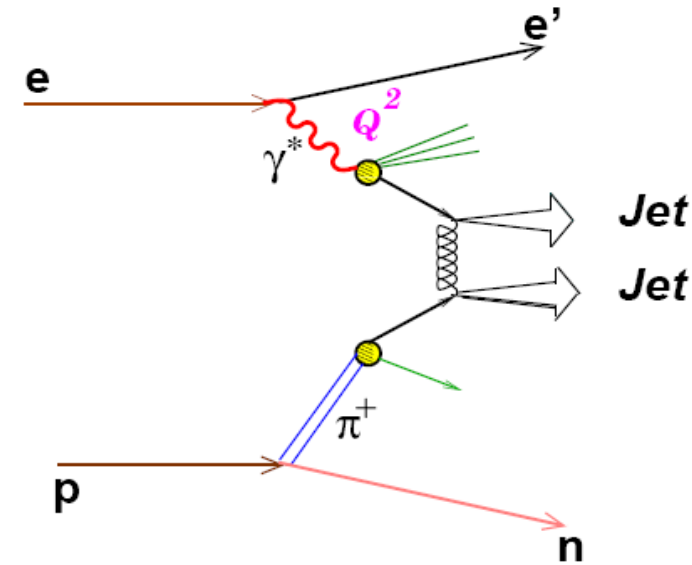
# Dijet photoproduction with LN

Study the jet production in event with leading neutron in the final state ( $\gamma^* p \rightarrow \text{jet} + \text{jet} + n + X$ )

In photoproduction ( $Q^2 \sim 0$ ) hard scale provided by jets with high  $E_T^{\text{jet}}$

Do we see suppression ?

Recall: factorisation breaking observed in diffractive dijet photoproduction at HERA: in  $\gamma p$  jets with  $E_T^{\text{jet}} > 5 \text{ GeV}$  are suppressed by factor 2 compared to DIS)



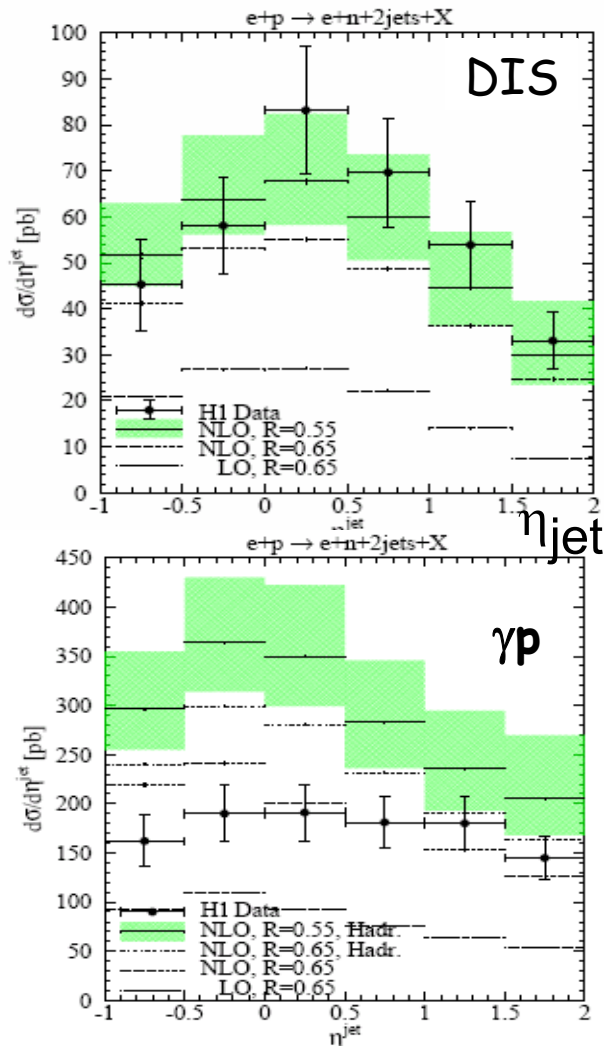
- $x_L$  distribution shows suppression at high  $x_L$   
→ phase space limitation (dijets in the final state leave little room for energetic neutrons)

