

α_s extraction from radiative quarkonium decays

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(work done with Nora Brambilla, Joan Soto and Antonio Vairo)

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Outline of the talk

- Reminder: Current α_s from PDG

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- The photon spectrum ($\Upsilon \rightarrow X\gamma$)

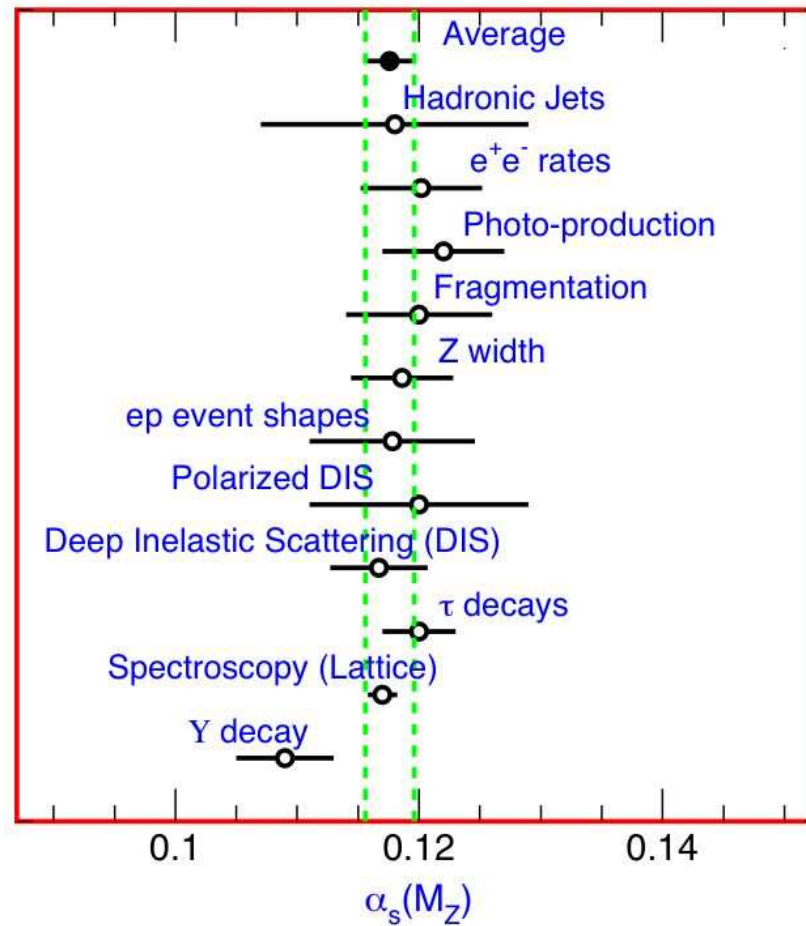
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- The extraction
- Conclusions

α_s from PDG06



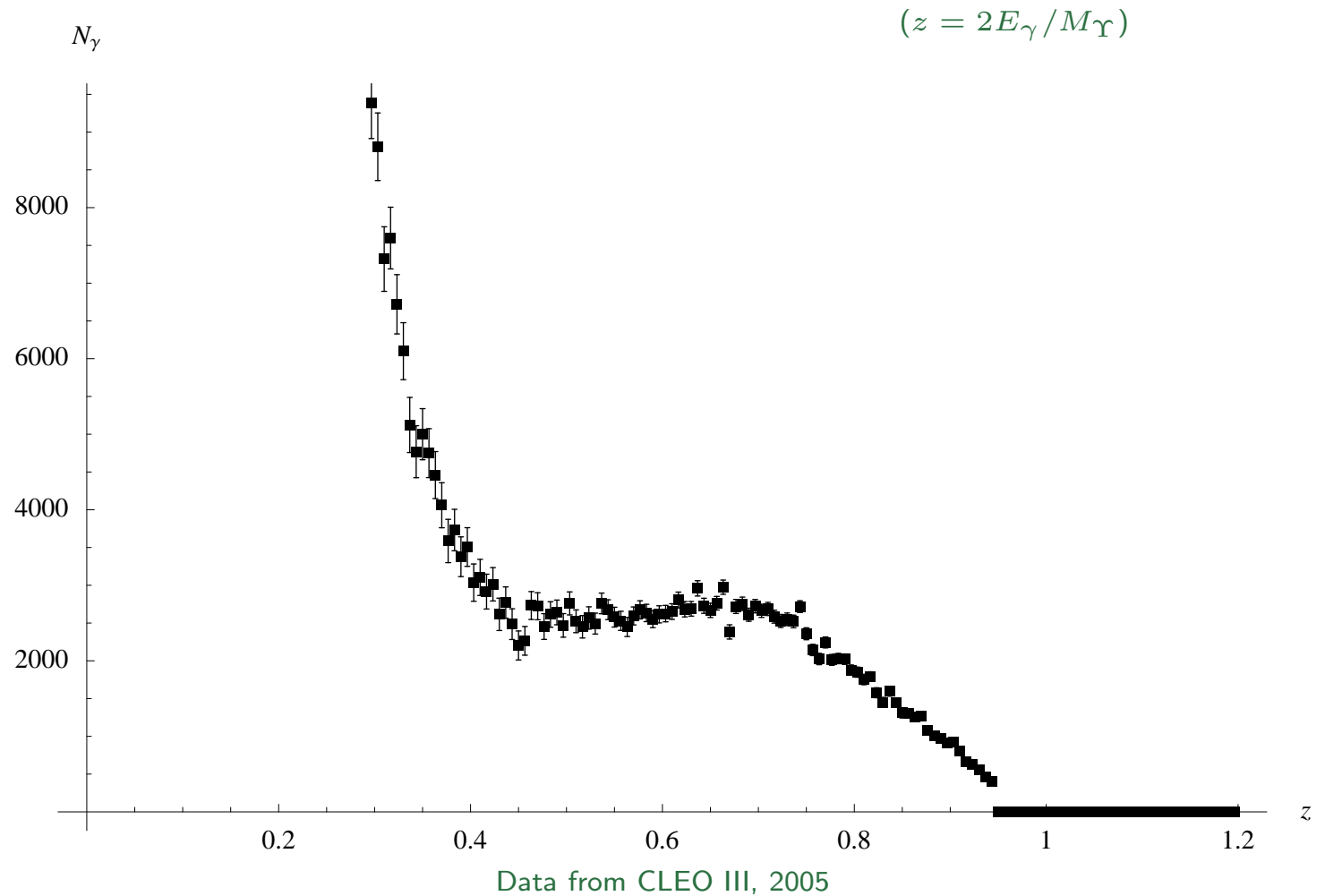
From W.-M. Yao et al., J. Phys. G 33, 1 (2006)

