## Recent Slice Energy Spread Measurements in the EuXFEL

## Injector

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## Overview

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## Motivation

- Uncorrelated (slice) energy spread is an important property in FELs, where high brightness is required.
- However, usually energy spread from injector is too low and must be increased using the laser heater (LH) to increase SASE performance due to the microbunching instability.
- Ideally increased to 8 keV based on simulations from Martin, but actual optimal value unknown.
- Regardless, understanding uncorrelated energy spread is necessary to improve machine reproducibility and understanding.
- Dynamics not understood, unable to recreate measured slice energy spread spread values with ASTRA simulation.


## Context

- Fermi, 2020: Intrabeam scattering measured, but result highly dependent on initial energy spread, which was treated as a free parameter.
- SwissFEL, 2020: energy spread measurement in the injector using an energy scan—15keV @ 200pC \& $>100 \mathrm{MeV}$.
- EuXFEL, 2021: Similar approaching involving a dispersion scan—6keV @ 250pC \& 130MeV. PITZ 2022 — 2 keV @ 250pC \& 20MeV. 10.1103/PhysRevAccelBeams.25.083401
- SwissFEL, 2022: Contribution of the microbunching instability and intra-beam scattering evidenced by varying the optics -> reduced by adjusting optics and disabling LH chicane. 10.1103/PhysRevAccelBeams.25.104401
- EuXFEL, 2022, 2023: Now measured 4 keV , reduced to $\sim 3.4 \mathrm{keV}$ with some manipulation of the solenoid current. Unpublished.
- The energy spread has stayed stable since then, at about $\sim 4 \mathrm{keV}$.

| Scenario | $\sigma_{E} / \mathrm{keV}$ | Solenoid / A | Gun Gradient | Gun Phase |
| :--- | :---: | ---: | ---: | ---: |
| Feb. 2021 Published Result | $5.842 \pm 0.003$ | 338 | 56.7 | -42.9 |
| Nov. 2022 Measurement 1 | $4.313 \pm 0.004$ | 326.6 | 54.7 | -43.1 |
| Nov. 2022 Measurement 2 | $3.635 \pm 0.004$ | 336 | 56.5 | -43.6 |
| Nov. 2022 Measurement 3 | $3.385 \pm 0.004$ | 335 | 56.5 | -41.6 |

## The Usual Approach

Only provides an upper limit on uncorrelated spread $\sigma_{E}$.

Neglected contributions to size:


- Betatronic
- TDS-induced energy spread
- Imaging resolution

In the injector these contributions can be larger than the slice energy spread's contribution!

Need to separate these effects...


## Separating the contributions to the slice size

Assuming no correlation between TDS-induced energy spread and "true" energy spread, the final slice

$$
\sigma_{M}^{2}=\sigma_{R}^{2}+\sigma_{B}^{2}+\frac{D^{2}}{E^{2}} \sigma_{E}^{2}+\frac{D^{2}}{E^{2}}(e k V)^{2} \sigma_{I}^{2}
$$ energy spread seen at the screen:

$$
\begin{aligned}
\sigma_{E_{\mathrm{final}}}^{2}=\sigma_{E}^{2}+ & (e k V)^{2} \sigma_{I}^{2} \\
& \text { Beamsize in the TDS }
\end{aligned}
$$

$$
\sigma_{B}^{2}=\frac{\beta_{x} \varepsilon_{n}}{\gamma_{0}} \quad \sigma_{I}^{2}=\frac{\varepsilon_{n}}{\gamma_{0}}\left(\beta_{y}^{0}+0.25 L^{2} \gamma_{y}^{0}-L c\right.
$$

