

Physics at HERA

Summer Student Lectures
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Overview

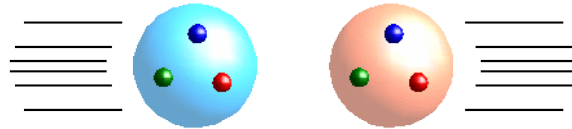
- Introduction to HERA
- Inclusive DIS & Structure Functions
 - formalism
 - HERA results
- High Q^2 & Electroweak Physics
- QCD: Jet Physics, Heavy Flavour Production
- Beyond the Standard Model
- (Diffraction)

Collider Types



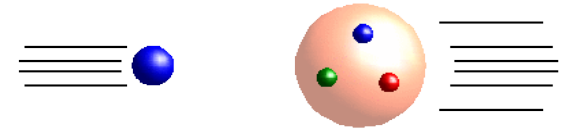
e^+e^-

- + clean initial and final state
- + small background
- limited energy
- LEP (200 GeV)
ILC (1 TeV)



$p^\pm p^\pm$

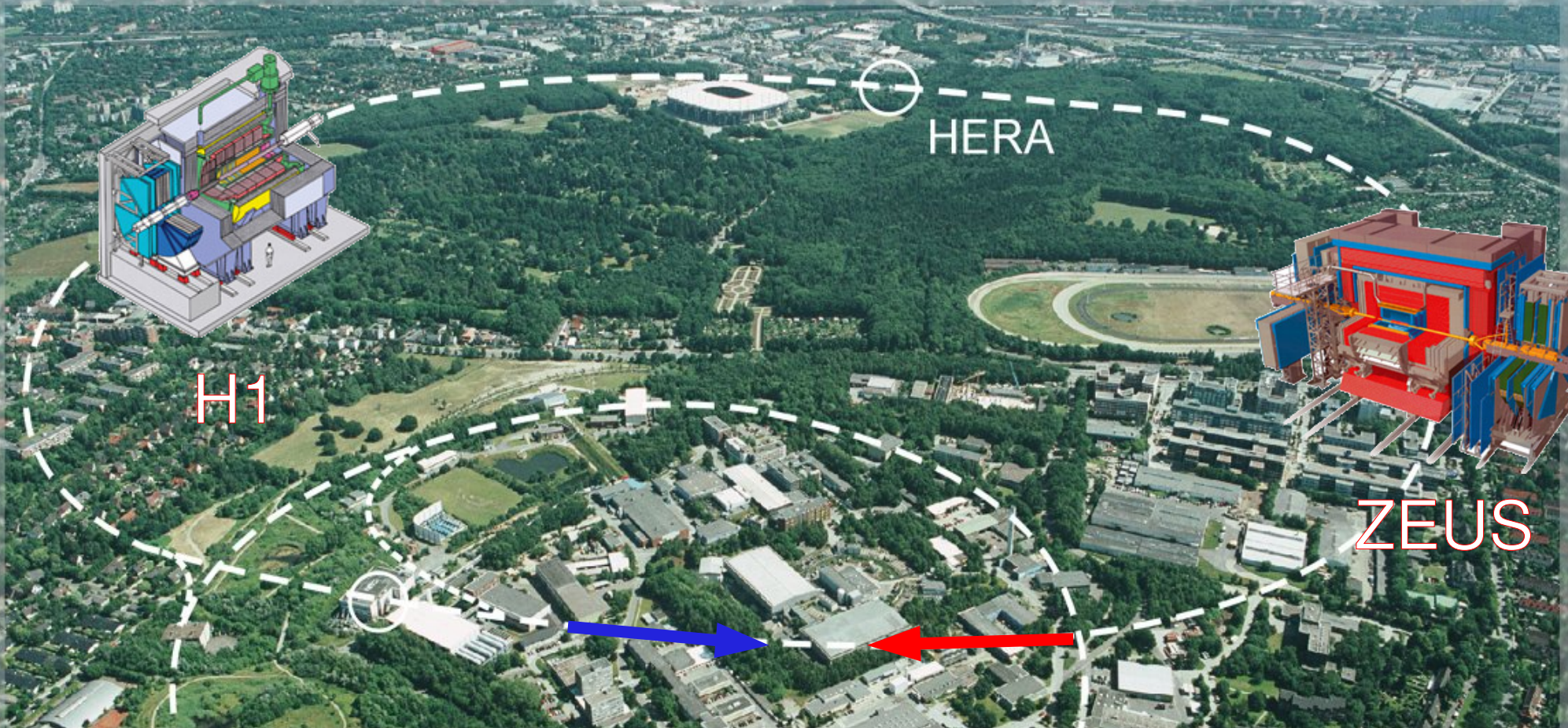
- + high energy
- complicated final state
- large background
- Tevatron (2 TeV)
LHC (14 TeV)



ep

- + unique initial state
- + electron as probe of proton structure
- two accelerators
- HERA (300 GeV)