The photon spectrum ($\Upsilon \rightarrow X\gamma$)

- Recent CLEO measurement of the photon spectrum

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Recent CLEO measurement of the photon spectrum



[D. Besson *et al.* [CLEO Collaboration], Phys. Rev. D **74** (2006) 012003 (hep-ex/0512061)]

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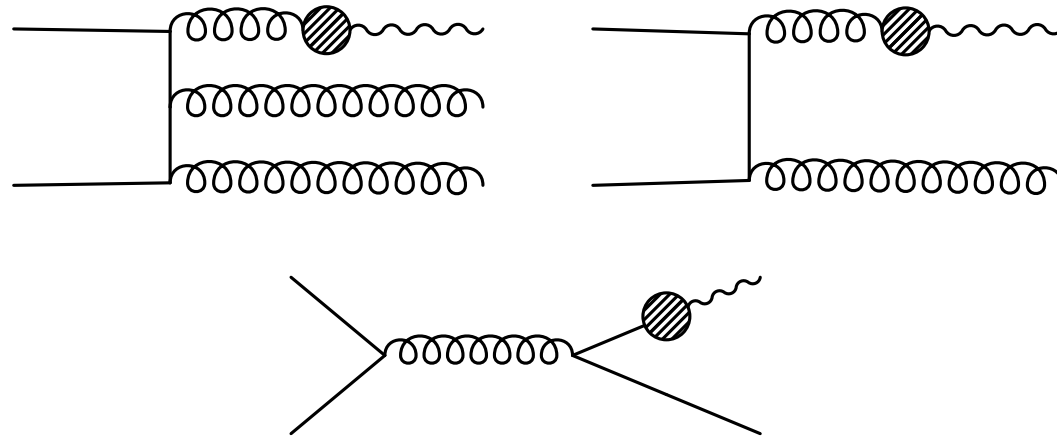
The photon spectrum ($\Upsilon \rightarrow X\gamma$)

- Recent CLEO measurement of the photon spectrum
- Two types of contributions to the photon spectrum
 - ◆ Fragmentation contributions
 - ◆ Direct contributions

Fragmentation contributions

Electromagnetic couplings to light quarks

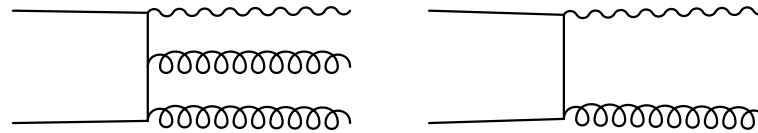
$$\frac{d\Gamma^{frag}}{dz} = \sum_a C_a \otimes D_{a \rightarrow \gamma}$$