## Measurement Technique



## Derived Values

For the two scans:

$$
\sigma_{E}=\frac{E_{0}}{D_{0}} \sqrt{A_{V}-A_{D}}
$$

$\sigma_{M}^{2}=\sigma_{R}^{2}+\sigma_{B}^{2}+\frac{D^{2}}{E^{2}} \sigma_{E}^{2}+\frac{D^{2}}{E^{2}}(e k V)^{2} \sigma_{I}^{2}$

$$
\sigma_{I}=\frac{E_{0}}{D_{0} e k} \sqrt{B_{V}}
$$

$$
\begin{aligned}
\sigma_{M}^{2} & =A_{V}+B_{V} V^{2} \\
\sigma_{M}^{2} & =A_{D}+B_{D} D^{2}
\end{aligned}
$$

$$
\begin{aligned}
\sigma_{B} & =\sqrt{B_{\beta} \beta_{x}^{0}} \\
B_{\beta} & =\sigma_{I}^{2}\left(\beta_{y}^{0}+0\right. \\
\sigma_{R} & =\sqrt{A_{D}-\sigma_{B}^{2}}
\end{aligned}
$$

$$
B_{\beta}=\sigma_{I}^{2}\left(\beta_{y}^{0}+0.25 L^{2} \gamma_{y}^{0}-L \alpha_{y}^{0}\right)
$$

## Injector Setup

- Calibrate the gun phase.
- Turn off AH1 for minimum chirp contribution to energy spread.
- Go on crest in A1, adjust the voltage so we are at 130 MeV at the screen.
- Apply special optics to maximise the ratio of the dispersion to the betatron contributions to the spot size.
- Match the central slice using Matthias tool (vital for this measurement)
- Measure the dispersion.
- Turn the laser heater off.



## Special optics for measurement from TDS to the screen



## Procedure

TDS No. $=13, \eta_{x}=1.181 \mathrm{~m}$, image 1 , before/after image processing

Take 5 background images at the start of the measurement, and then 30 images:

1. At each TDS $V$ in the voltage scan.
2. At each $D_{x}$ at the screen in dispersion scan.
3. At each $\beta_{x}$ at the screen in the beta scan.

Then:

1. Subtract background.
2. Mask to remove isolated blobs.
3. Pick the largest connected non-zero pixel blob to be "the beam".
4. Fit a Gaussian to 10 slices centred on



 highest energy slice \& average.

## New Tooling

## Measurement

- Taking $\sim 400$ images across $\sim 15$ machine setpoints without special tooling would take about an hour. Then, offline analysis takes another 10 minutes.
- To this end we have developed a GUI to do the parameter scanning.
- Live image processing panel.
- Live slice width scan plots (e.g. TDS voltage against central slice width).
- Immediate results at the end of the scan.
- Bottom line: results in a few minutes.



## New Tooling

TDS Calibration

- Additionally the TDS needs to be calibrated for each measurement campaign to get the voltage $\rightarrow$ also takes up to an hour.
- Also developed an injector TDS calibration GUI
- Automatically scan TDS "amplitudes" and phases to build mapping of amplitudes [\%] to voltages [MV] $\rightarrow$ Gives all downstream longitudinal calibrations!
- Also useful for non-invasive emittance and bunch length measurements.
- Bottom lline: calibrate the injector TDS in a few minutes with no human intervention.

B2 calibrator development ongoing. Non-invasive development ongoing.

Extracted zero-crossing gradients at each amplitude

## Experiment details

- We did two scans over a single four hour shift.
- Charge scan: $50 \mathrm{pC}, 150 \mathrm{pC}, 250 \mathrm{pC}$ and 350 pC
- $R_{56}$ scan: $-1 \mathrm{~mm},-2.5 \mathrm{~mm},-4.336 \mathrm{~mm}$ (nominal), -5.5 mm and -7 mm .
- We rematched the projected emittance every single time,
- then measured the central slice emittances (showing good matching already).
- We additionally scanned the beta function at the screen (not shown here).
- We used the new online data data taking and analysis software for these scans, as well as for calibrating the TDS.


## Charge Scan

- We scanned from 50 pC to 350 pC in steps of 100 pC .
- Unable to resolve the TDS contribution to the energy spread at 50 pC , so assumed to be zero.
- No obvious reason to consider energy spread from TDS to be zero, related to the camera sensitivity?



## $R_{56}$ Scan

- We scanned the $\mathrm{R}_{56}$ of the laser heater chicane from • Design setting ( -4.336 mm ) not scanned in order -7 mm to -1 mm (the minimum). with the other setpoints.
- We rematched the beam each time, first projected, - Strong dependence on LH $R_{56}$ is apparent, from then the central slice. 4.5 keV at maximum to 2.4 keV at minimum.