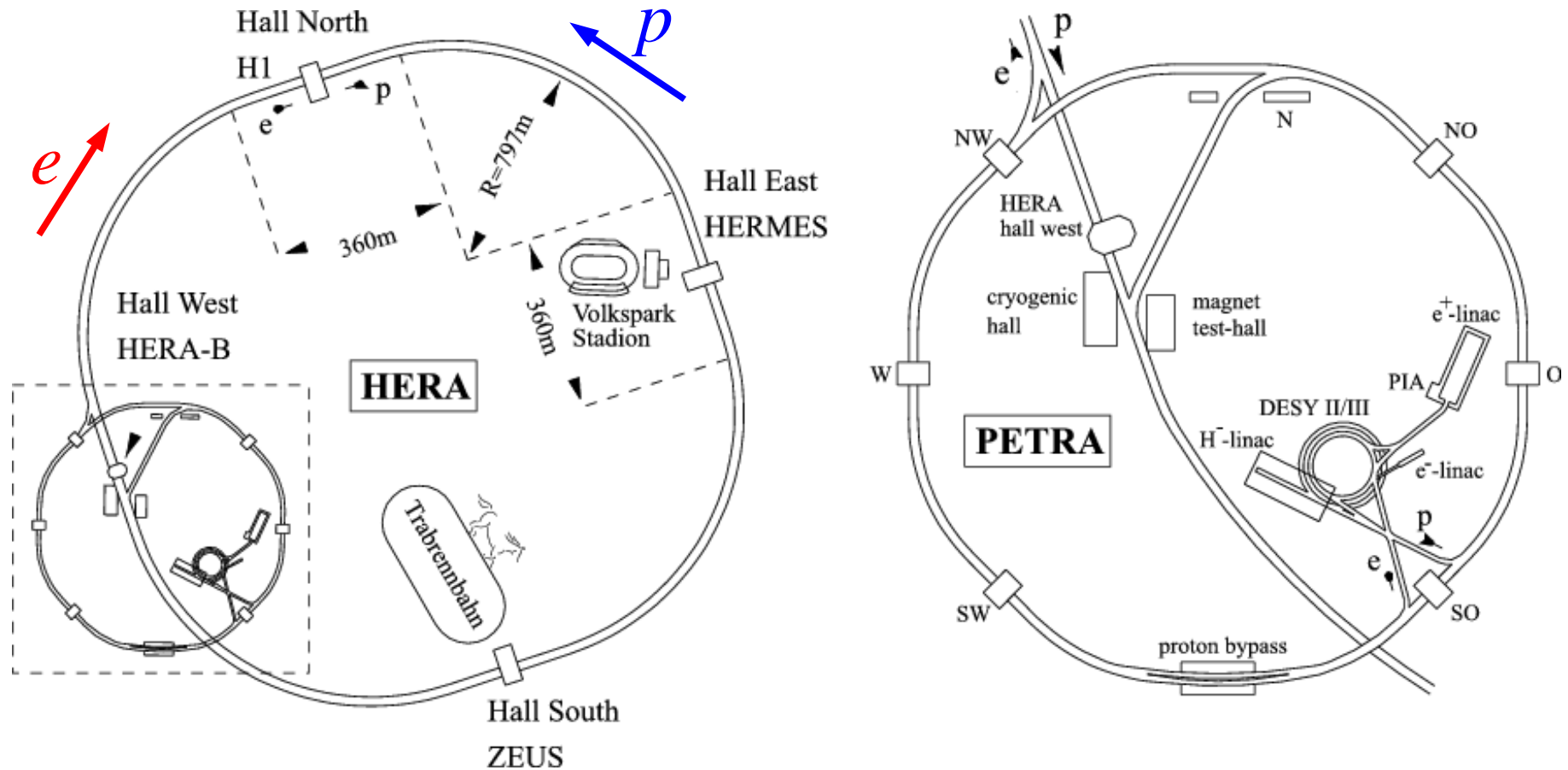
HERA



p
920 GeV

e
27.6 GeV

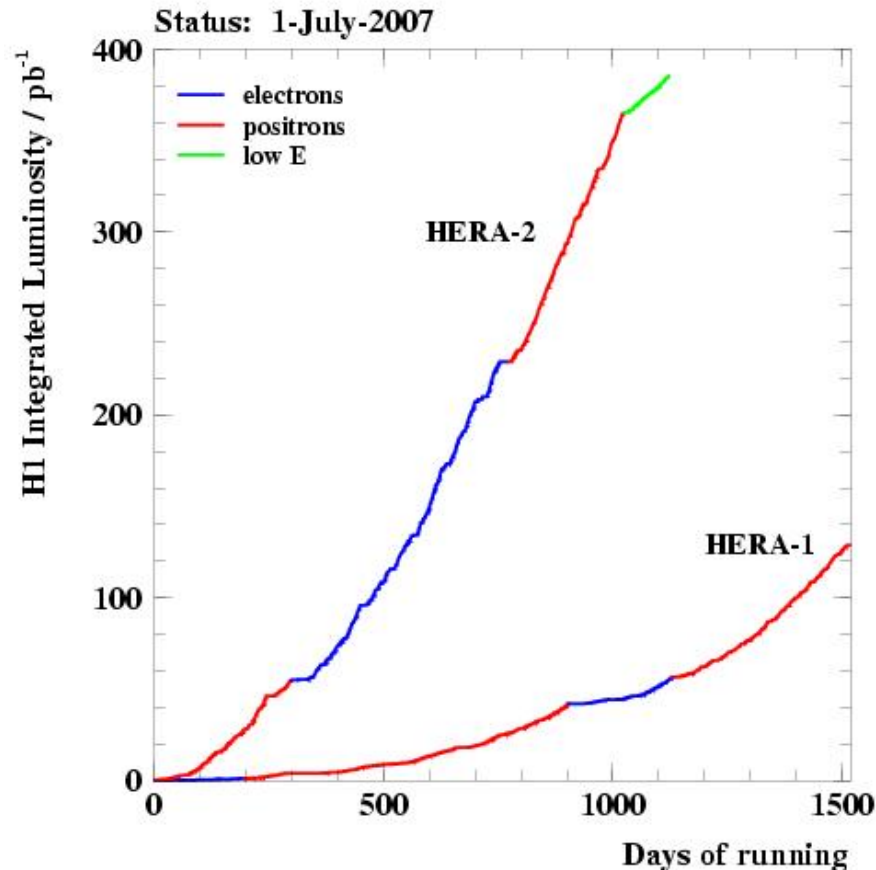
HERA & its Pre-Accelerators



circumference: 6.3 km

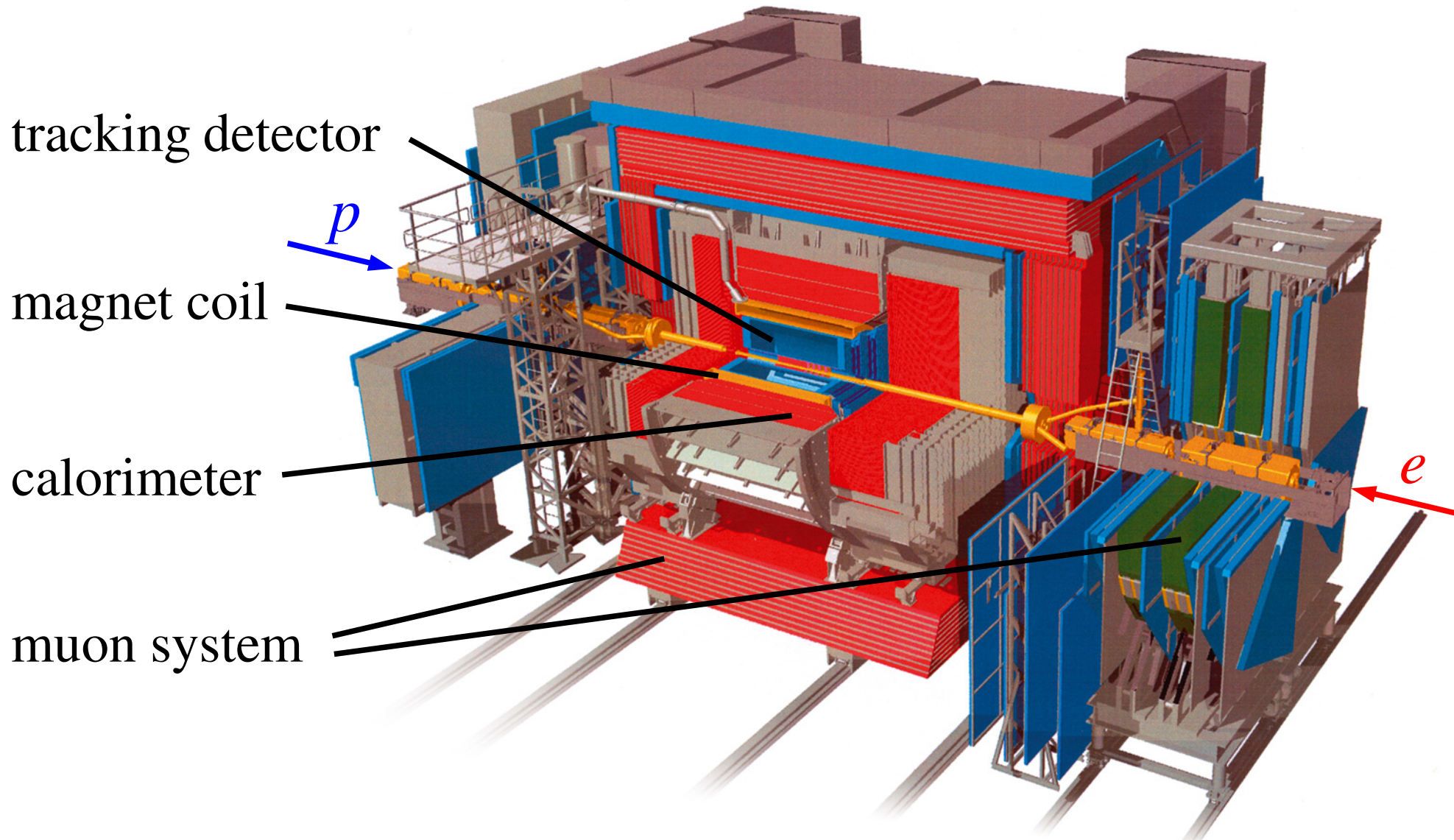
bunch crossing rate: 10.4 MHz

Collected Luminosity

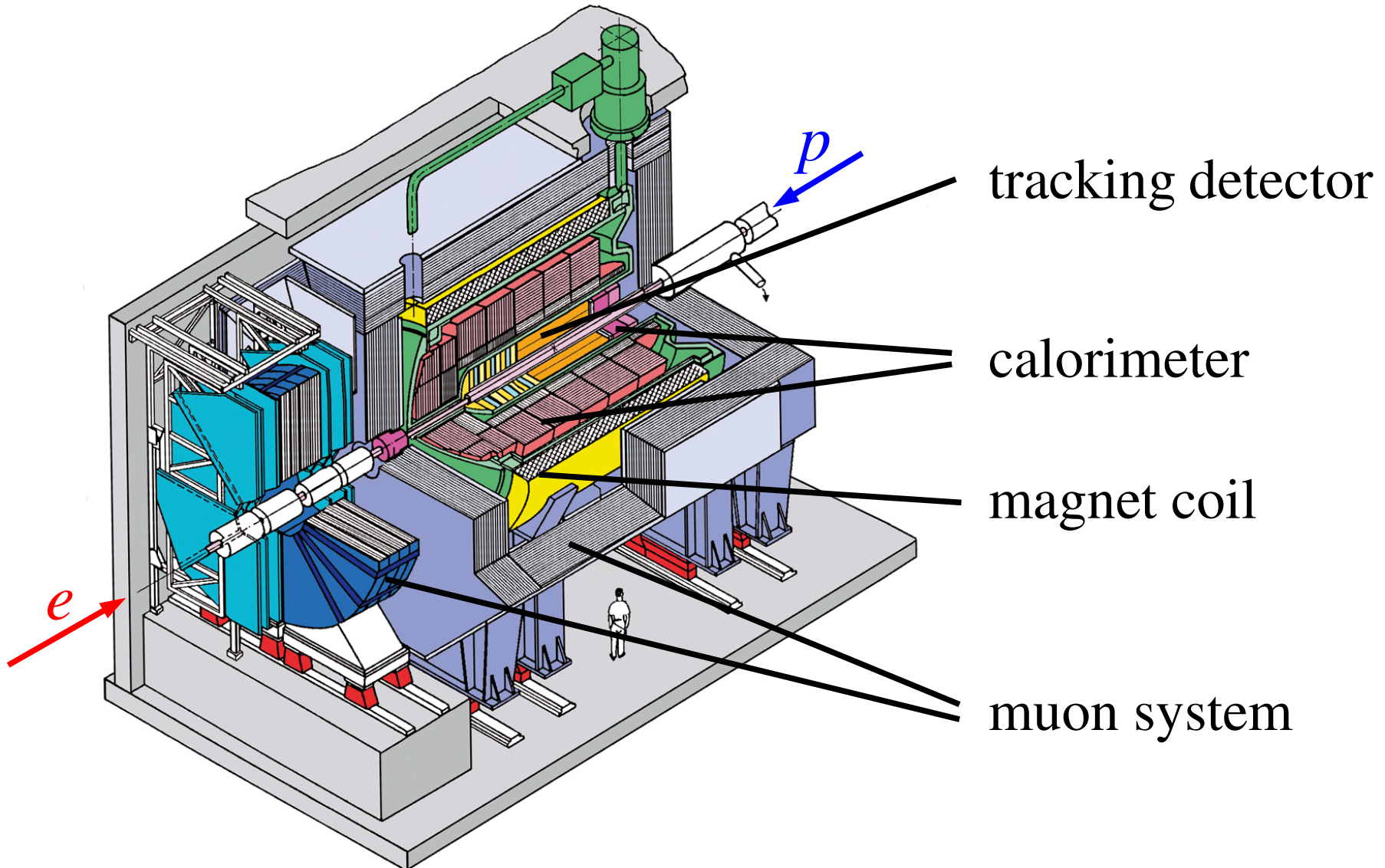


- HERA operated 1992-2007
- lumi upgrade in 2001
 - higher luminosity
 - e polarization
 - detector upgrades
- in total $\sim 500 \text{ pb}^{-1}$ of high energy data collected per experiment
- last months devoted to low p energy (460, 575 GeV)

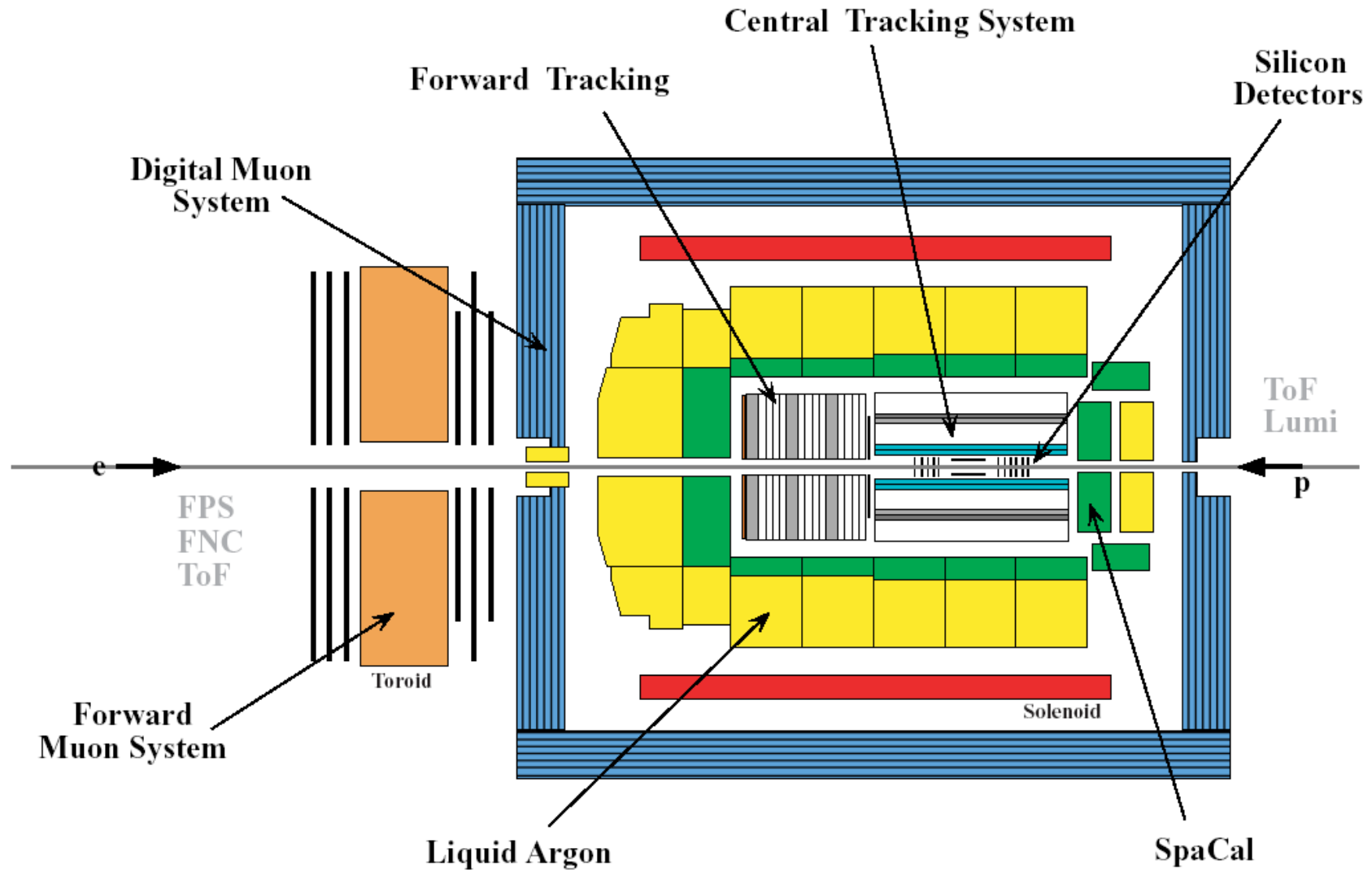
ZEUS Detector



H1 Detector



Schematic View of the H1 Detector



Physics Topics at HERA

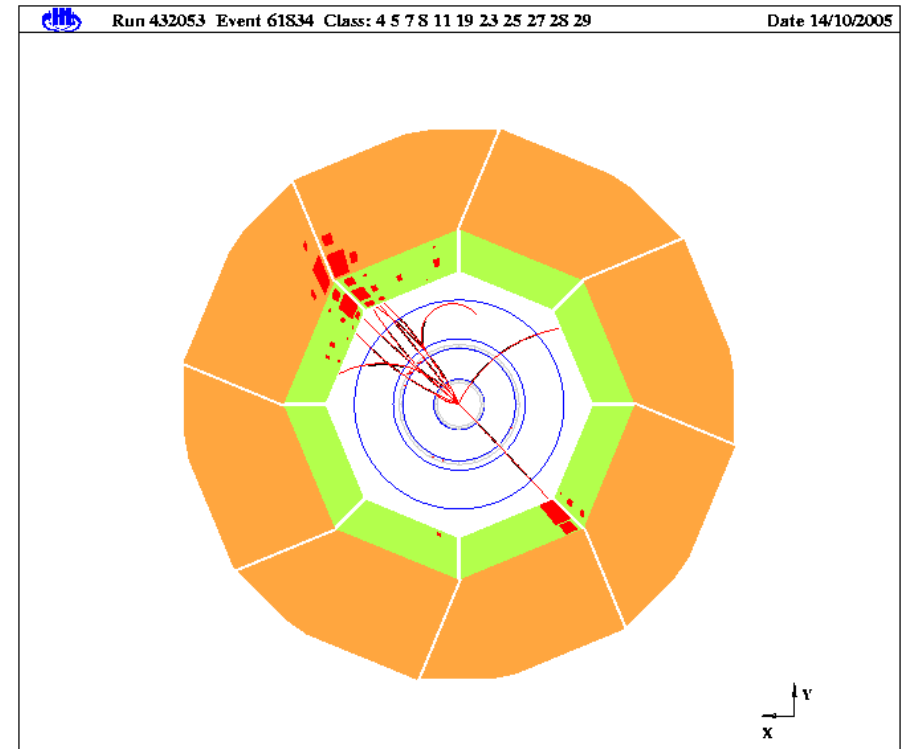
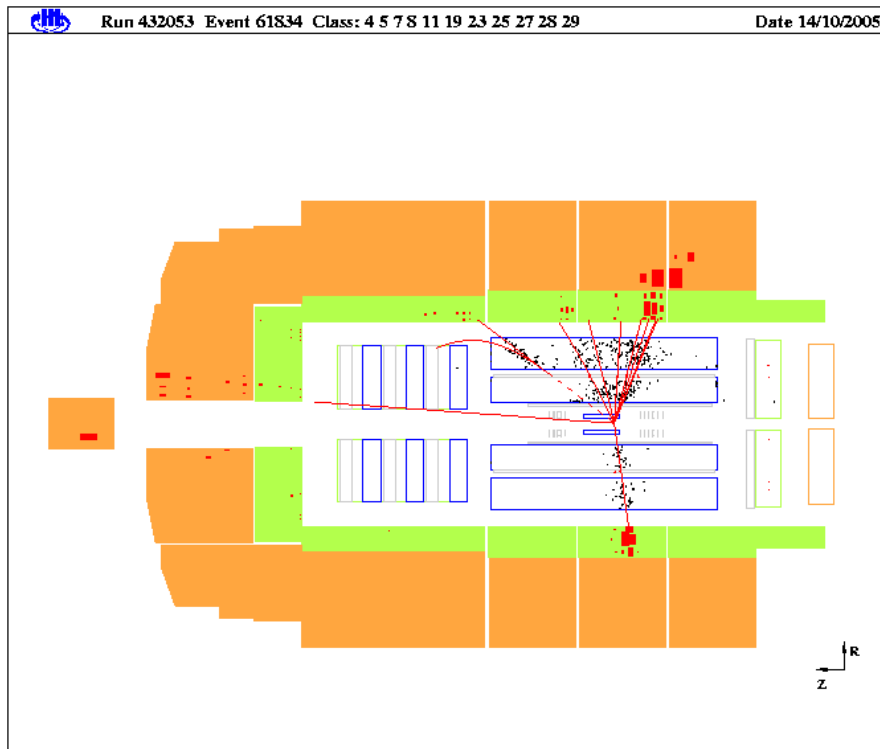
expected

- proton structure
 - structure functions
 - parton densities
 - α_s
- photon structure
- perturbative QCD
 - jets
 - heavy quarks
- electroweak