This type of contributions become important in the low z region of the spectrum

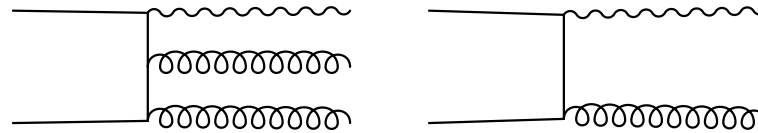
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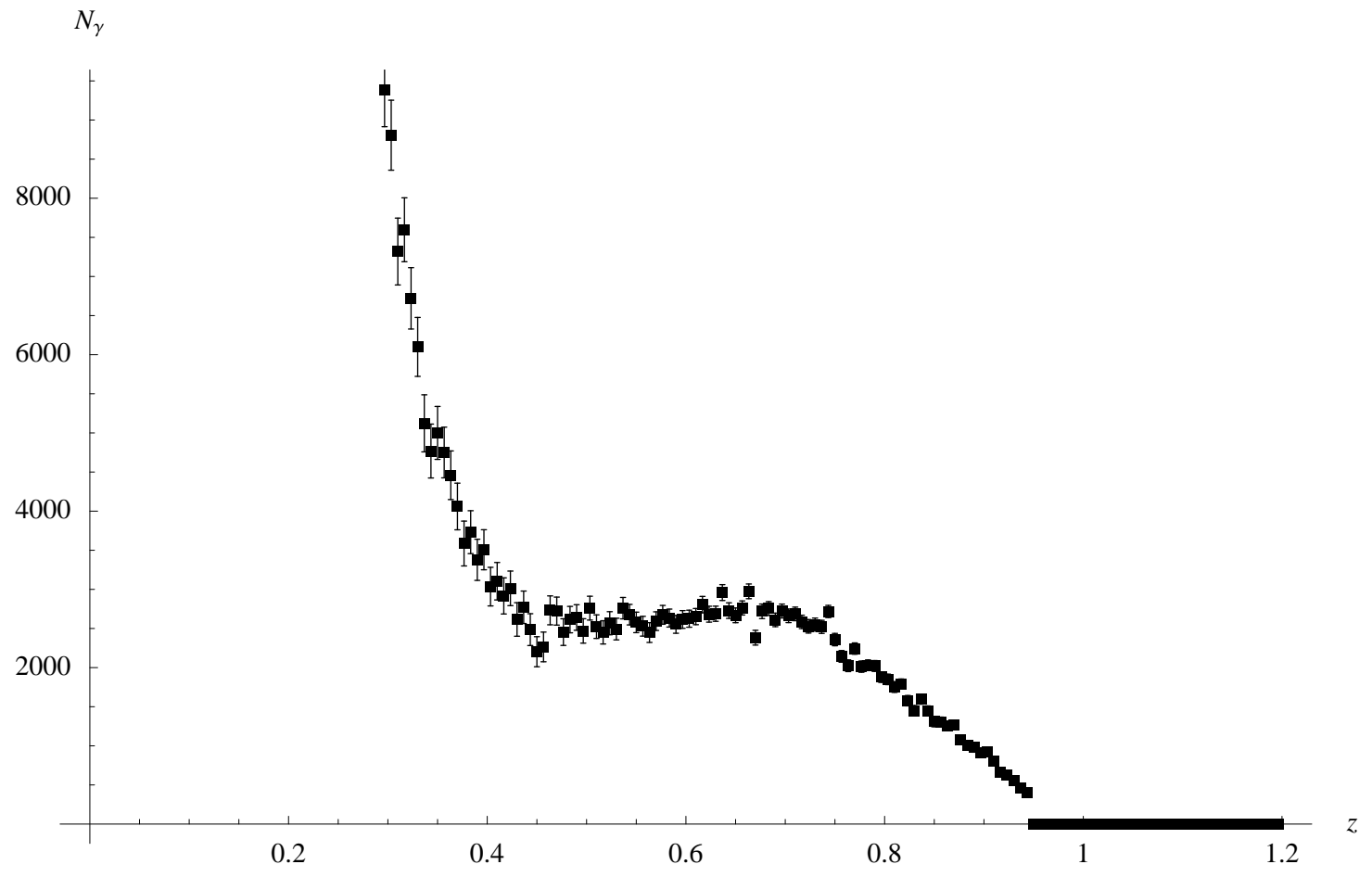