## Measurement Details

## Possible Dispersion Leakage from the LH Chicane

- We did not remeasure the dispersion at each setpoint of the laser heater $\mathrm{R}_{56} \mathrm{scan}$.
- Possible: we have some leaked dispersion from the chicane at the screen.
- First two dipoles have same power supplies, second two are each independent.
- Consider: Identical 15\% error in first two dipoles and independent 15\% errors in second pair, perhaps some sort of "worst case".
- Simulated in OCELOT.
- Dispersion at the screen is tolerant to even large changes in the LH dipoles.
- Nevertheless should measure the dispersion in future experiments.
- Translates max $\pm 10 \%$ error in energy spread.



## Emittance Measurements

- Matthias matching/measurement tool is vital to our measurement.
- We match projected and then measure slice—and normally this gives reasonable matching, but we extract much smaller emittances from our scans.
- We see measured central slice emittances that are larger than the corresponding projected emittances even with apparently "good" matching.
- Is there any way we can reduce the error bars here?

ref


## Constant Contributions to Measured Widths

- In principle both are known:
- Screen resolution $\left(\sigma_{R}\right)$ from dedicated optics simulations (Artem)

$$
A_{D}=\sigma_{R}^{2}+\sigma_{B}^{2}
$$

- Betatron contribution $\left(\sigma_{B}\right)$ from slice emittance measurement and linear optics (Matthias)
- $\sigma_{R}=28 \mu \mathrm{~m}$ (from Artem and previous scans); $\beta_{x}=0.6 \mathrm{~m}$ and $\varepsilon_{x}=0.43 \mathrm{~mm} \cdot \mathrm{mrad}$ (Matthias matching tool)
- Constant terms remain constant across the scan, but $\sim 7 \%$ larger than expected (from Matthias+Artem).



The Matthias Tool

## Deriving the emittance assuming the screen resolution

- Fitting $A_{D}$ and using a pre-calculated value of $\sigma_{R}$, we can extract a value for the slice emittance.
- We independently measure, using Matthias' matching tool, $(0.43 \pm 0.05) \mathrm{mm} \cdot \mathrm{mrad}$ for the slice emittance.

$$
\varepsilon_{n}=\frac{\gamma_{0}}{\beta_{x}}\left(A_{D}-\sigma_{R}^{2}\right)
$$

- Not entirely in disagreement with our derived value.
- Can we reduce the uncertainties in the emittance measurement?



## Conclusion

- The slice energy spread shows wide tunability based on the charge and laser heater chicane $\mathrm{R}_{56}$.
- (2.4-4.5) keV from -1mm to -7mm $\mathrm{R}_{56}$ @ 250 pC
- (3.0-4.5) keV from 50pC to 350pC @ nominal $\mathrm{R}_{56}$.
- In combination with the longer term decrease from $\sim 6 \mathrm{keV}$ to 4 keV in the last two years for unknown reasons, it is clear that there is still a lot to be understood about the LPS dynamics in the injector.
- Important to understand to boost reproducibility.
- The charge scan points at possible IBS contributions to the slice energy spread.
- The $\mathrm{R}_{56}$ dependence needs more understanding and explanation. MBI?
- We now have a fast and easy to use data taking and online analysis tool which gives fast results. Room for further studies of other parameters, e.g. bunch length.


## Outlook

- Upcoming blue week in October for additional chicane studies.
- Aiming to take part in the upcoming Nepal laser commissioning--bunch length scans (not possible with laser 2).
- We would still like to provide an absolute calibration of the laser heater.
- IBS simulation work ongoing with Erion to ground these results more in theory.
- Further work on the GUI and tooling to improve usability and speed.
- Broader GUI development will feed back into parasitic bunch length and slice emittance measurements for use in normal operation.


## Final Slide

- Thanks also to:
- Frank (beam energy coming out of the gun measurement)
- Torsten (useful and interesting conversations regarding laser heater)
- Weilun (night shift moral support)
- Questions?


## Thank you

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