- b-slopes (DIS &  $\gamma p$ -dijets) are slightly different at high  $x_L$

# Comparison of $\gamma^*p \rightarrow 2\text{jets} + n + X$ with theory

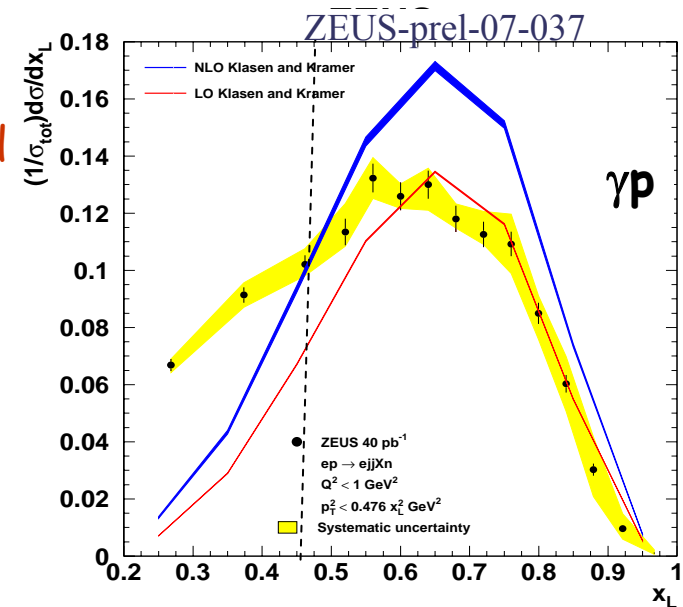
Calculations (Klasen & Kramer, Eur.Phys.J.C49:957-965,2007)

- normalise NLO (fix pion PDF, adjust pion flux) to H1-DIS data  $\gamma^*p \rightarrow jj+n+X$  (Eur. Phys. J. C41 (2005) 27)
- compare to H1- $\gamma p$  data ( $\gamma p \rightarrow jjnX$ ), look for suppression

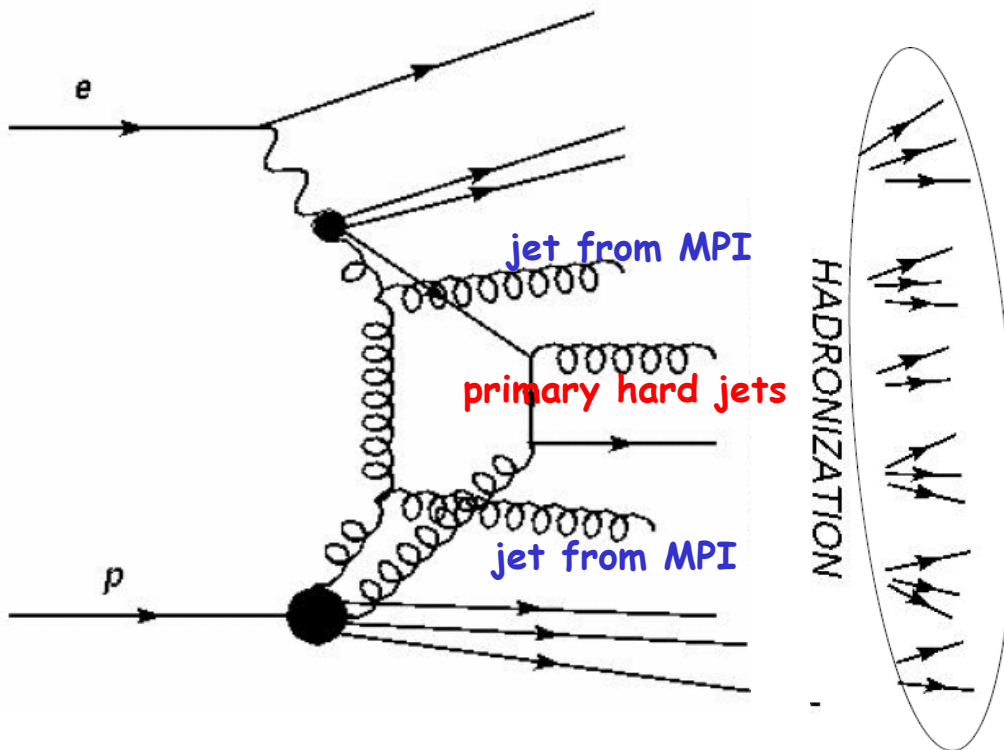


NLO vs H1 photoproduction data ( $E_T^{\text{jet}} > 7 \text{ GeV}$ ) needs  $\sim 0.48$  suppression of resolved component (or 0.64 global suppression)