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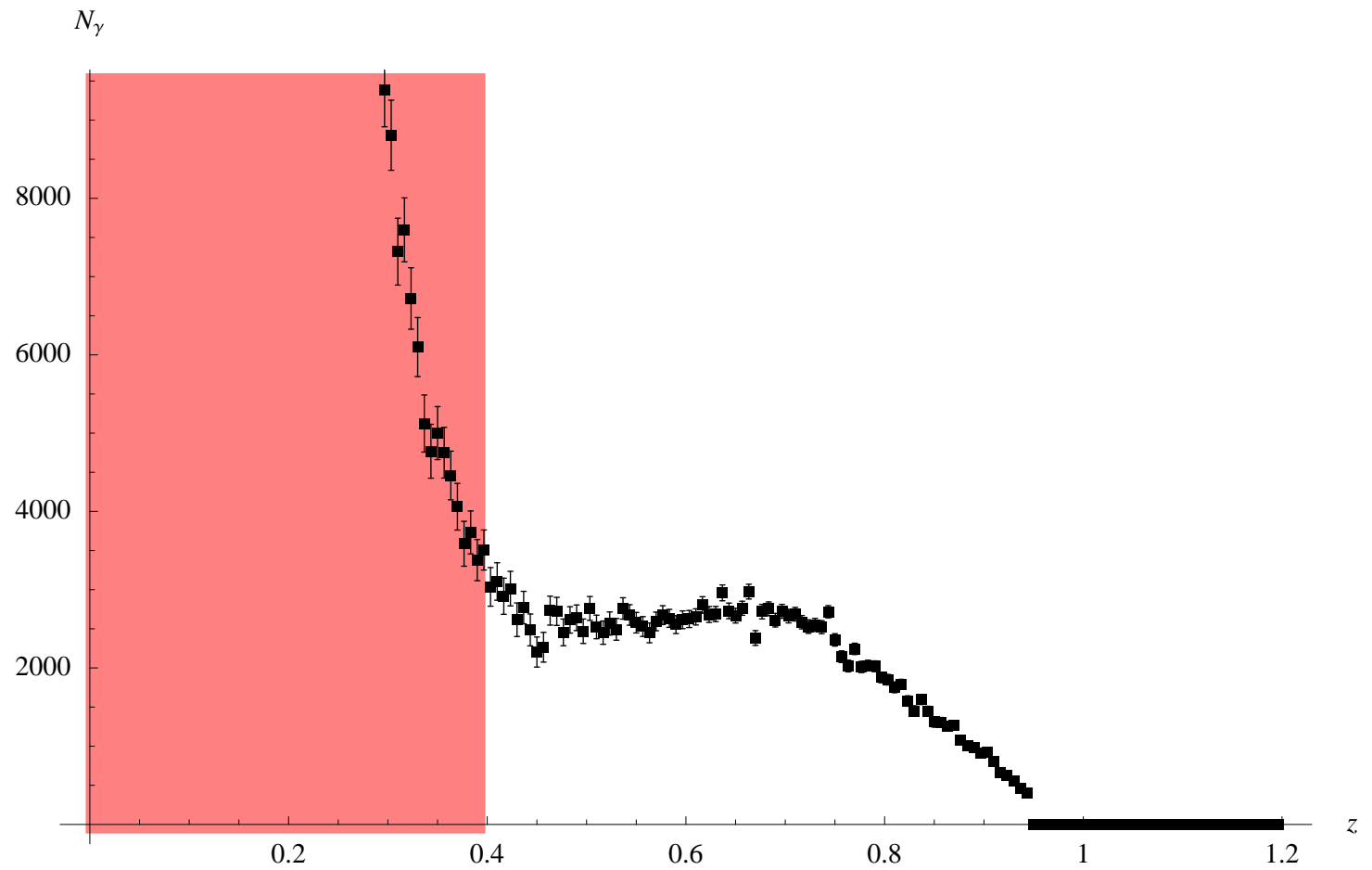
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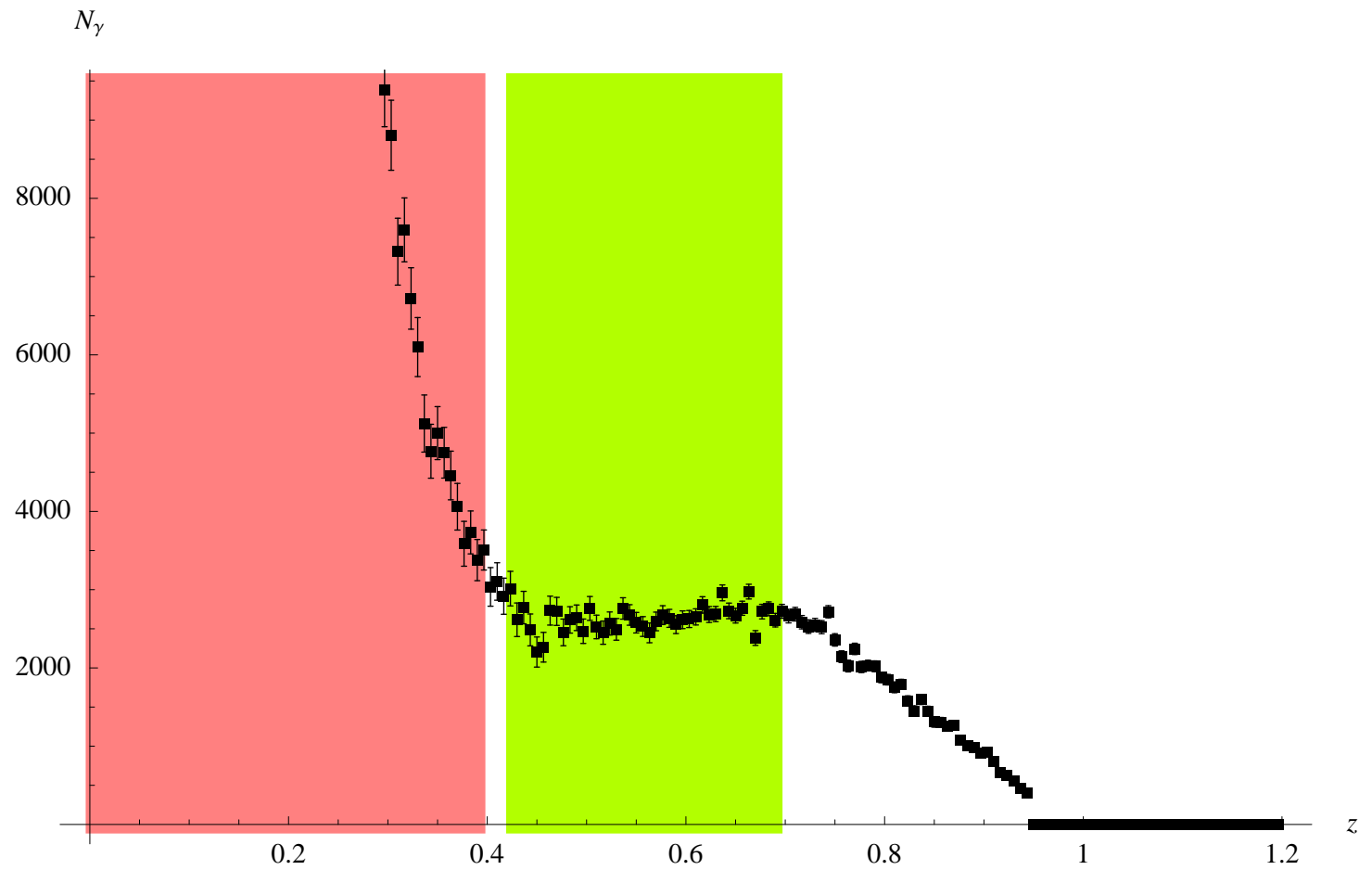
Valid in the central z region