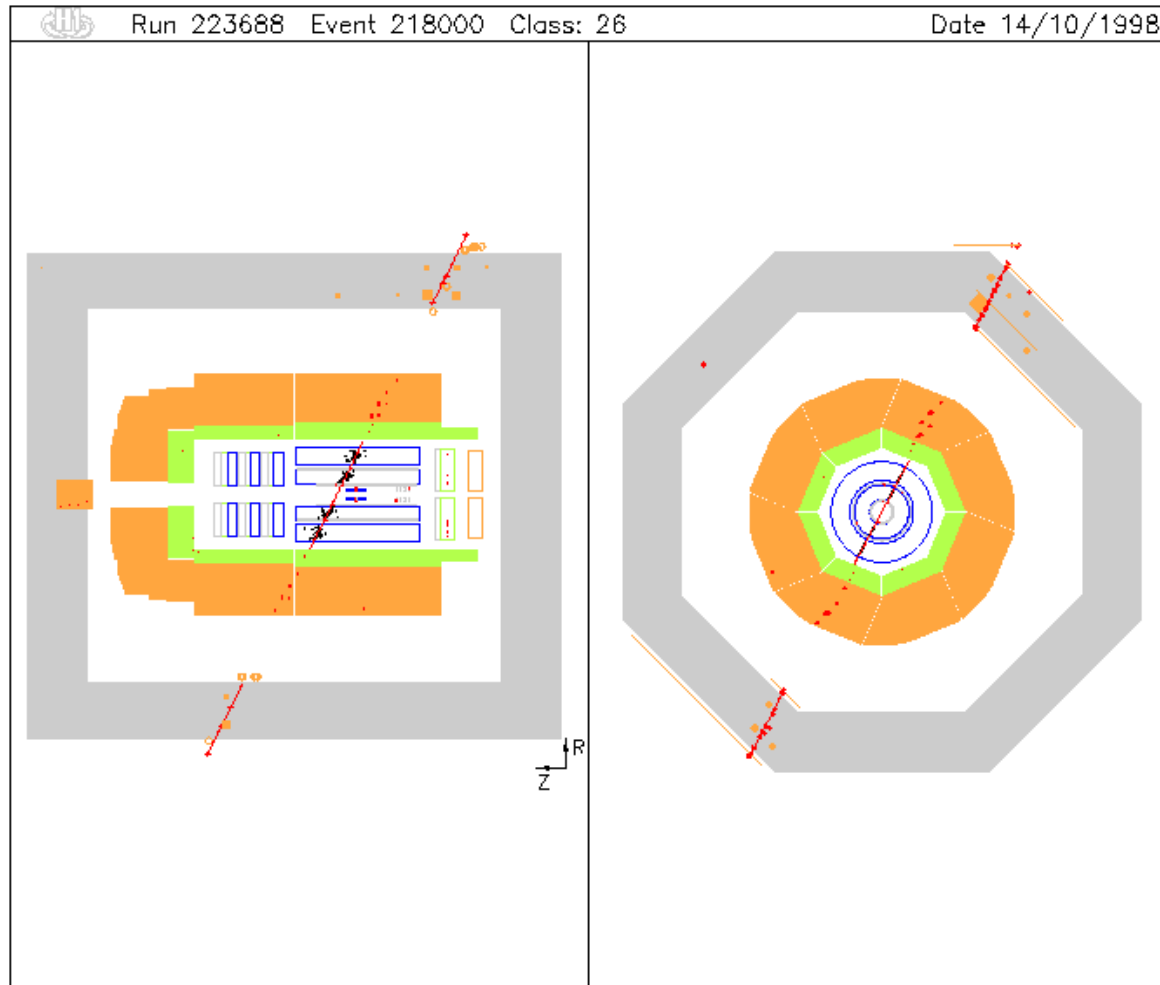
not (so) expected

- exotics (beyond the standard model)
 - SUSY
 - leptoquarks
 - ...
- diffraction

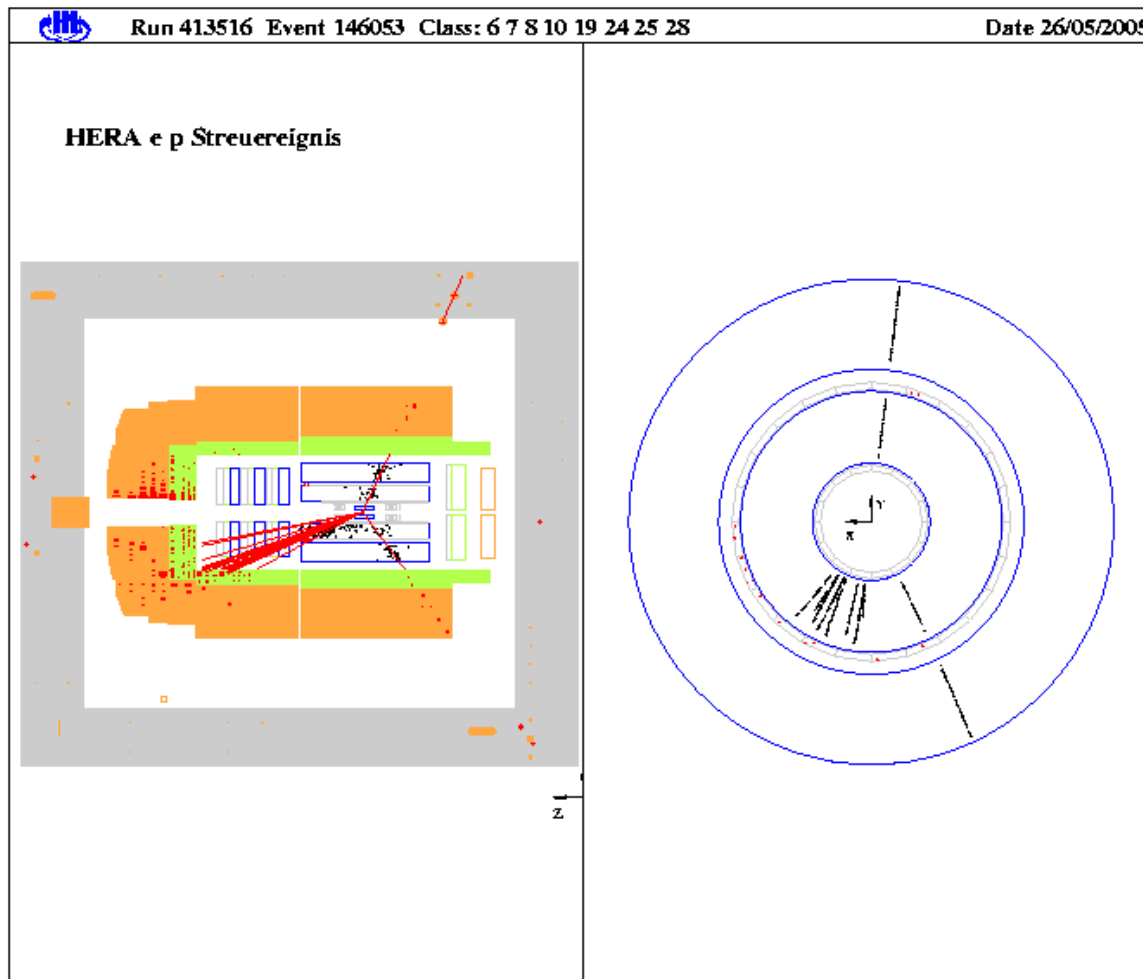
Some Events...



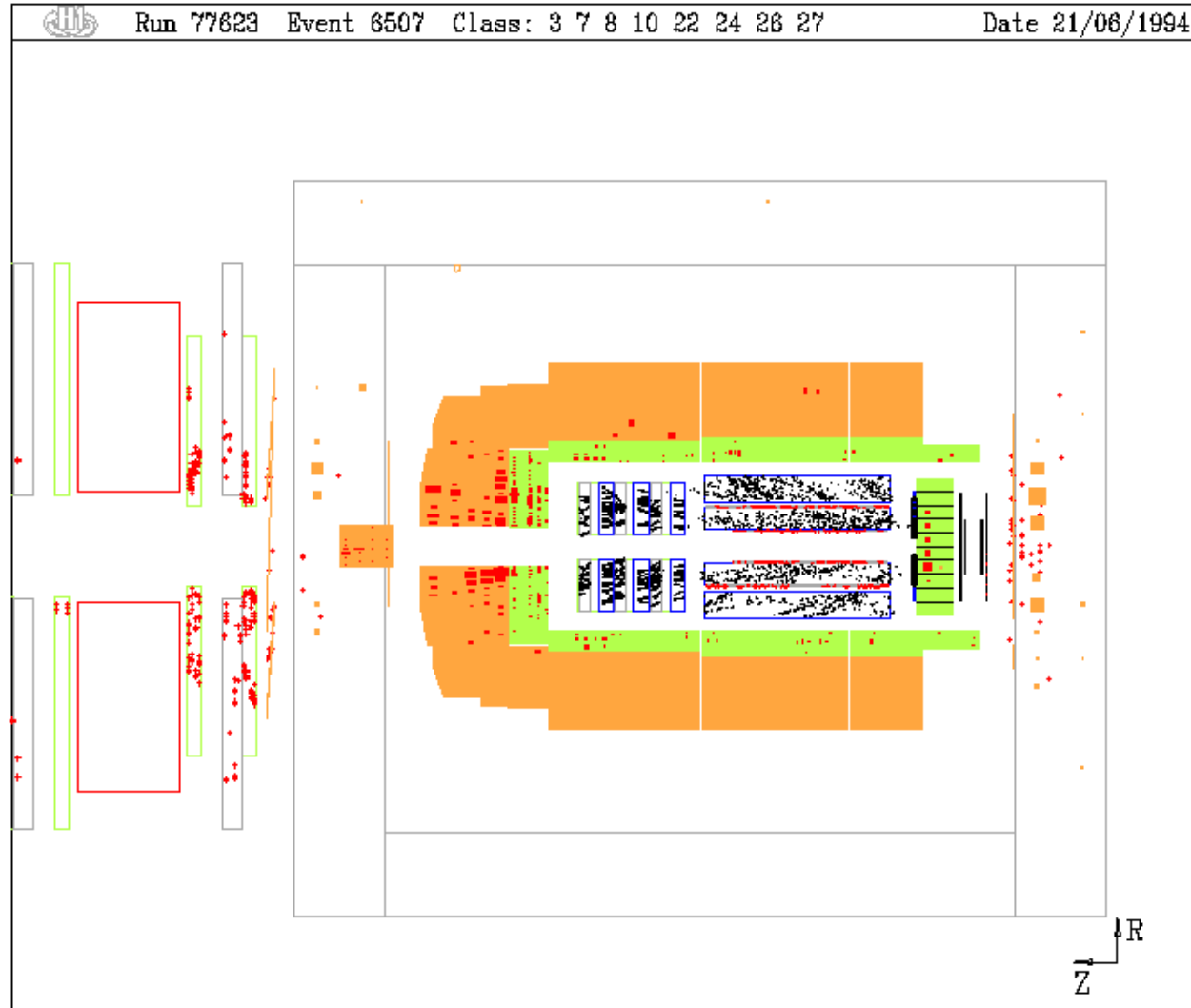
Some Events...



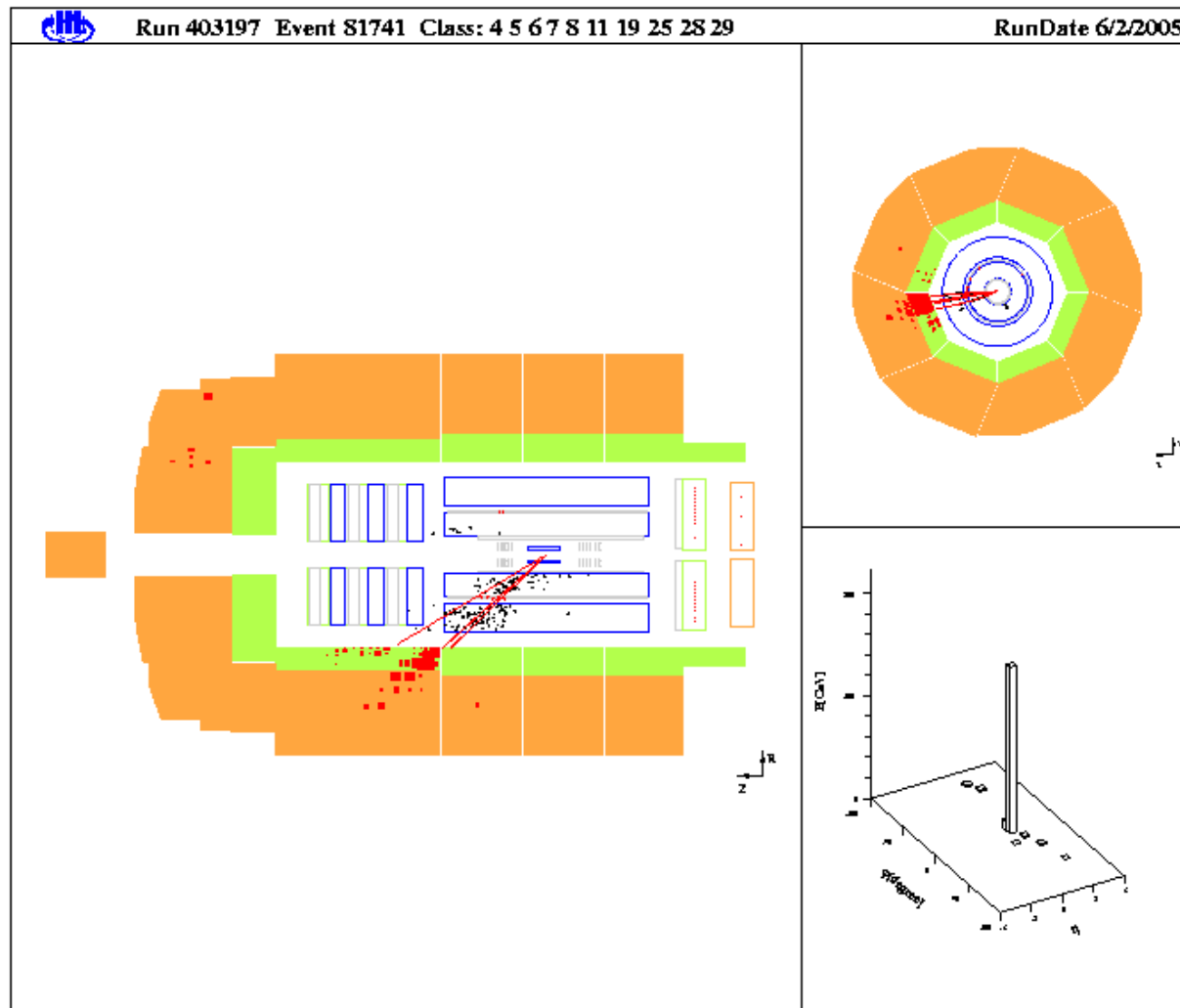
Some Events...



Some Events...



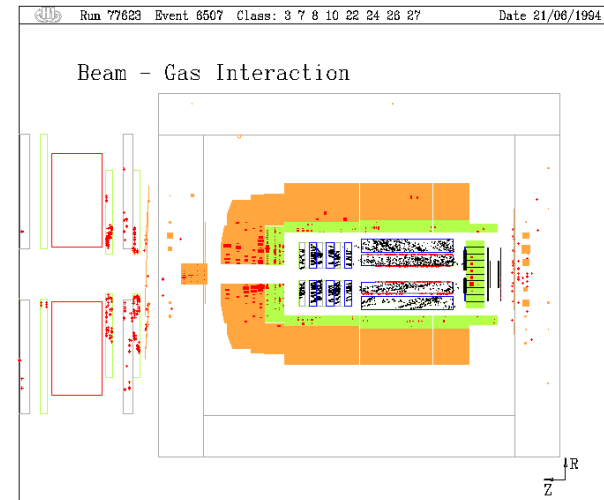
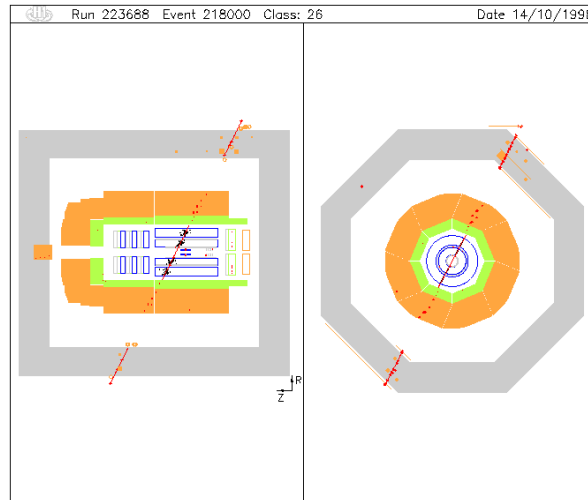
Some Events...