NLO overestimates also the ZEUS  $\gamma p$  dijet data ( $E_T^{\text{jet}} > 7.5 \text{ GeV}$ ) for  $x_L > 0.5$



# Underlying event / Multi-parton interactions



In addition to the primary hard parton-parton interaction with large  $p_T$ :

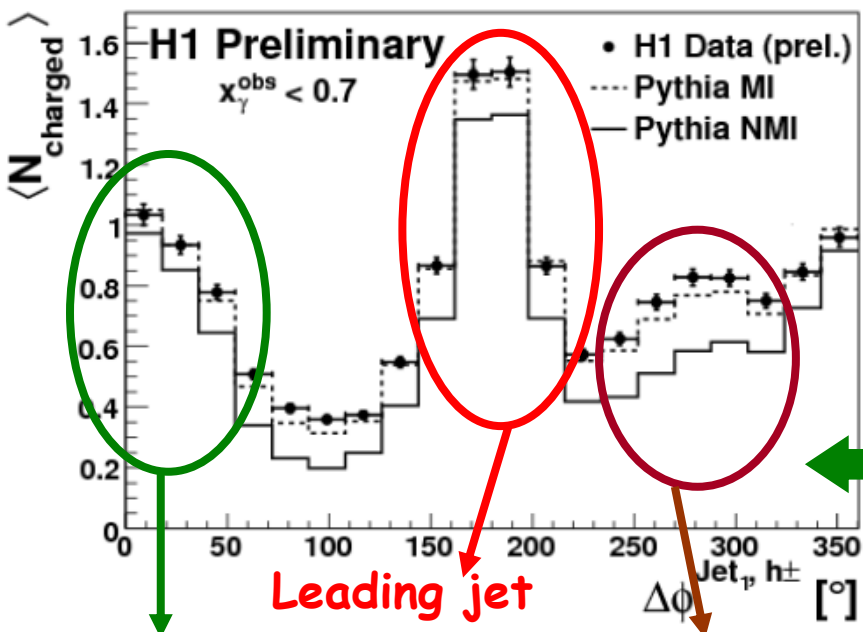
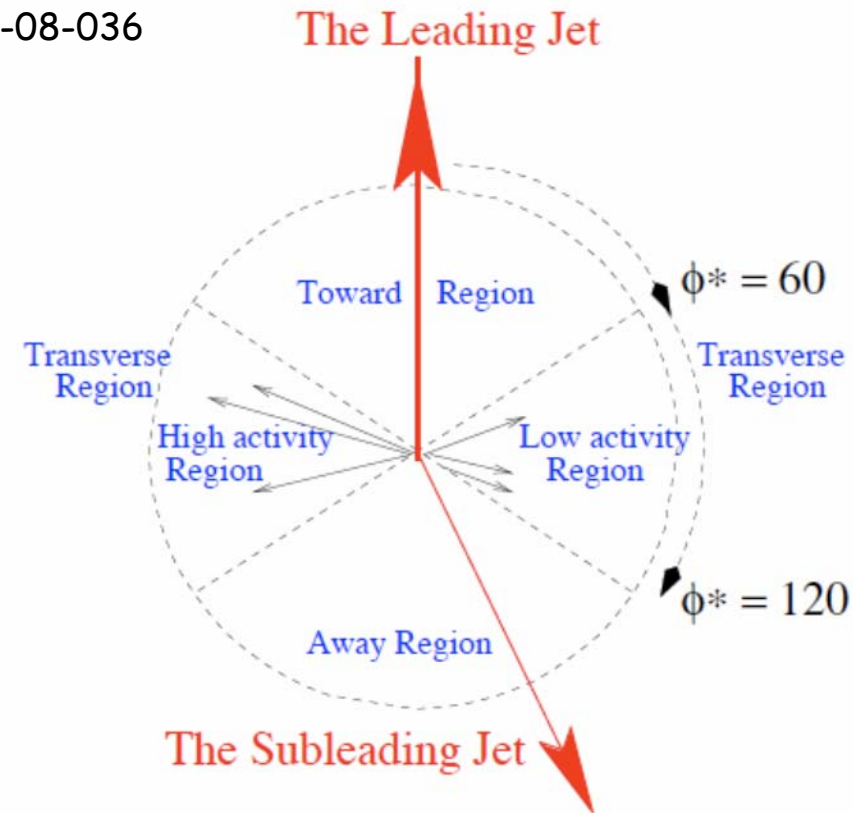
- interactions with lower  $p_T$  (remnant interactions)
- additional hard parton-parton interactions
  - higher particle multiplicity, jet multiplicity, energy offset
  - important for analyses involving jets (pedestal under the jet) !
  - may fake a discovery signal !
- also to take into account
  - higher order QCD corrections (e.g. parton showers)
  - effects of fragmentation
  - beam remnants