- For $z \rightarrow 0$ the photon can only cause transitions within the bound state. But this is the fragmentation dominated region
- For $z \rightarrow 1$ NRQCD expansion breaks down. Collinear degrees of freedom become relevant
 - ◆ Large $\log(1 - z)$ need to be resummed
Photiadis '85; Bauer et al. '01; Fleming and Leibovich '02 '04
 - ◆ Shape functions must be introduced. Rothstein and Wise '97
Can be calculated assuming Coulombic state X.G.T. and Soto '04

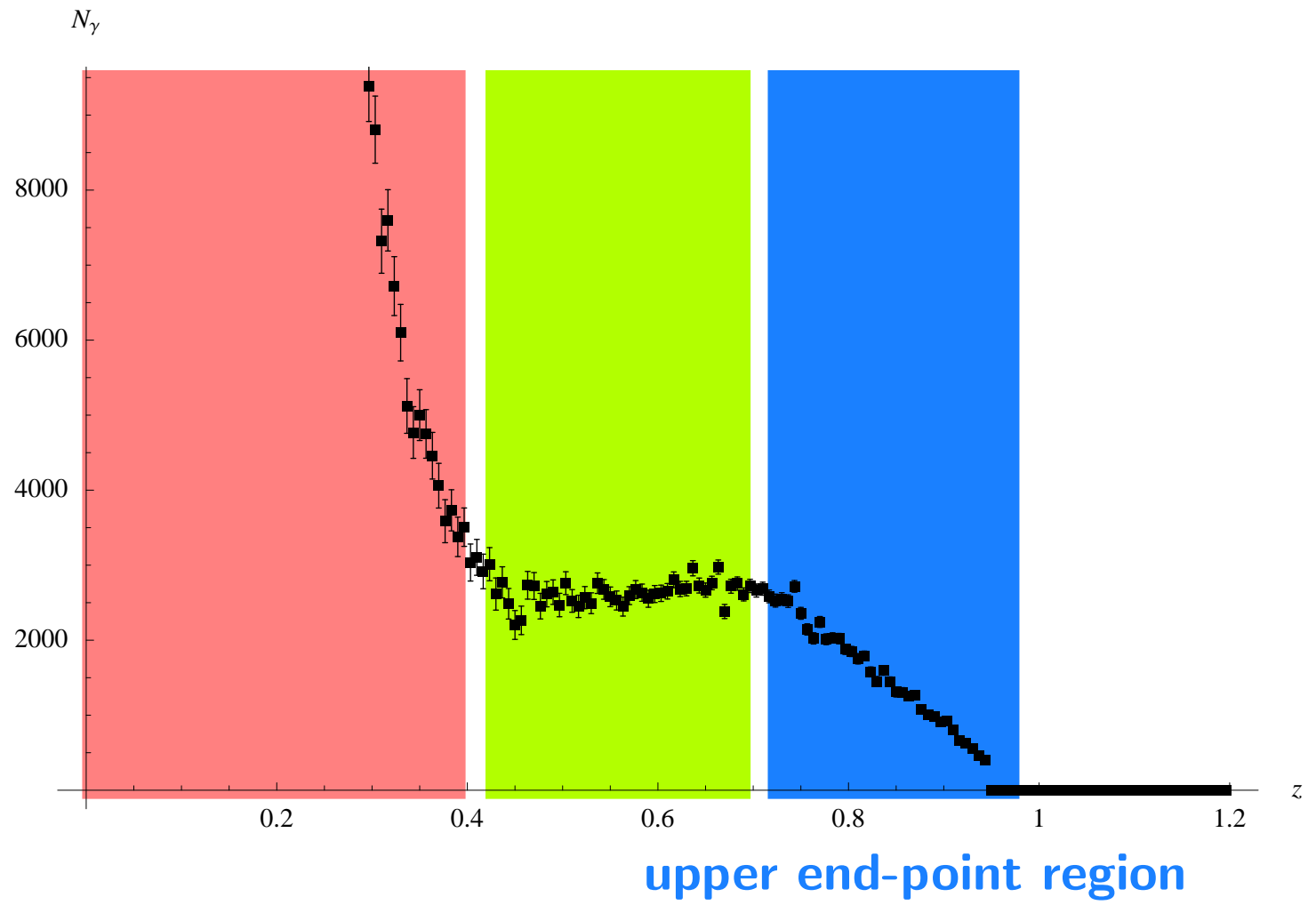




Fragmentation region



NRQCD (central) region

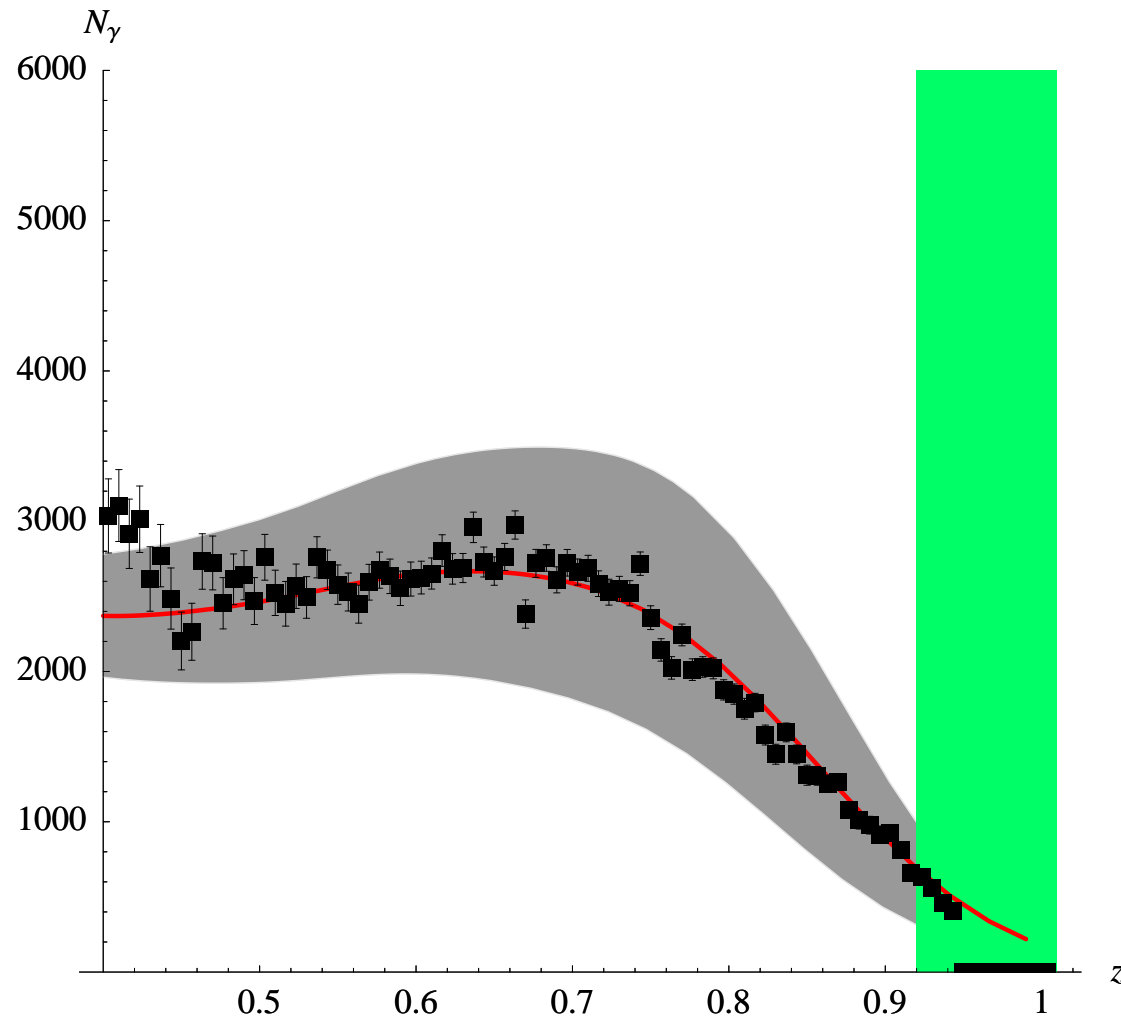


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- Compare then with a theoretical expression for R_γ to extract α_s

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We will use the same counting employed for the calculation of the spectrum to extract α_s

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- ◆ We include all the terms up to $\mathcal{O}(v^2)$: radiative corrections, relativistic corrections and octet operators
- Inclusion of all those terms well known to be potentially important
- Data is now very precise and we have all the necessary theoretical ingredients to include all the pieces