Event Rates

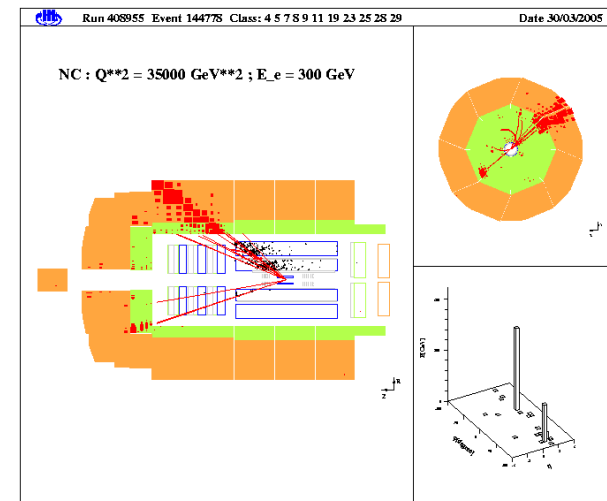
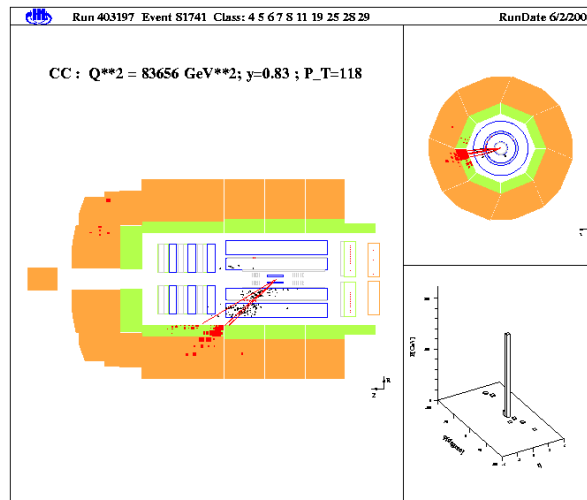
some kHz

some Hz



some min^{-1}

some hour^{-1}

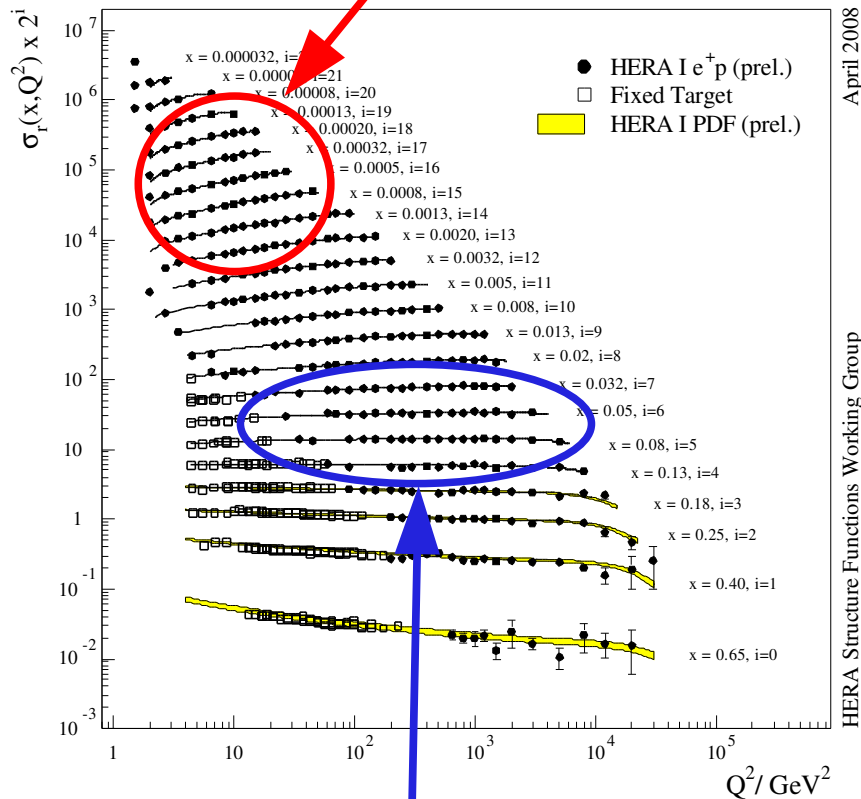


ep Scattering & Structure Functions

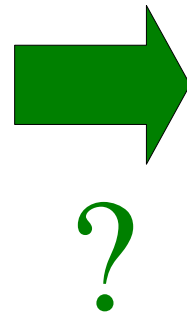
„The“ HERA Textbook Plots

gluons ?

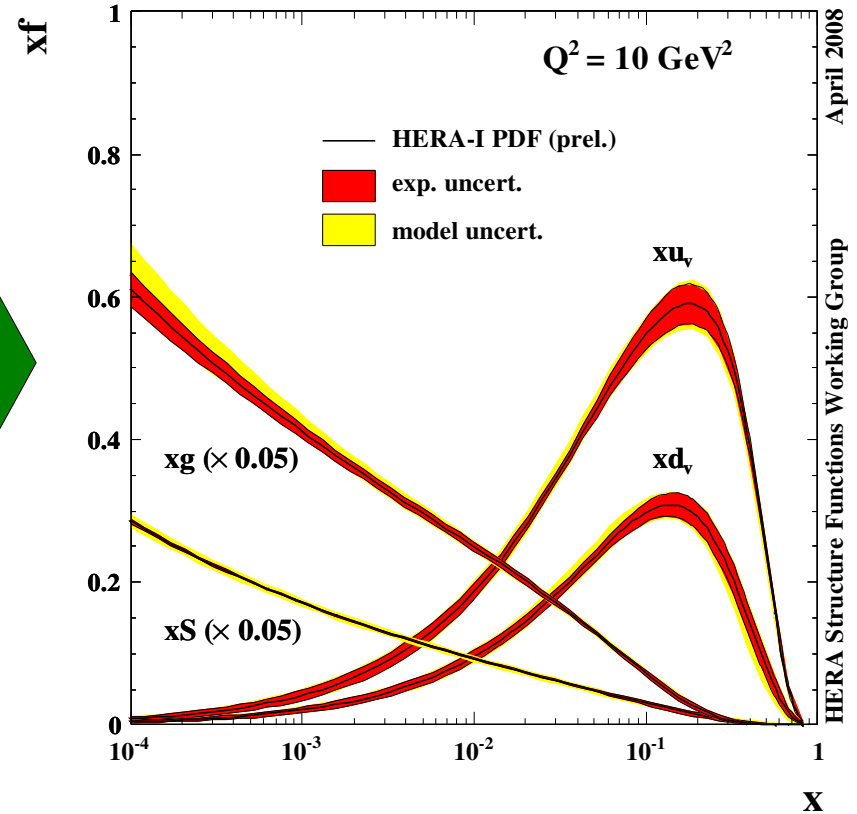
H1 and ZEUS Combined PDF Fit



quarks ?

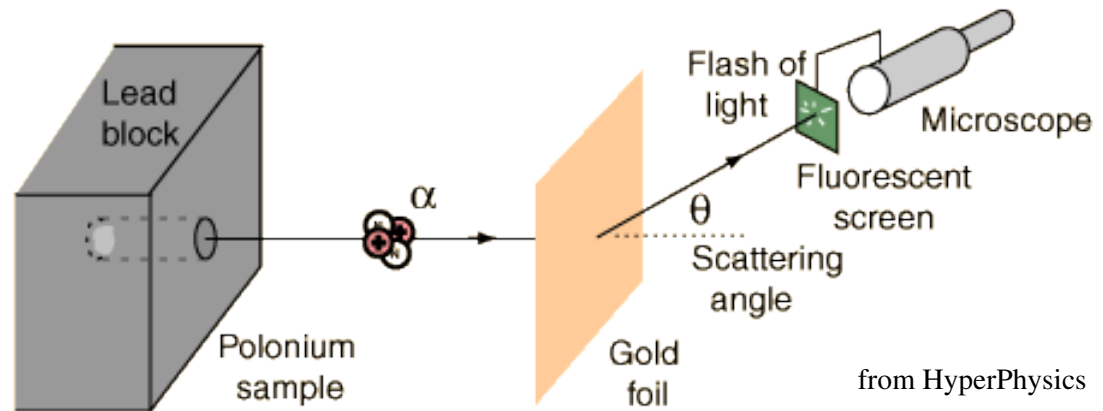


H1 and ZEUS Combined PDF Fit



Rutherford Scattering

- first scattering experiment
- existence of the nucleus



$$\frac{d\sigma}{d\Omega} = \left(\frac{1}{4\pi\epsilon_0} \frac{Z_1 Z_2 e^2}{4E_{kin}} \right)^2 \frac{1}{\sin^4 \frac{\Theta}{2}}$$

assumes

- Coulomb potential
- no spins
- no recoil

Elastic Electron Scattering

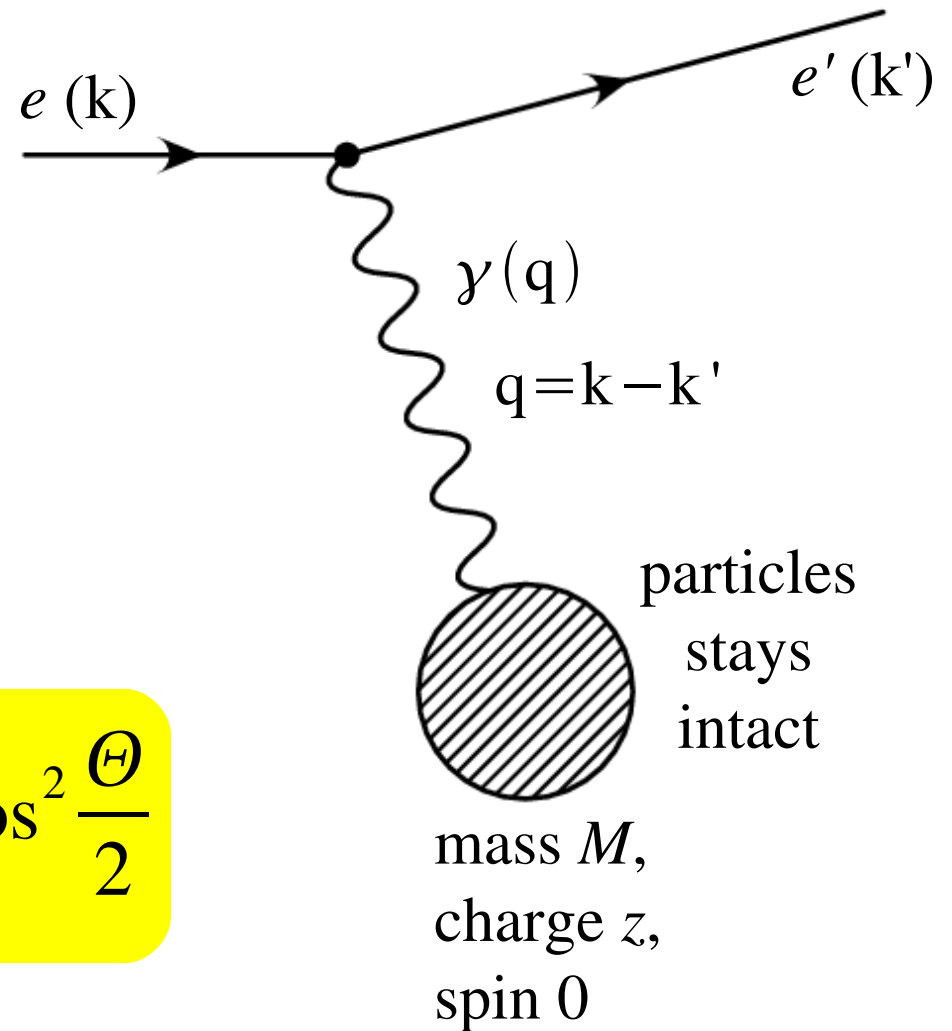
variables:

- $q = k - k'$
 - $Q^2 = -q^2$
 $= 4 E E' \sin^2(\Theta/2)$
 - $E' = \frac{E}{1 + (2 E/M) \sin^2(\Theta/2)}$
- only one independent!

$$\frac{d\sigma}{dQ^2} = \frac{4\pi\alpha^2 z^2}{Q^4} \left(\frac{E'}{E}\right)^2 \cos^2 \frac{\Theta}{2}$$

Coulomb-
Potential $\sim 1/r$

recoil



Elastic Electron Scattering: Cross Section

- Mott Scattering: electron on a pointlike charged particle with spin 0

$$\left(\frac{d\sigma}{dQ^2} \right)_{\text{Mott}} = \frac{4\pi\alpha^2}{Q^4} \left(\frac{E'}{E} \right)^2 \cos^2 \frac{\Theta}{2}$$

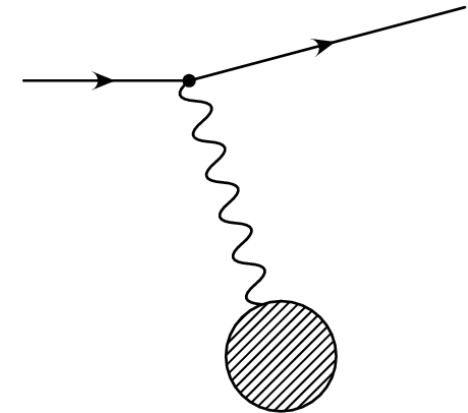
- Dirac Scattering: electron on a pointlike charged particle with spin 1/2

$$\left(\frac{d\sigma}{dQ^2} \right)_{\text{Dirac}} = \left(\frac{d\sigma}{dQ^2} \right)_{\text{Mott}} \left[1 + 2\tau \tan^2 \frac{\Theta}{2} \right] \quad \text{with} \quad \tau = \frac{Q^2}{4M^2}$$

- electron on proton: „form factors“ needed:

$$\left(\frac{d\sigma}{dQ^2} \right)_{ep} = \left(\frac{d\sigma}{dQ^2} \right)_{\text{Mott}} \left[\frac{G_E^2(Q^2) + \tau G_M^2(Q^2)}{1 + \tau} + 2\tau G_M^2(Q^2) \tan^2 \frac{\Theta}{2} \right]$$

→ protons are not pointlike!