**Understanding and modeling of Underlying event and multi-parton interactions crucial for all precision measurements!**

# MPI studies in photoproduction ( $Q^2 < 10^{-2} \text{ GeV}^2$ )

H1-prelim-08-036

- 2-jets with  $P_{T}^{\text{jet}} > 5 \text{ GeV}$  and  $|\eta^{\text{jet}}| < 1.5$
- charged particles with  $P_T > 150 \text{ MeV}$  and  $|\eta| < 1.5$
- define two transverse regions:
  - high activity region with higher  $P_T^{\text{sum}} = \sum_i^{\text{tracksp}} P_{T,i}$
  - low activity region



Charged particle multiplicity vs angle between leading jet and charged particles

Away region

High activity transverse region

measure charged particle multiplicity  
 → expect sensitivity to MPI in transverse regions

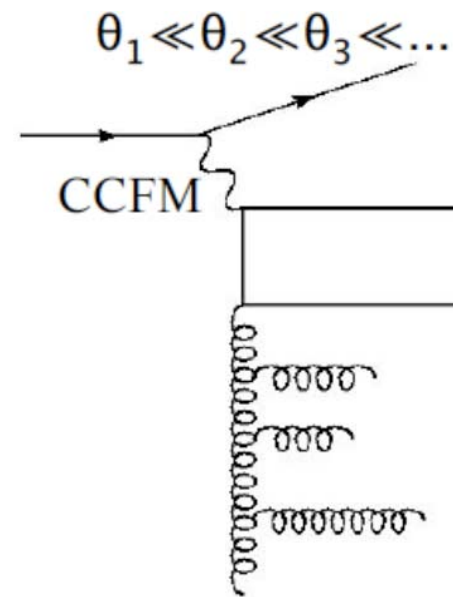
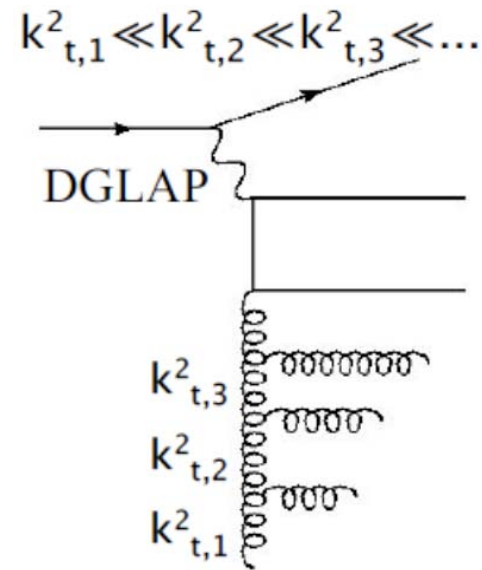


- **Pythia**

- direct+resolved processes in LO
- matched DGLAP parton shower
- with/without MPI (additional "semi-hard" interactions down to  $P_{T,min}=1.2 \text{ GeV}$ )

- **Cascade**

- off-shell LO ME for direct processes
- matched with CCFM parton showers
- $k_+$  un-integrated gluon densities (set1 , set2)
- no resolved (VDM) photon, no MPI