Schematically

$$\frac{\Gamma_{gg\gamma}}{\Gamma_{ggg}} = \frac{C_{\gamma O_1(^3S_1)} O_1(^3S_1) + C_{\mathcal{P}_1(^3S_1)} \mathcal{P}_1(^3S_1) + C_{\gamma O_8(^1S_0, ^3P_0)} O_8(^1S_0, ^3P_0)}{C_{O_1(^3S_1)} O_1(^3S_1) + C_{\mathcal{P}_1(^3S_1)} \mathcal{P}_1(^3S_1) + C_{O_8(^1S_0, ^3P_0, ^3S_1)} O_8(^1S_0, ^3P_0, ^3S_1)}$$

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Octet matrix elements:

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Octet matrix elements:

- $O_8(^1S_0)$ and $O_8(^3P_0)$ have been estimated in the continuum (weak coupling)

X.G.T. and Soto '04

- $O_8(^3S_1)$ and $O_8(^1S_0)$ have been calculated on the lattice

Bodwin, Lee and Sinclair '05

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 - ◆ Uses all the weak-coupling expressions available and lattice calculation for $\mathcal{O}_8(^3S_1)$
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The two procedures give very similar results. We take the average as the final value

Error estimation

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$$\left(\mathcal{R}_O = \frac{\langle O \rangle}{m \Delta d \langle O_1(^3S_1) \rangle} \right)$$