Electric Form Factor of the Proton

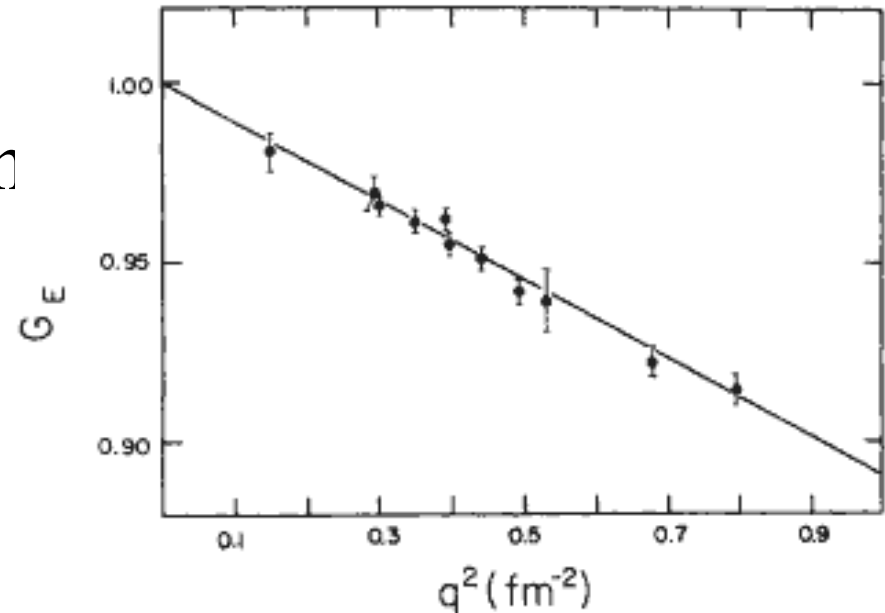
- describes the charge distribution in the proton (Fourier transform)
- measured:

- $G_E(0) = 1$

- $G_M(0) = 2.79$

- $G_E(Q^2), G_M(Q^2) \propto \left(1 + \frac{Q^2}{0.71 \text{ GeV}^2}\right)^{-2}$

→ elastic scattering only import at low Q^2

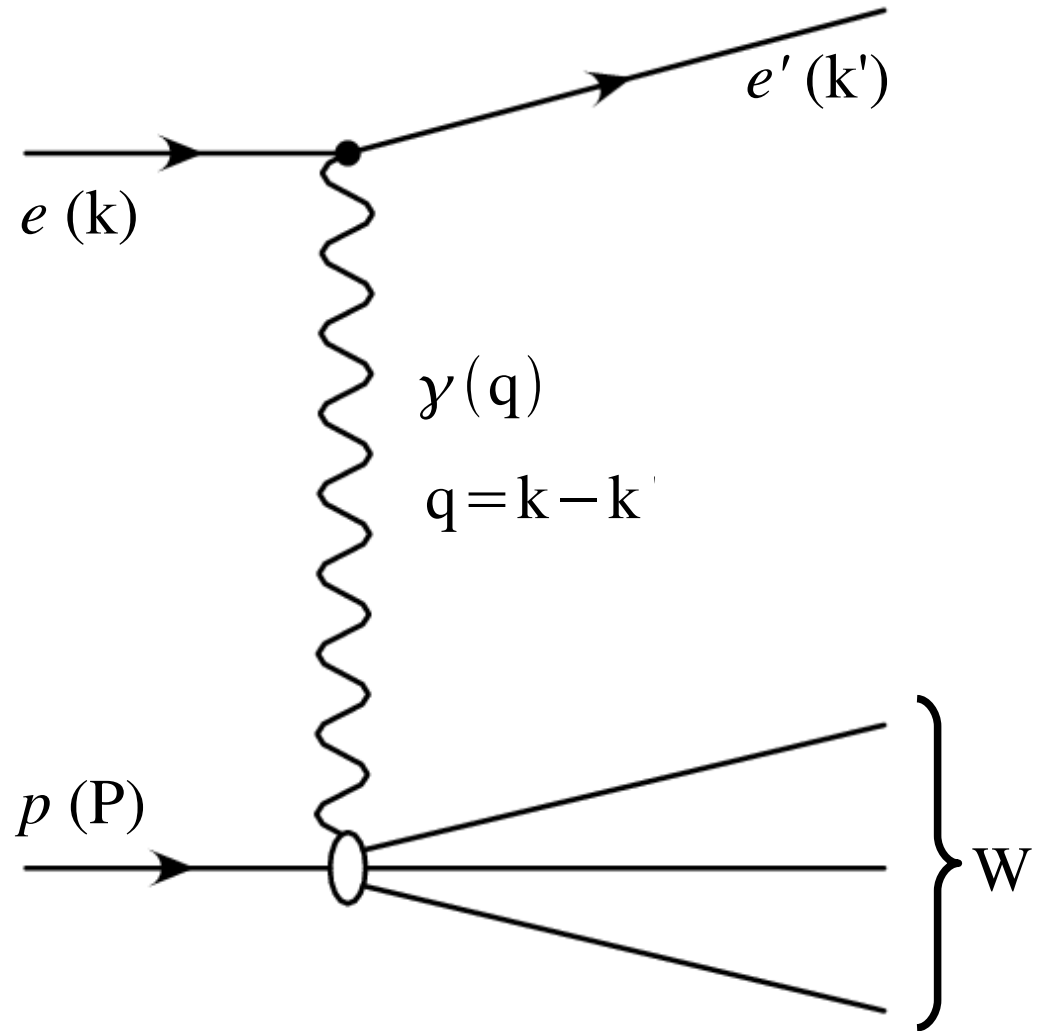


from J.J. Murphy et al., „Proton form factor from 0.15 to 0.79 fm⁻²“

Inelastic Electron Scattering

variables:

- $q = k - k'$
 - $Q^2 = -q^2$
 - $s = (P + k)^2$
 - $W^2 = (P + q)^2$
 $= M^2 + 2q \cdot P - Q^2$
 - $y = q \cdot P / k \cdot P$
- two independent!



elastic: $W = M$

inelastic: $W > M$

Inelastic Electron Proton Scattering

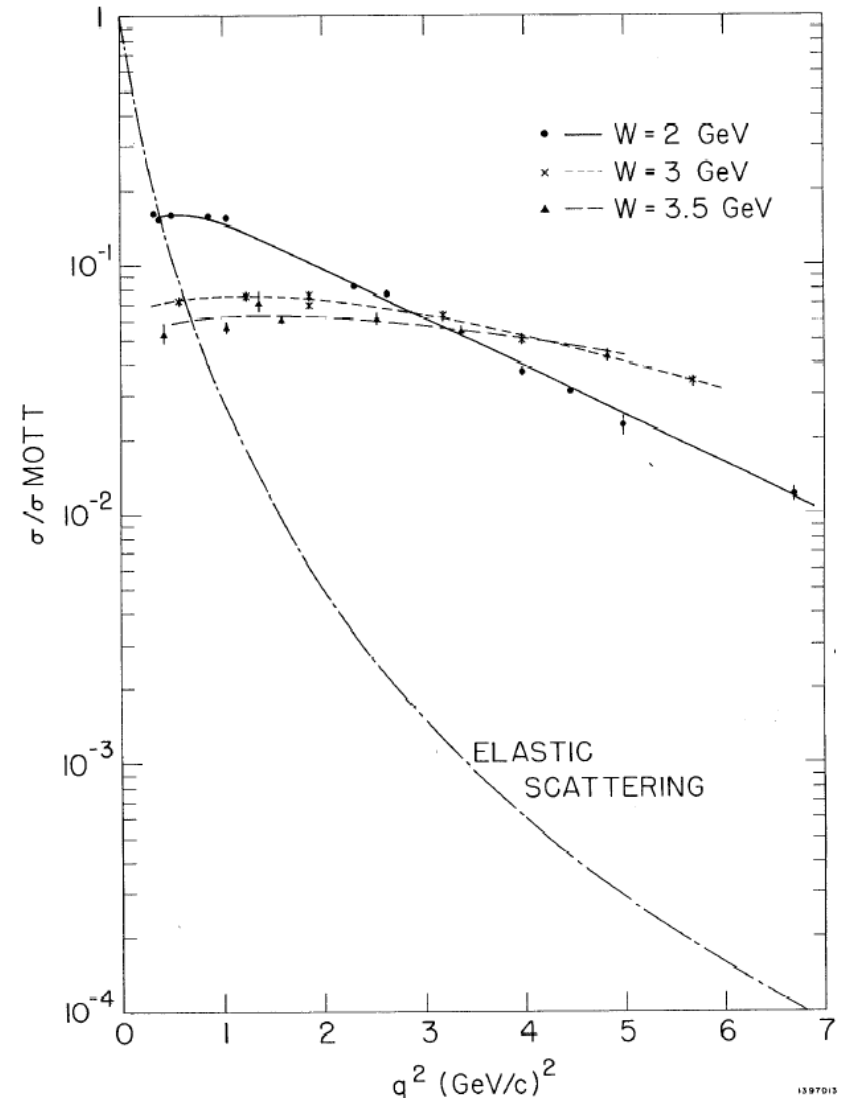
- inelastic scattering:
 $W > M_p$
- ratio to Mott cross section
nearly flat in Q^2

SLAC-PUB-650
August 1969
(EXP) and (TH)

OBSERVED BEHAVIOR OF HIGHLY INELASTIC
ELECTRON-PROTON SCATTERING

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E. D. Bloom, D. H. Coward, H. DeStaebler,
J. Drees, L. W. Mo, R. E. Taylor
Stanford Linear Accelerator Center, † Stanford, California 94305

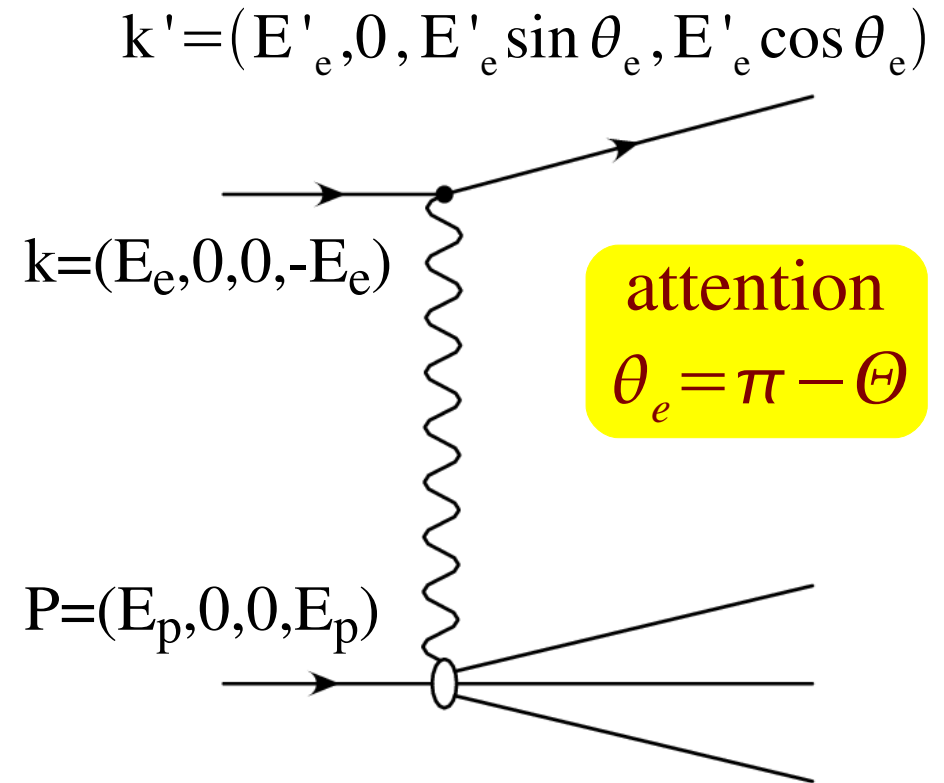


Deep Inelastic Scattering (DIS)

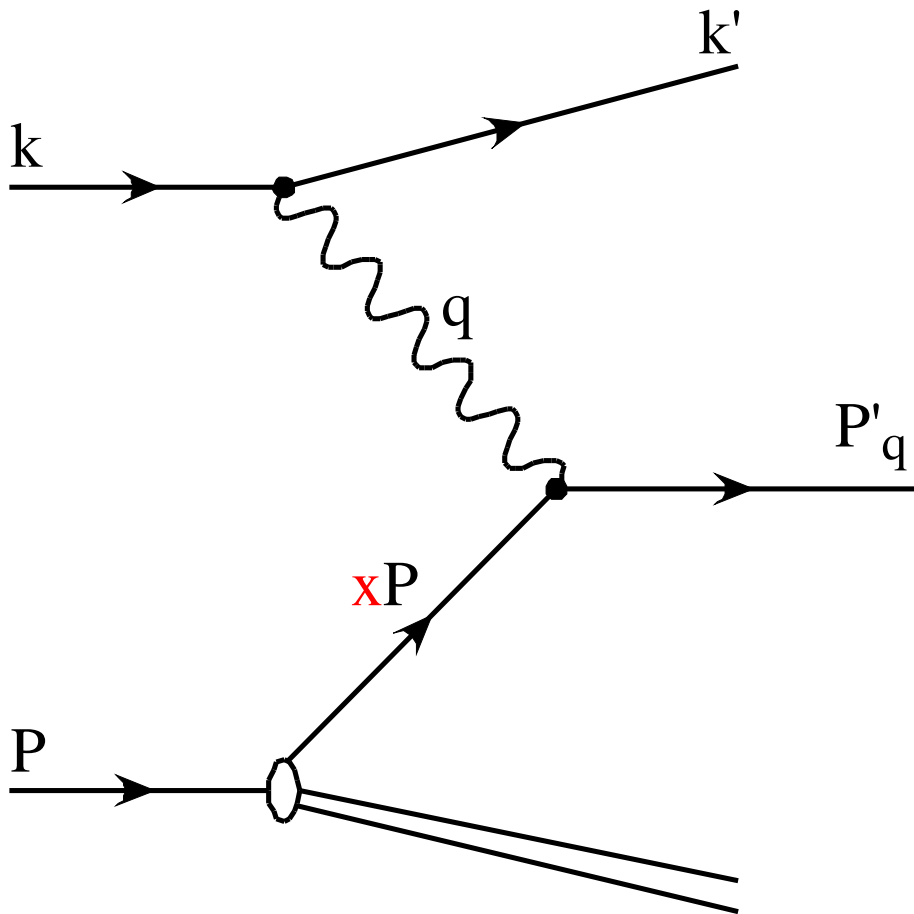
- deep: $Q^2 > M_p$
- inelastic: $W > M_p$
- for HERA: $m_e, M_p \ll W$
 → neglect m_e, M_p

- $s = 4 E_p E_e$
- $Q^2 = 2 E_e E'_e (1 + \cos \theta_e)$
- $y = 1 - \frac{E'_e}{E_e} \sin^2 \frac{\theta_e}{2}$
- $W = y \sqrt{s} - Q^2$

- one more variable: $x = Q^2 / (2 P \cdot q) = Q^2 / ys$



DIS: What is x ?



x can be interpreted as the momentum fraction of the struck parton of the proton:

$$P'_q = q + xP$$

$$(q + xP)^2 = -Q^2 + 2x q \cdot P + (xP)^2$$

$$(q + xP)^2 = (xP)^2 = (m_q)^2$$

$$x = \frac{Q^2}{2q \cdot P} = \frac{Q^2}{ys}$$

inelastic proton scattering is scattering on a parton of the proton!