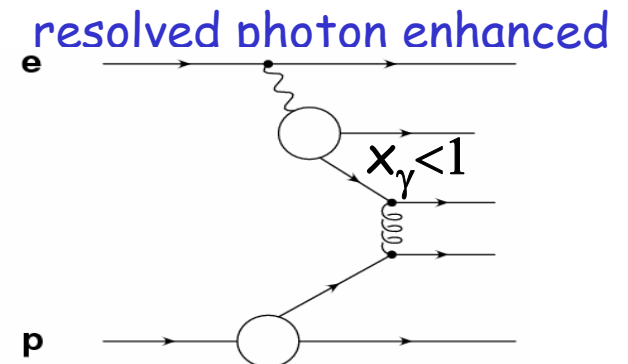
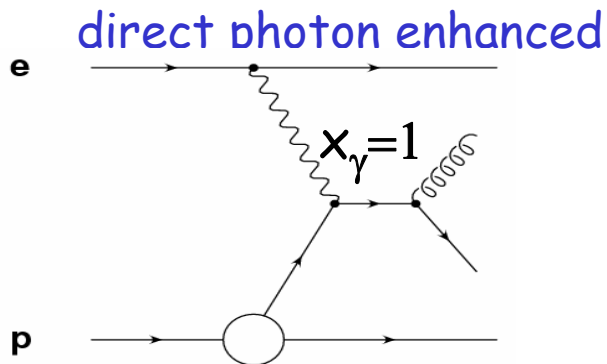
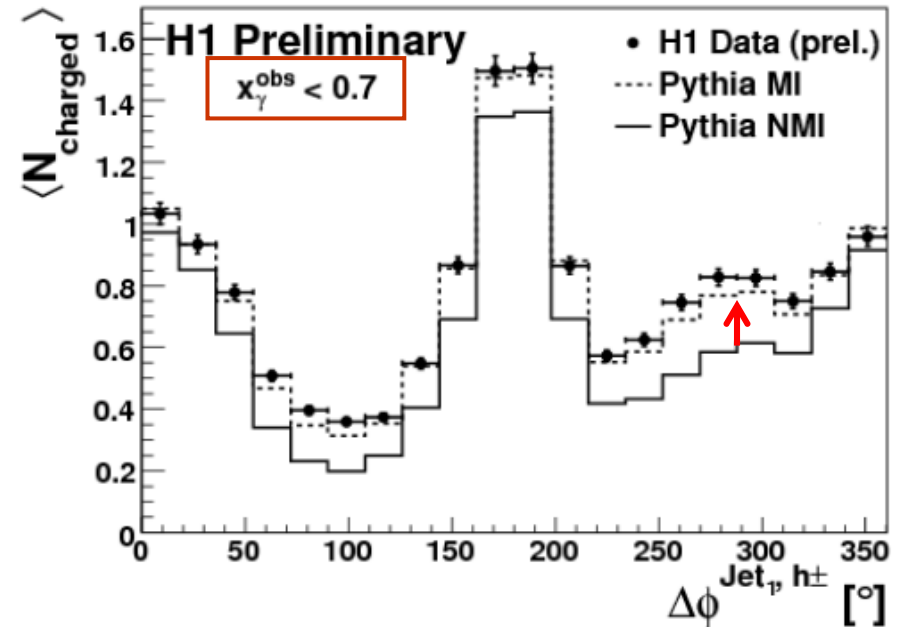
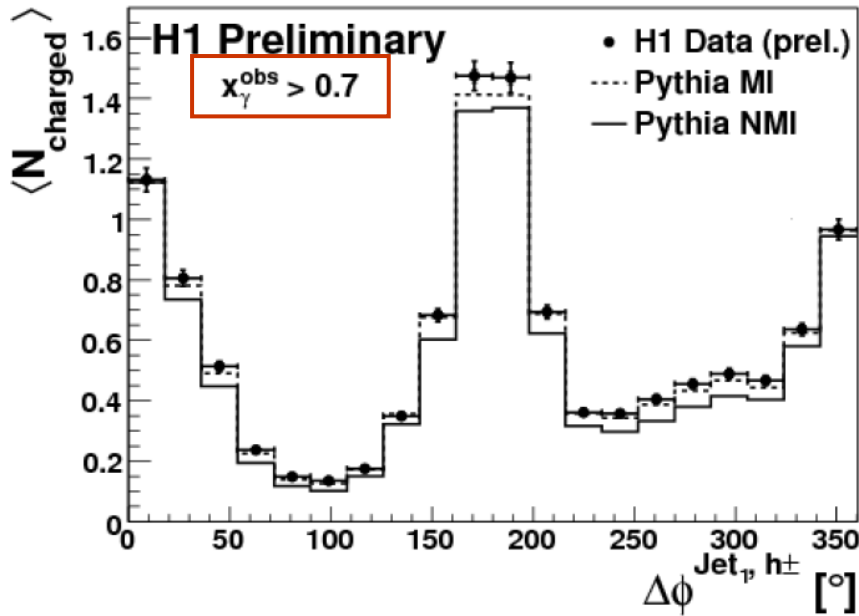