■ C (for continuum)

$$0.18 \leq \alpha_S(m_b v) \leq 0.38$$

$$0.32 \leq \alpha_S(m_b v^2) \leq 1.3$$

$$0 \leq \mathcal{R}_{O_8(^3S_1)} \leq 1.6 \times 10^{-4}$$

■ L (for lattice)

$$0 \leq \mathcal{R}_{O_8(^1S_0)} \leq 4.8 \times 10^{-3}$$

$$0 \leq \mathcal{R}_{O_8(^3S_1)} \leq 1.6 \times 10^{-4}$$

$$-2.4 \times 10^{-4} \leq \mathcal{R}_{O_8(^3P_0)} \leq 2.4 \times 10^{-4}$$

$$-0.052 \leq \mathcal{R}_{P_1(^3S_1)} \leq -0.035$$

Plus errors associated to higher order terms (v^3) and experimental errors

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N.Brambilla, X.G.T., J.Soto, A.Vairo '07

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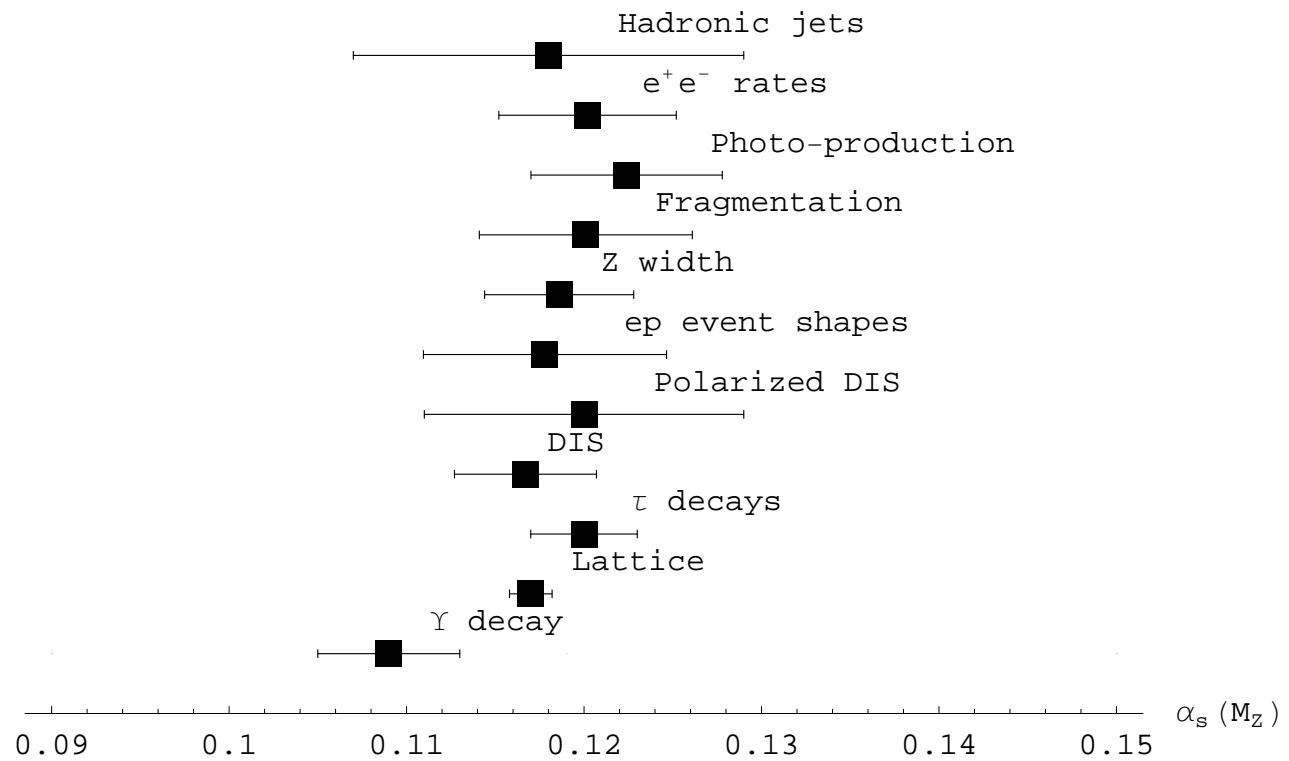
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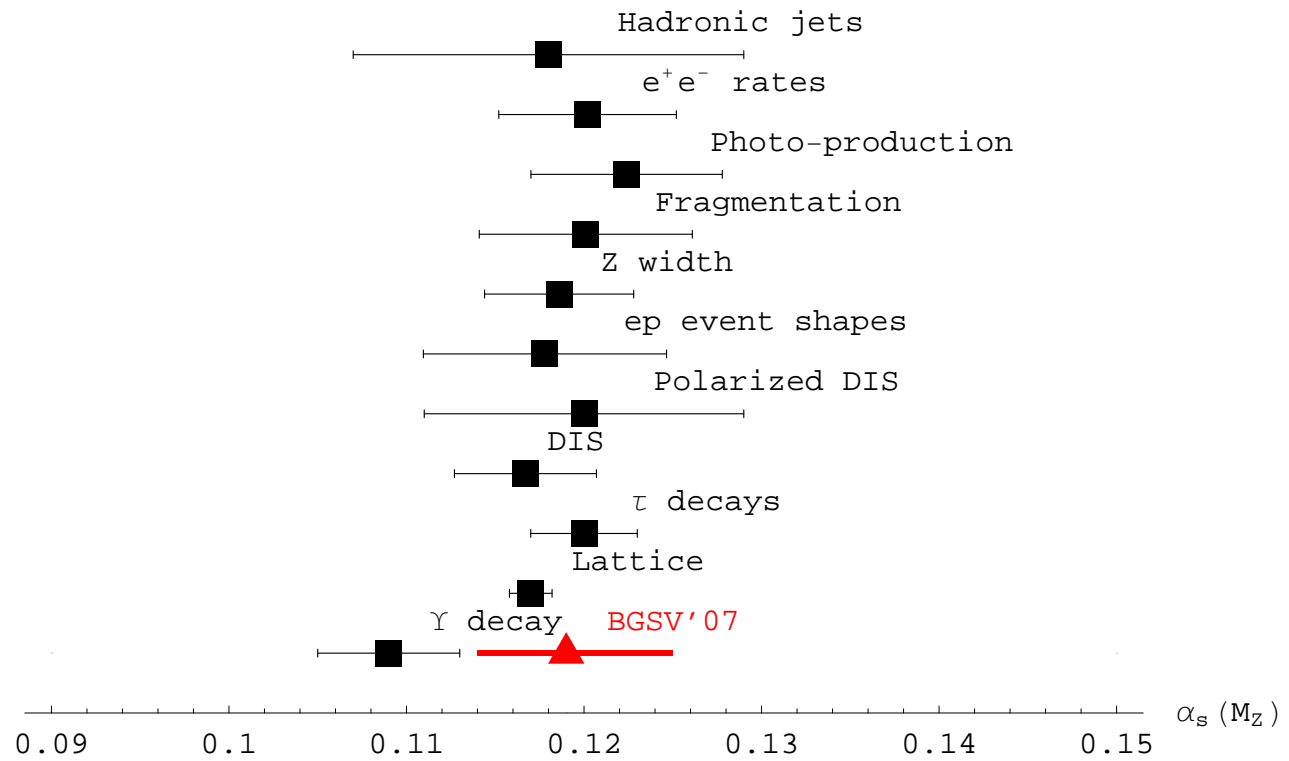
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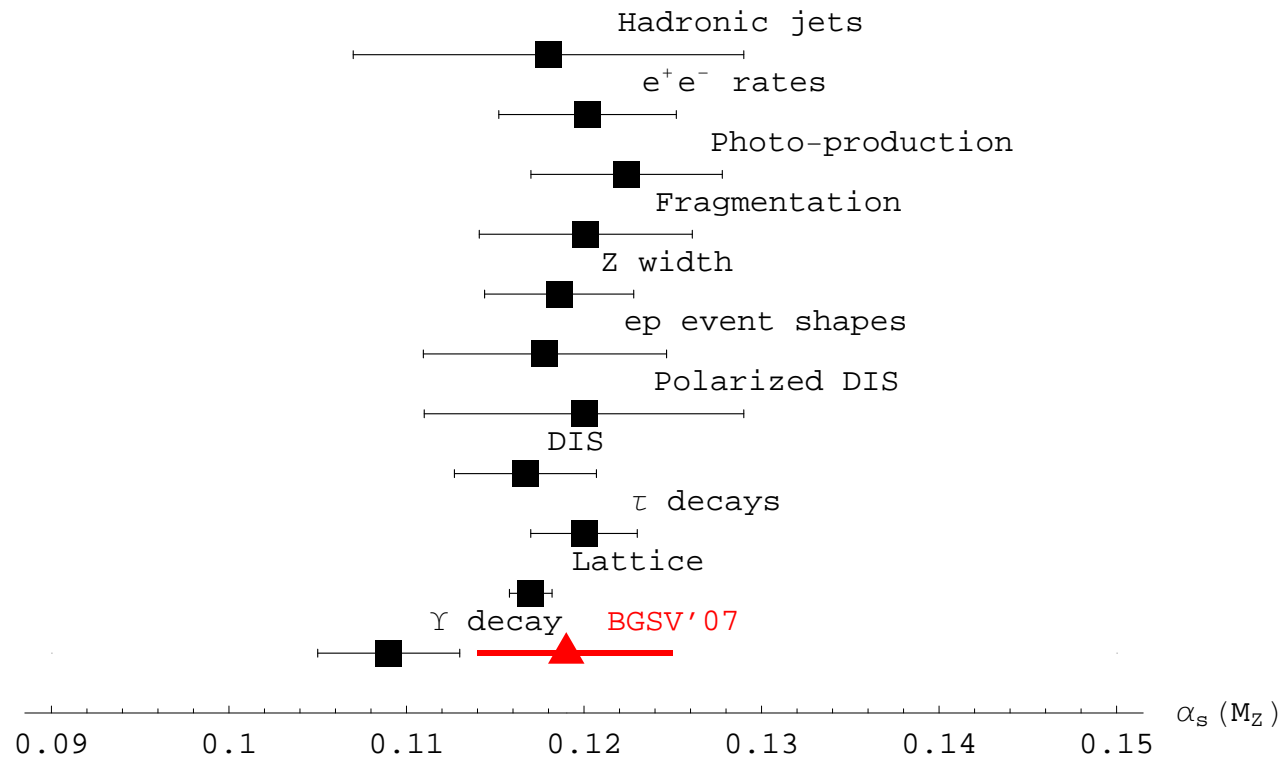
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Recall, PDG average $\rightarrow \alpha_s(M_Z) = 0.1176 \pm 0.0020$







Of course, if the new value is used in the average, the average will move up

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- The new extracted value of α_s shows better agreement with the other determinations