Structure Functions F_1 & F_2

- the DIS cross section can be written as

$$\begin{aligned}\frac{d^2 \sigma}{dx dQ^2} &= \frac{4 \pi \alpha^2}{Q^4} \frac{1}{x} \left[(1-y) F_2(x, Q^2) + \frac{y^2}{2} 2x F_1(x, Q^2) \right] \\ &= \frac{4 \pi \alpha^2}{Q^4} \frac{1}{x} \frac{E'}{E} \left[F_2(x, Q^2) \cos^2 \frac{\Theta}{2} + \frac{Q^2}{2x^2 M_p^2} 2x F_1(x, Q^2) \sin^2 \frac{\Theta}{2} \right]\end{aligned}$$

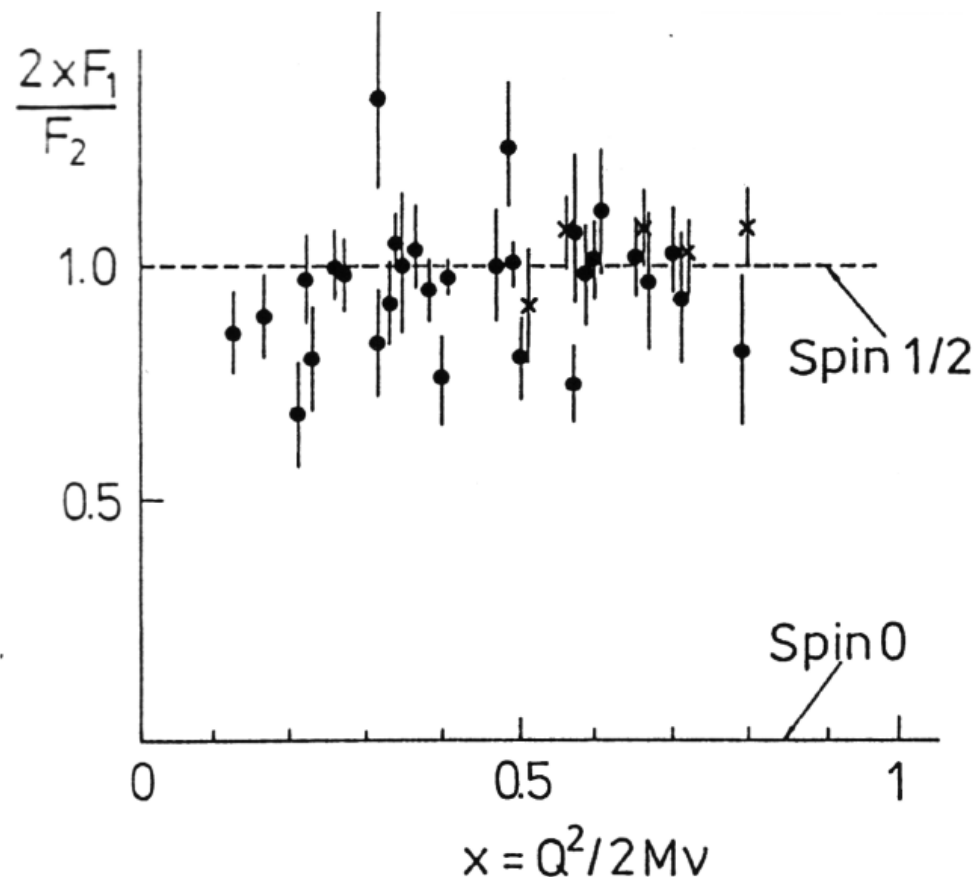
- comparison with Dirac formula

$$\left(\frac{d\sigma}{dQ^2} \right)_{\text{Dirac}} = \frac{4 \pi \alpha^2 z^2}{Q^4} \left(\frac{E'}{E} \right)^2 \left[\cos^2 \frac{\Theta}{2} + \frac{Q^2}{2M^2} \sin^2 \frac{\Theta}{2} \right]$$

- F_2 corresponds to **electric** field of the parton
- F_1 corresponds to **spin** of the parton

Parton Spin

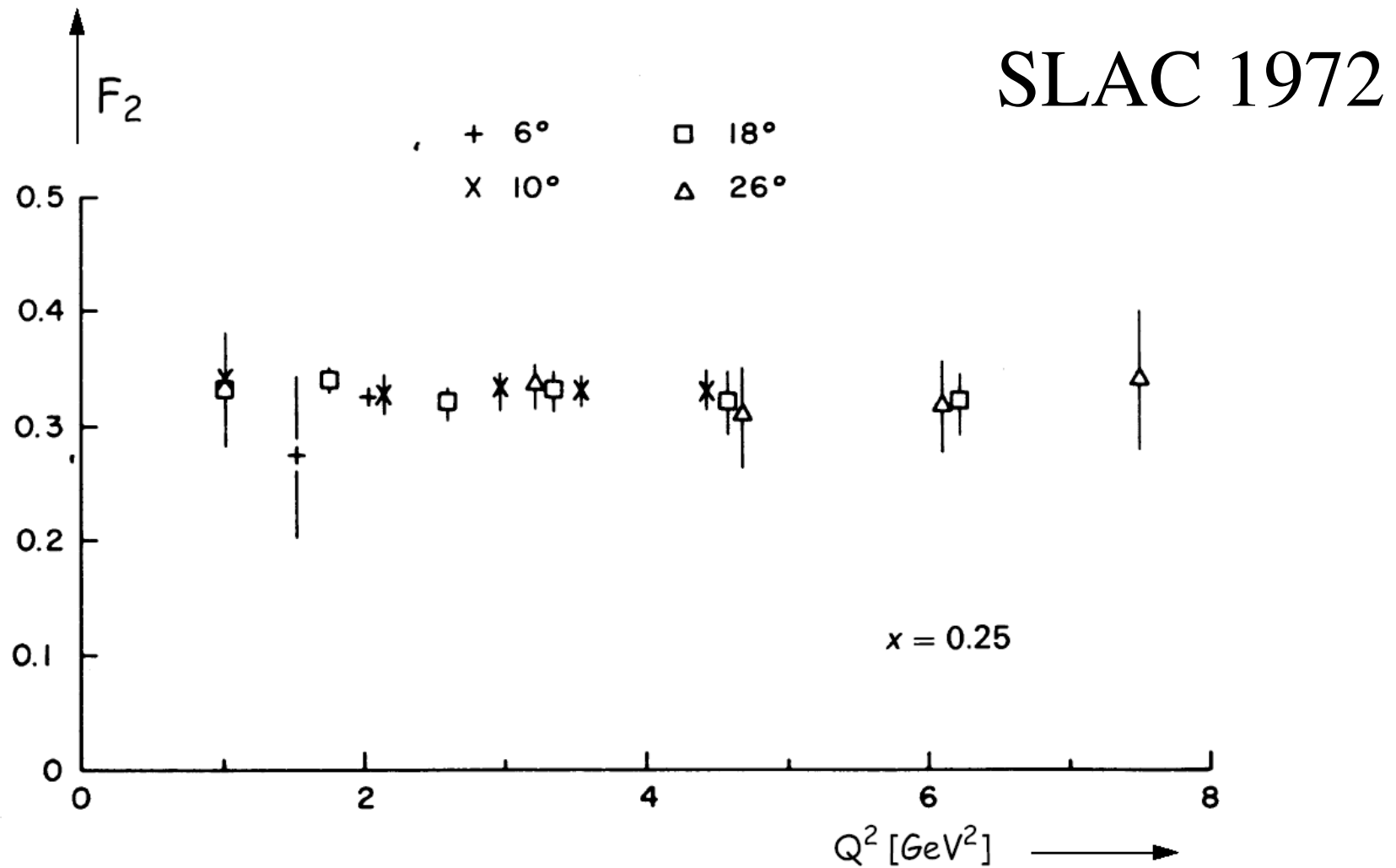
- parton spin $\frac{1}{2}$: $2 \times F_1 = F_2$ (Callan Gross)
- parton spin 0: $2 \times F_1 = 0$



partons
have spin $\frac{1}{2}$

from P. Schmüser, „Feynman-Graphen und Eichtheorien für Experimentalphysiker“

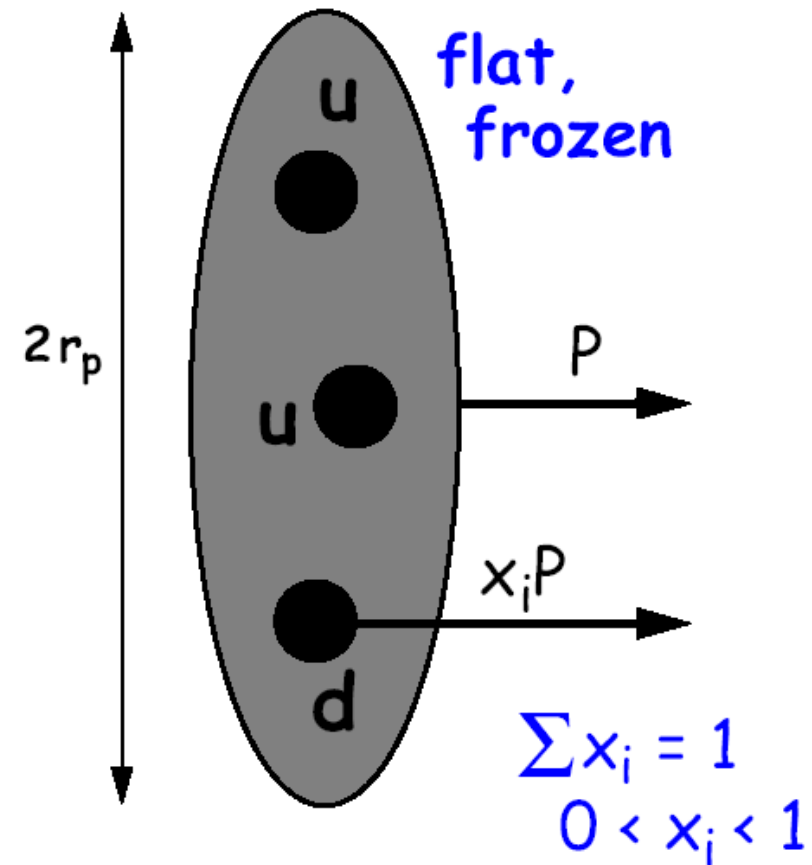
Scaling: F_2 independent of Q^2



independent of Q^2 , we always see the same partons (=quarks)

(Naive) Quark Parton Model

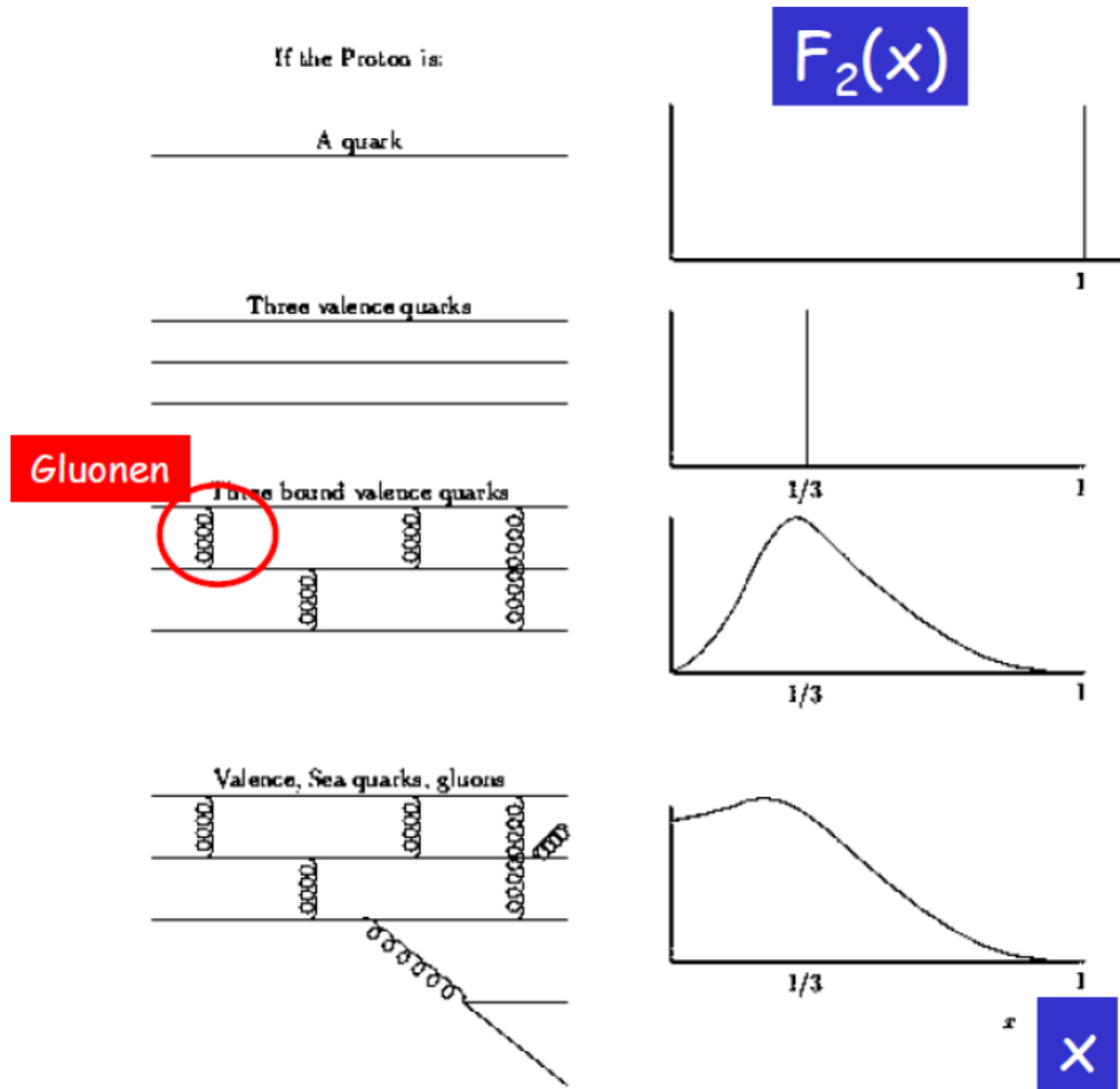
- proton consists of 3 partons, identified with the QCD quarks
- during the interaction proton is „frozen“
- electron proton scattering is sum of incoherent electron quark scatterings
- proton structure is defined by **parton distributions**