# Charged particle multiplicity vs $\Delta\phi$

$\Delta\phi$  - angle between leading jet and charged particles

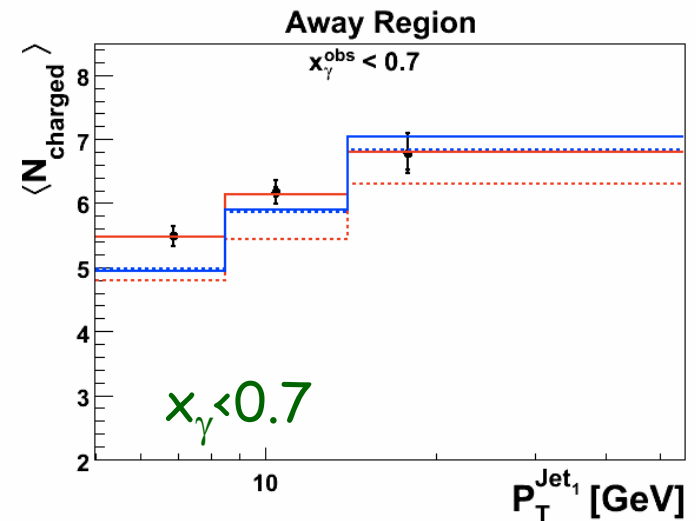
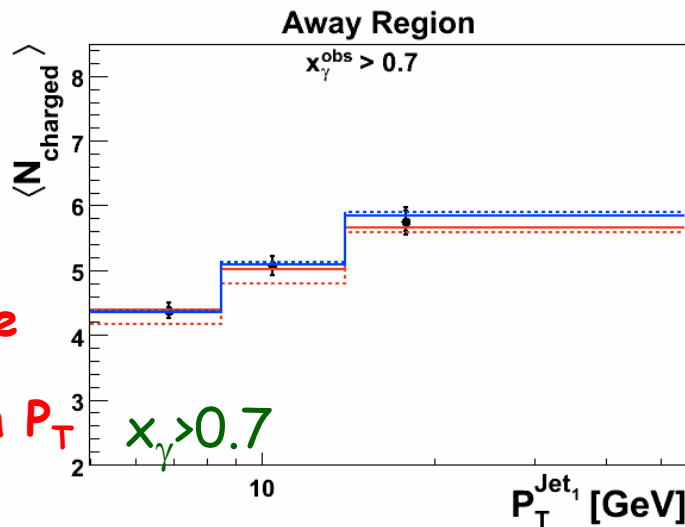
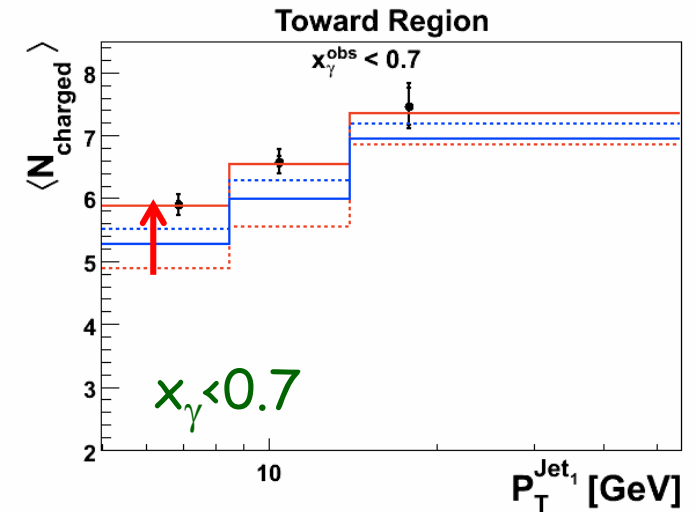
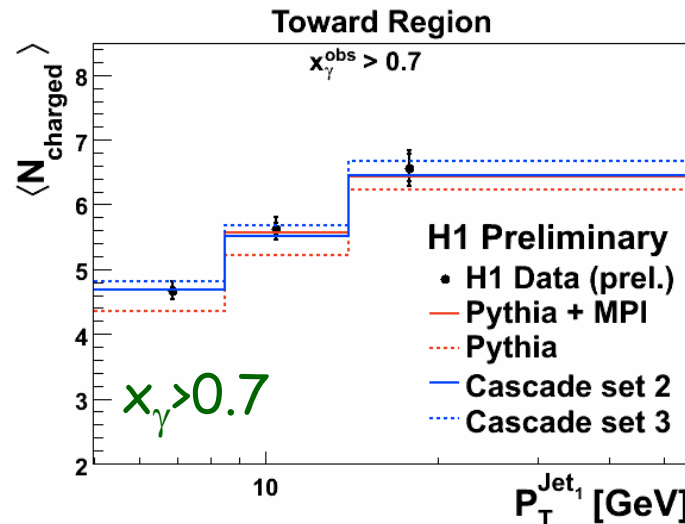
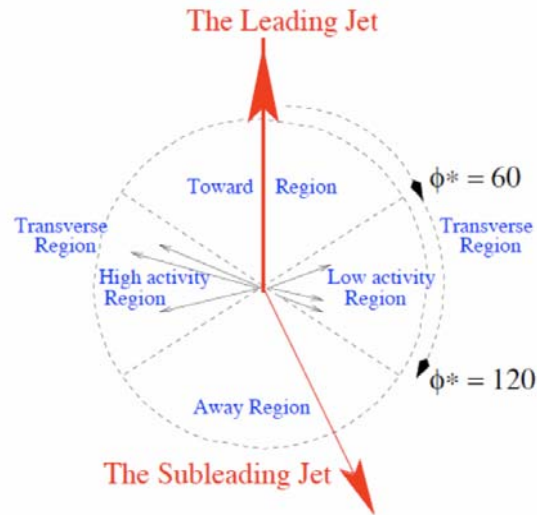
H1-prelim-08-036



