

## QFT II exercises - sheet 0

### Exercise 1

- a Show, assuming  $a \neq 0$  and  $J \neq 0$  that

$$\int_{-\infty}^{\infty} dx e^{-\frac{1}{2}ax^2 + Jx} = \sqrt{\frac{2\pi}{a}} e^{\frac{J^2}{2a}}$$

- b Show, assuming  $A \neq 0$  and  $J \neq 0$  and with summation convention as well as a suitable constraint on the matrix  $A$  that  $Z[J] :=$

$$\int_{-\infty}^{\infty} dx_1 \dots dx_n e^{-\frac{1}{2}(x_i A_{ij} x_j) + J_i x_i} = \sqrt{\frac{(2\pi)^n}{\det(A)}} e^{\frac{1}{2} J_i J_j (A^{-1})_{ij}}$$

*Hint:* Diagonalize  $A$ .

- c Compute

$$\frac{1}{Z[0]} \frac{\partial^4 Z[J]}{\partial J_i \partial J_j \partial J_k \partial J_l} \Big|_{J=0} = \int_{-\infty}^{\infty} dx_1 \dots dx_n x_i x_j x_k x_l e^{-\frac{1}{2}(x^T A x)}$$

- d Formulate a simple 'Feynman rule' which reproduce the result of c).

### Exercise 2

A one dimensional harmonic oscillator with external force  $J(t)$  is described by the Lagrangian

$$L = \frac{1}{2} \dot{x}^2 - \frac{1}{2} \omega^2 x^2 + Jx$$

- a Give the action in the form

$$S = \int dt \left[ \frac{1}{2} x A x + Jx \right]$$

and determine the differential operator  $A$ , assuming surface terms vanish

- b The Green's function  $G(t - t')$  of this operator satisfies

$$A(t)G(t - t') = \delta(t - t')$$

Show

$$S = \int dt \left[ \frac{1}{2} x' A x' - \frac{1}{2} \int dt' J(t) G(t - t') J(t') \right]$$

for a suitably shifted  $x' = x + X$ . *Hint:*  $X$  must be linear in  $J$  and involve an integral

- c Give the Fourier representation of  $G(t - t')$

### Exercise 3

The action for a gauge boson  $A_\mu$  is given by

$$S = -\frac{1}{4} \int d^4x F_{\mu\nu} F^{\mu\nu}, \quad F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

a By using partial integration write the action in the form

$$S = \frac{1}{2} \int d^4x A_\mu D^{\mu\nu} A_\nu$$

and determine the differential operator  $D^{\mu\nu}$

b By using the Ansatz  $A_\mu(x) = \int \frac{d^4k}{(2\pi)^4} \tilde{A}_\mu e^{-ik \cdot x}$  show that the action in Fourier space takes the form

$$S = \frac{1}{2} \int \frac{d^4k}{(2\pi)^4} \tilde{A}_\mu(k) \tilde{D}^{\mu\nu} \tilde{A}_\nu(-k)$$

and compute  $\tilde{D}^{\mu\nu}(k)$ .

c The Greens function  $G_{\nu\rho}(x-y)$  of  $D^{\mu\nu}$  is defined by

$$D^{\mu\nu} G_{\nu\rho}(x-y) = i\delta^\mu_\rho \delta(x-y) \quad (1)$$

Show that  $G_{\nu\rho}(x-y)$  is ill defined by acting with  $\partial_\mu$  on equation (1)

d By using the Ansatz  $G_{\nu\rho}(x-y) = \int \frac{d^4k}{(2\pi)^4} \tilde{G}_{\nu\rho} e^{ik(x-y)}$  determine the analog of equation (1) for  $\tilde{G}_{\nu\rho}$ . How does the problem of c) show up in this equation?

e Add to the action  $S$  a term

$$\delta S = -\frac{1}{2\xi} \int d^4x (\partial_\mu A^\mu)^2$$

and recompute  $D^{\mu\nu}$  and  $\tilde{D}^{\mu\nu}$ . Is the problem of c) still there?

f Determine  $\tilde{G}_{\nu\rho}$  with the help of the Ansatz

$$\tilde{G}_{\nu\rho} = a(k^2)\eta_{\nu\rho} + b(k^2)k_\nu k_\rho$$

and compute  $a$  and  $b$ .

g BONUS: instead of e), take

$$\delta S = -\frac{1}{2} \int d^4x (q_\mu A^\mu)^2$$

for some light-like vector  $q$  ( $q^2 = 0$ ). Determine  $\tilde{G}_{\nu\rho}$  with the help of the Ansatz

$$\tilde{G}_{\nu\rho} = a(k^2)\eta_{\nu\rho} + b(k^2)(q_\nu k_\rho + q_\rho k_\nu)$$