$$F_2(x, Q^2) = x \sum e_q^2 q(x)$$

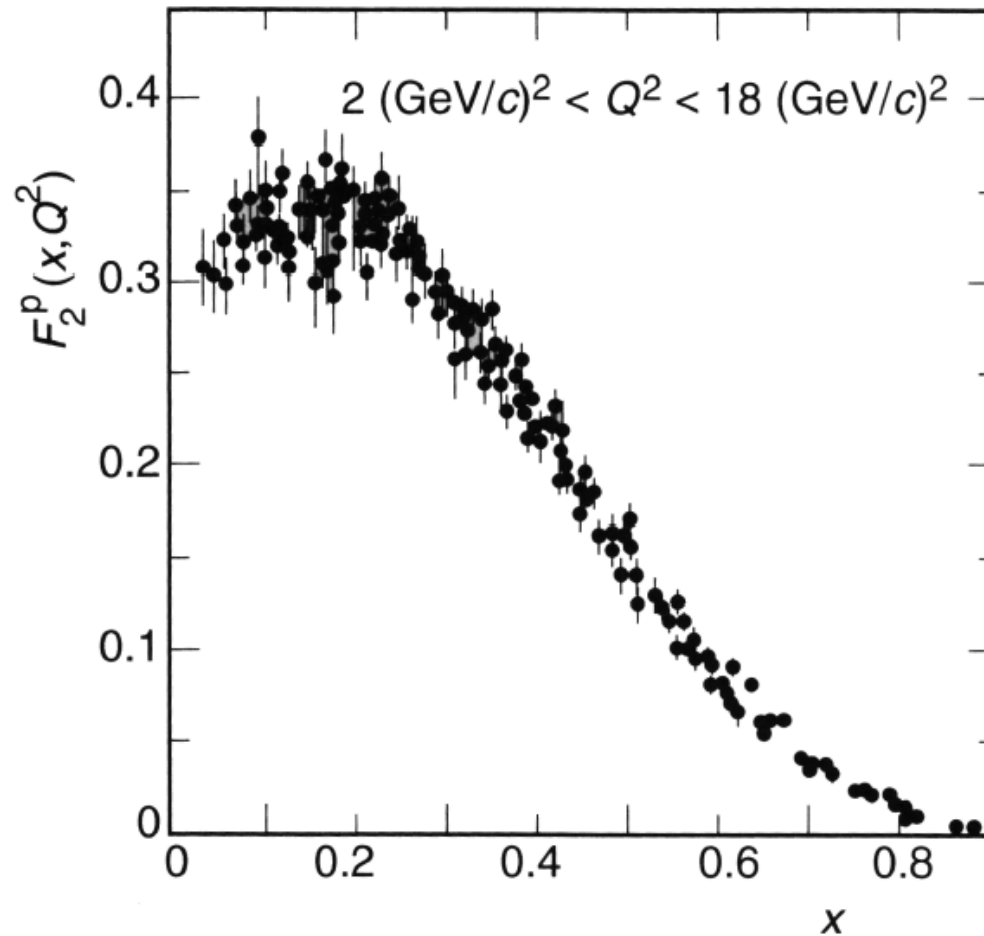
How does $F_2(x)$ look like?

How does $F_2(x)$ look like?



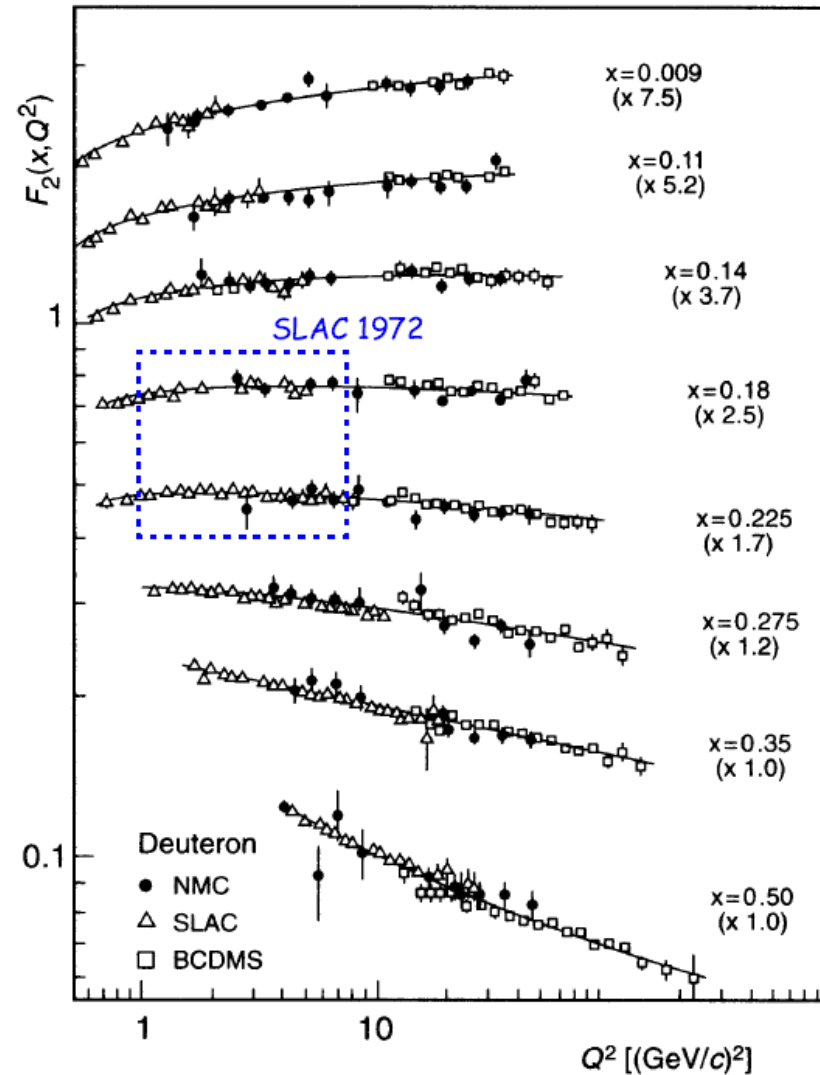
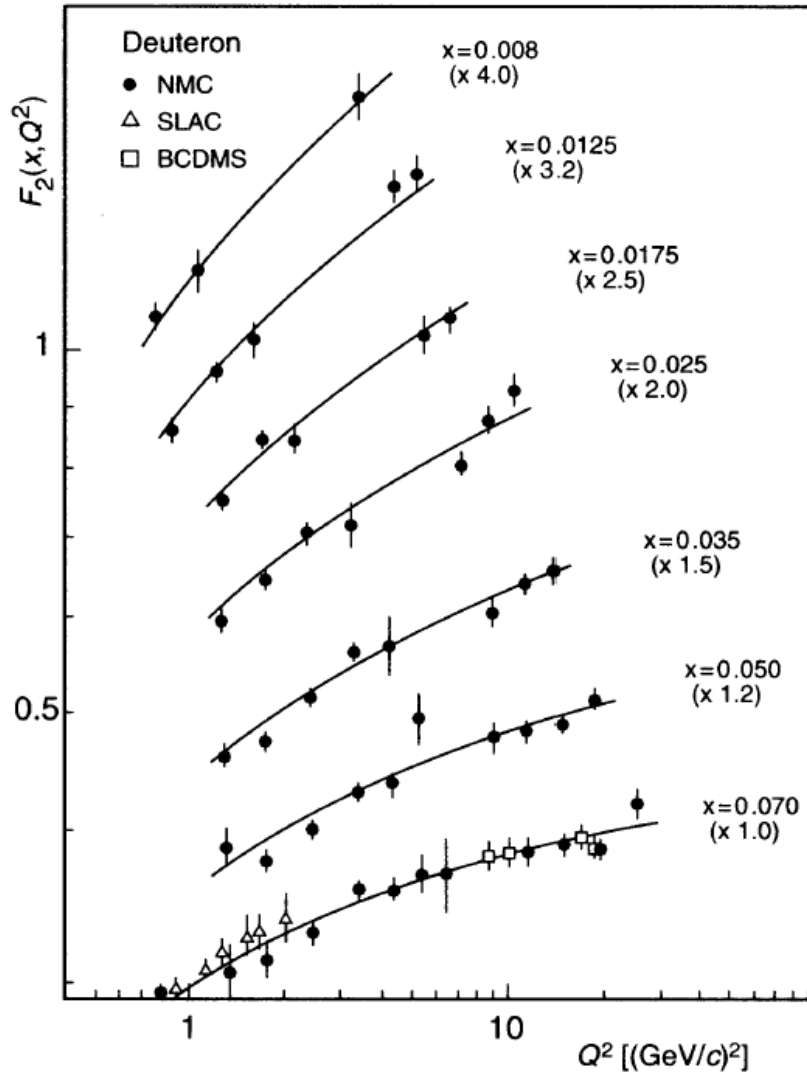
How does $F_2(x)$ look like?

what happens
at low x ?



from Povh et al., „Teilchen und Kerne“

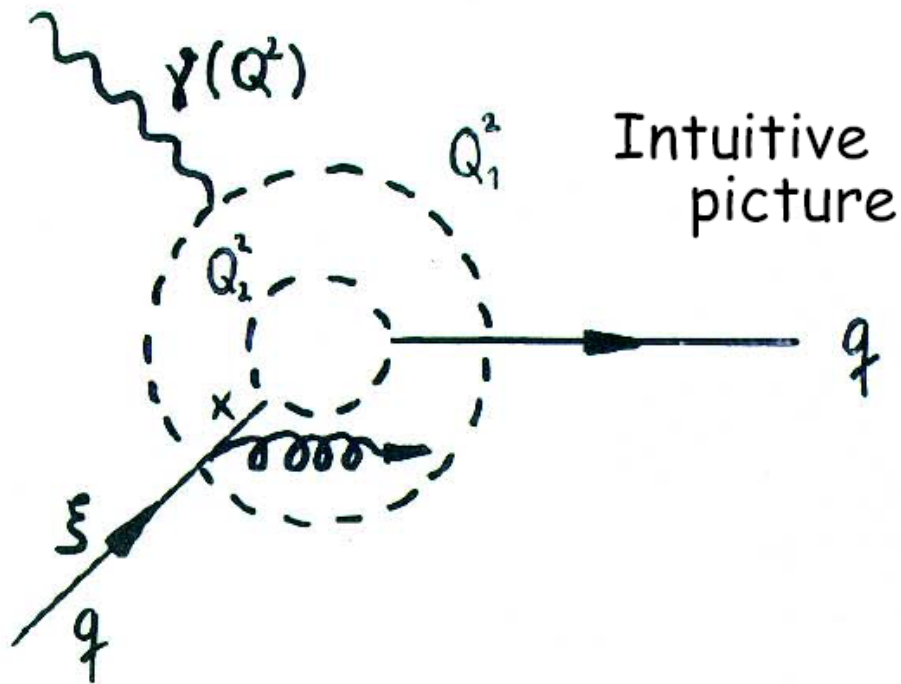
Scaling Violations



at smaller & larger x , the amount of quarks depends on Q^2 !

Parton Evolution

- number of partons changes with Q^2
- Q^2 can be interpreted as resolving power: $Q^2 \propto (\hbar/\lambda)^2$



small Q^2 :

- many partons with large x
- (nearly) no partons at low x

large Q^2 :