Data at low  $x_\gamma$  described by Pythia when multi-parton interactions are included

# Charged particle multiplicity in Toward and Away regions

H1-prelim-08-036



models are similar  
at high  $x_\gamma$   
higher multiplicity  
at  $x_\gamma < 0.7$

MPI in Pythia contribute  
more at low  $P_T^{\text{jet}1}$

'pedestal' decrease with  $P_T$

CASCADE (no MPI) not too bad

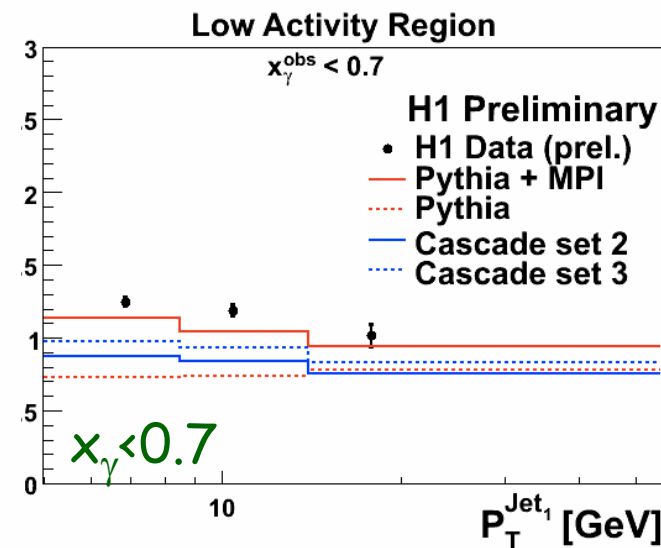
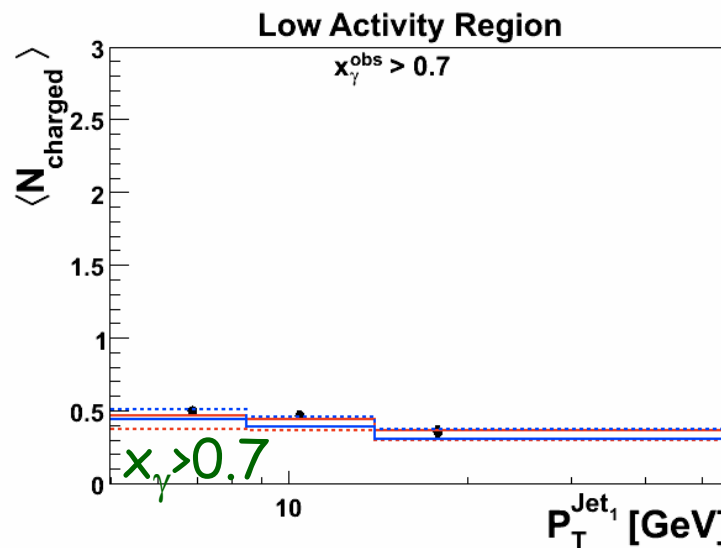
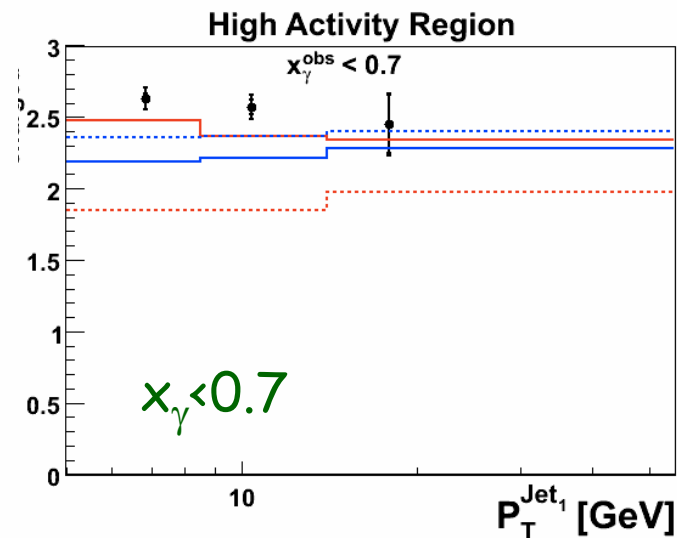
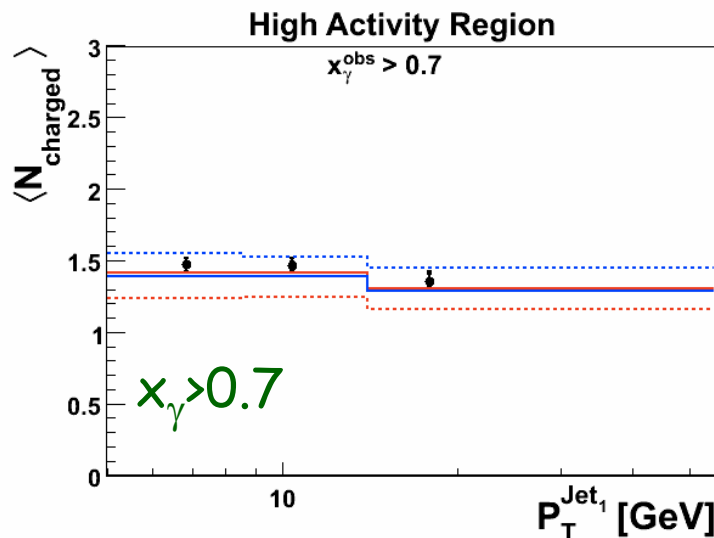
# Charged particle multiplicity in transverse regions

models are similar  
at high  $x_\gamma$

higher multiplicity  
at  $x_\gamma < 0.7$

**PYTHIA+MPI**  
provides reasonable  
description of data;

CASCADE is not too  
bad (worse than  
PYTHIA with MPI but  
better than PYTHIA  
without MPI)



**Data described by Pythia only when  
multi-parton interactions are included**

# Summary

## Leading Baryons are good ground to study soft vs. hard physics

- Precise measurements of LB  $x_L$  and  $p_T^2$  presented in  $\gamma p$ , DIS,  $\gamma p$  with dijets;
- Standard fragmentation models do not describe the data;  
Models with virtual particle exchange describe data better;
- For LN production pion structure function estimated, compared with parameterisations
- $F_2^{LN}/F_2$  ratio is mostly independent of  $x$  and  $Q^2$
- neutron energy spectrum in  $\gamma p$  compatible with effects of absorption and migration;  
suppression in  $\gamma p$  at low  $x_L$ , high  $p_T$
- better agreement with data if account for absorption and for  $(\rho, a_2)$  exchanges
- suppression seen also in photoproduction of dijets

## Multiple interactions are very relevant for hadronic interactions

- description of average charged particle multiplicity in resolved photon events with hard jets requires MPI (at least when using PYTHIA);  
interestingly, the model without MPI, CASCADE, is not too bad