- less partons with large x
- more partons at low x

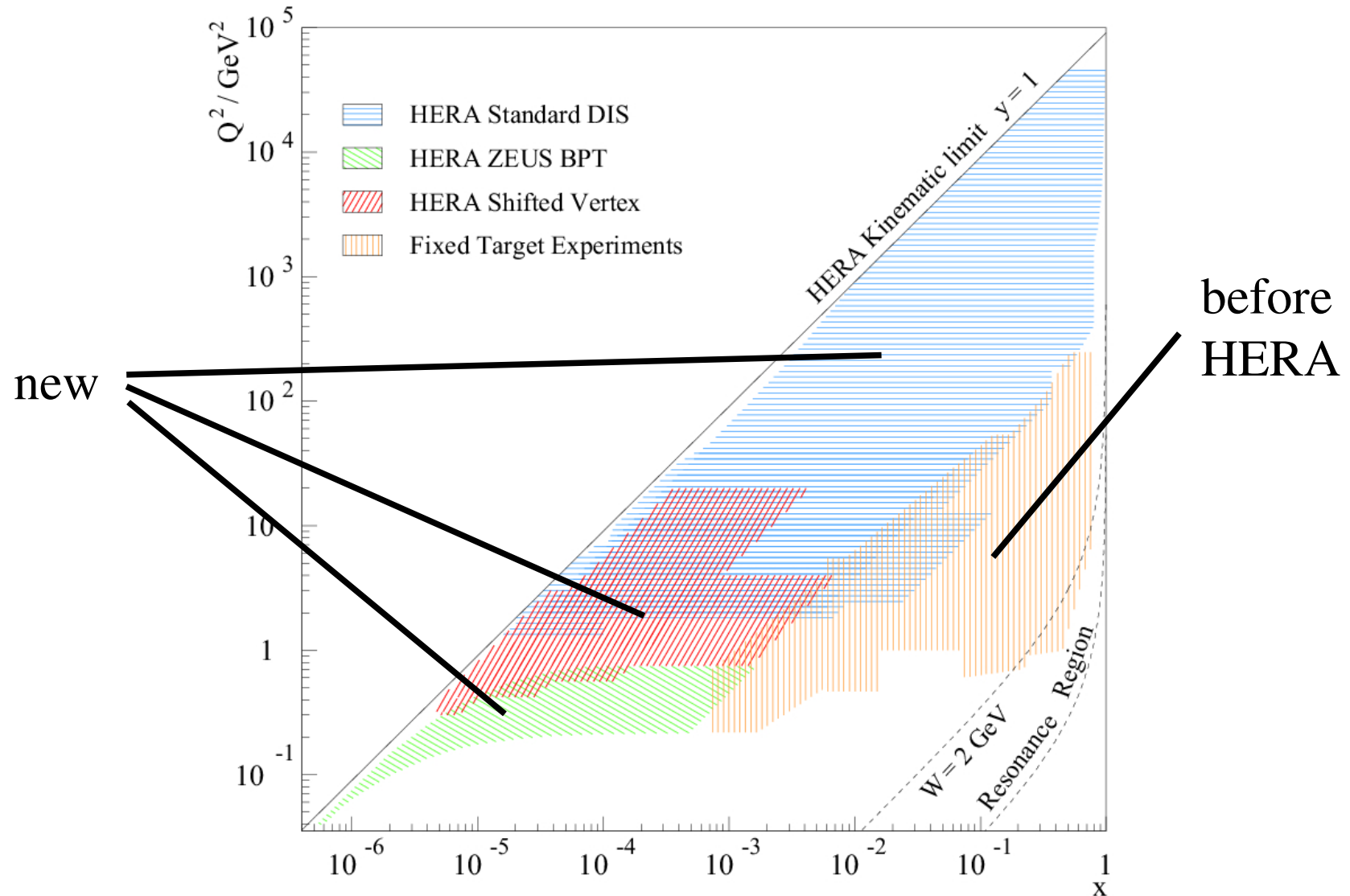
DGLAP Evolution Equations

$$\frac{\partial}{\partial \log Q^2} \begin{bmatrix} q(x, Q^2) \\ g(x, Q^2) \end{bmatrix} = \frac{\alpha_s}{2\pi} \begin{bmatrix} P_{q/q} \left[\begin{array}{c} \gamma \\ x \end{array} \right] & P_{q/g} \left[\begin{array}{c} \gamma \\ x \end{array} \right] \\ P_{g/q} \left[\begin{array}{c} \gamma \\ x \end{array} \right] & P_{g/g} \left[\begin{array}{c} \gamma \\ x \end{array} \right] \end{bmatrix} \otimes \begin{bmatrix} q(x, Q^2) \\ g(x, Q^2) \end{bmatrix}$$

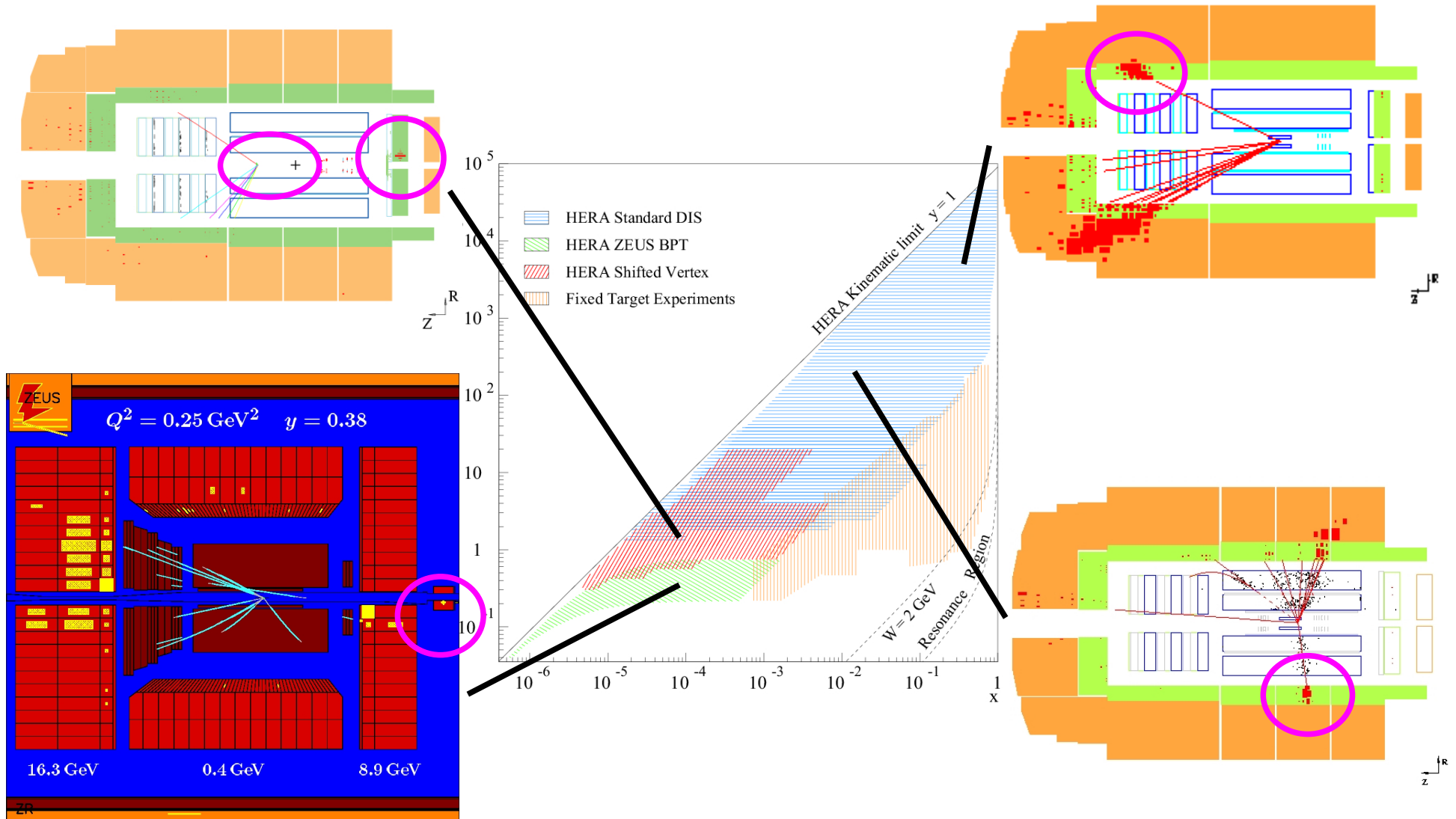
$P \otimes f(x, Q^2) = \int_x^1 \frac{dy}{y} P(x/y) f(y, Q^2)$

- Q^2 dependence of quark densities $q(x, Q^2)$ and gluon density $g(x, Q^2)$ is predicted
- no prediction for the x dependence \rightarrow initial condition needed

HERA Kinematic Range



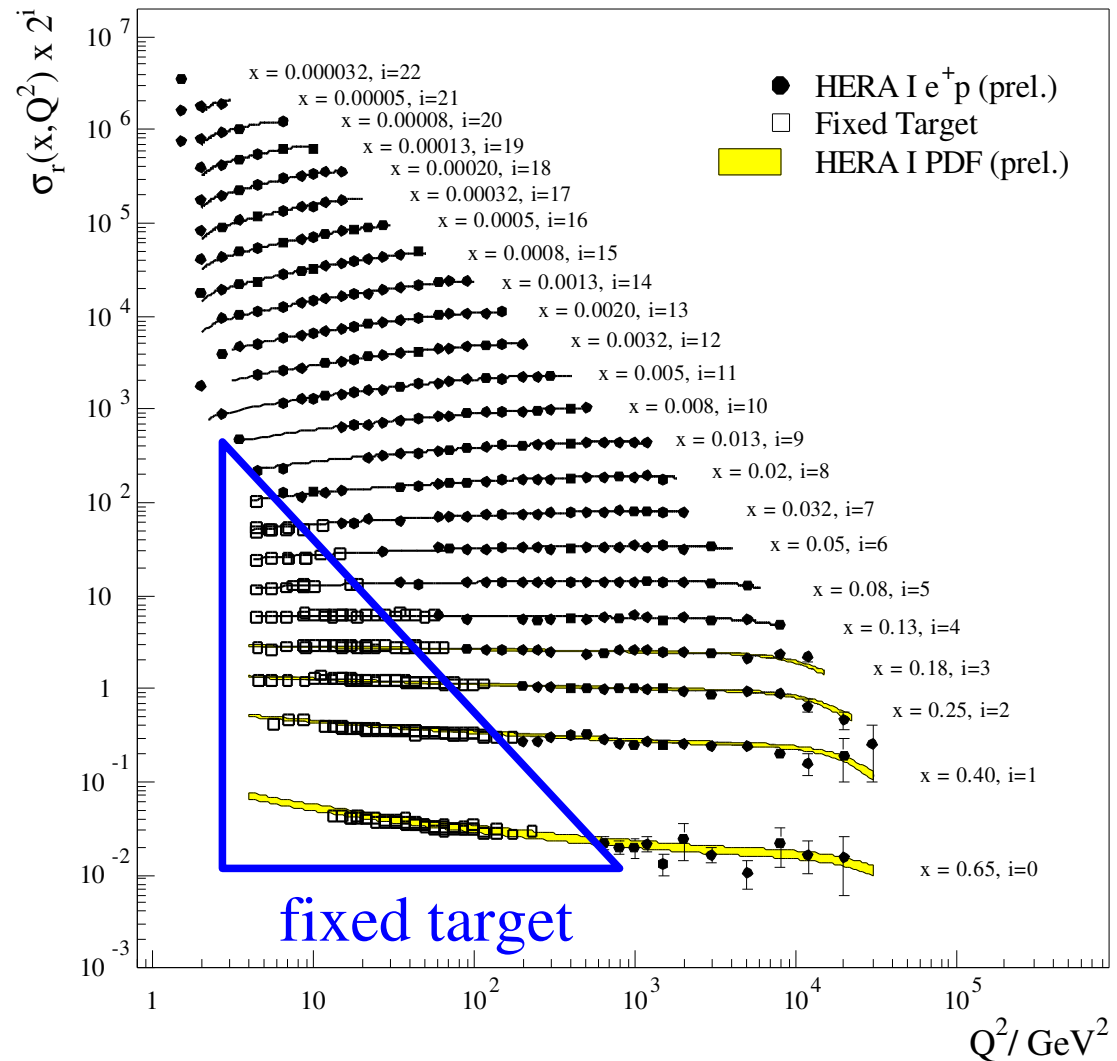
Events in Different Regions



F_2 vs. Q^2

- HERA data cover huge range: 5 orders in Q^2 and 4 orders in x
- approximate scaling at large x
- clear scaling violations at small x

H1 and ZEUS Combined PDF Fit



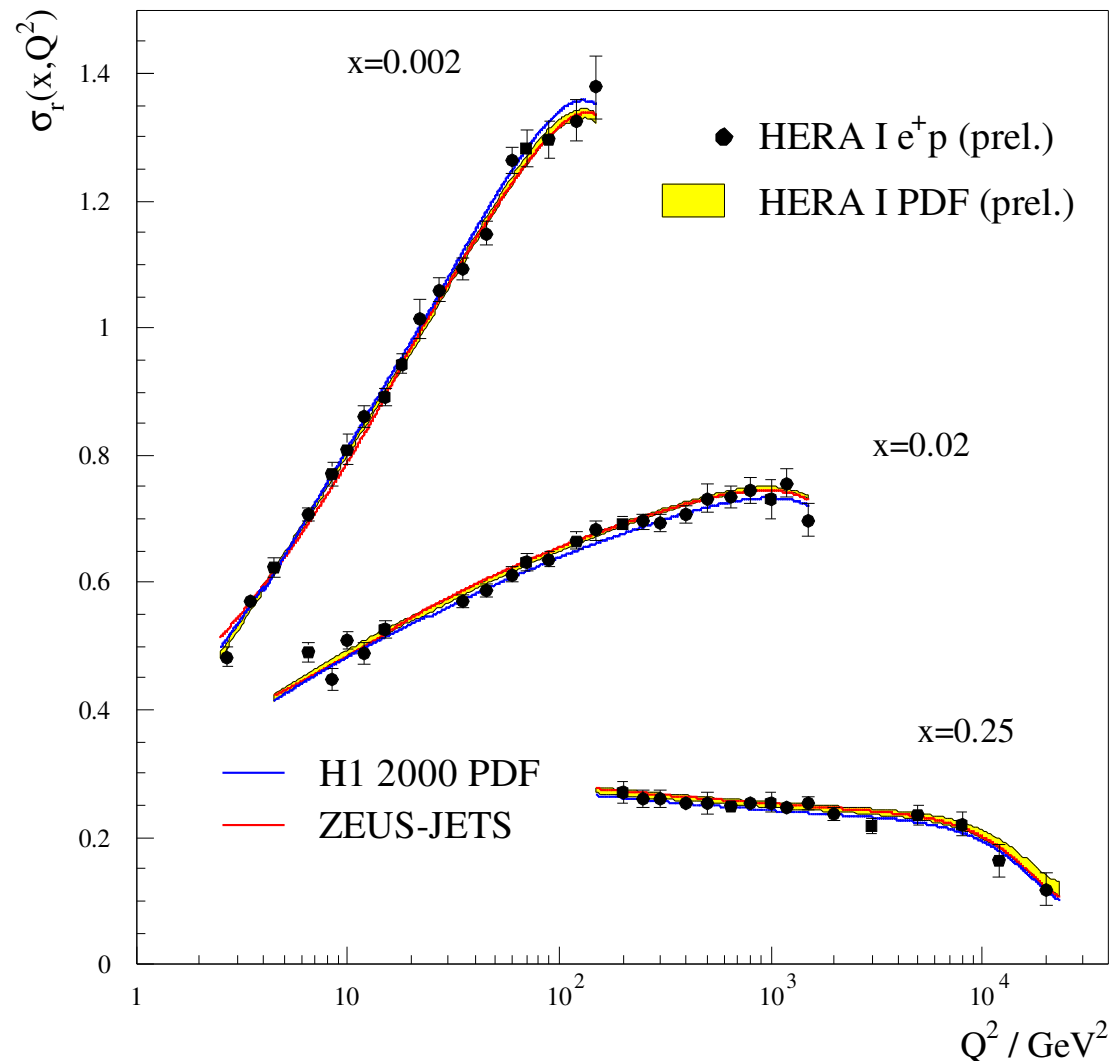
April 2008

HERA Structure Functions Working Group

F_2 vs. Q^2 : example bins

- clear scaling violations at small x
- approximate scaling at large x

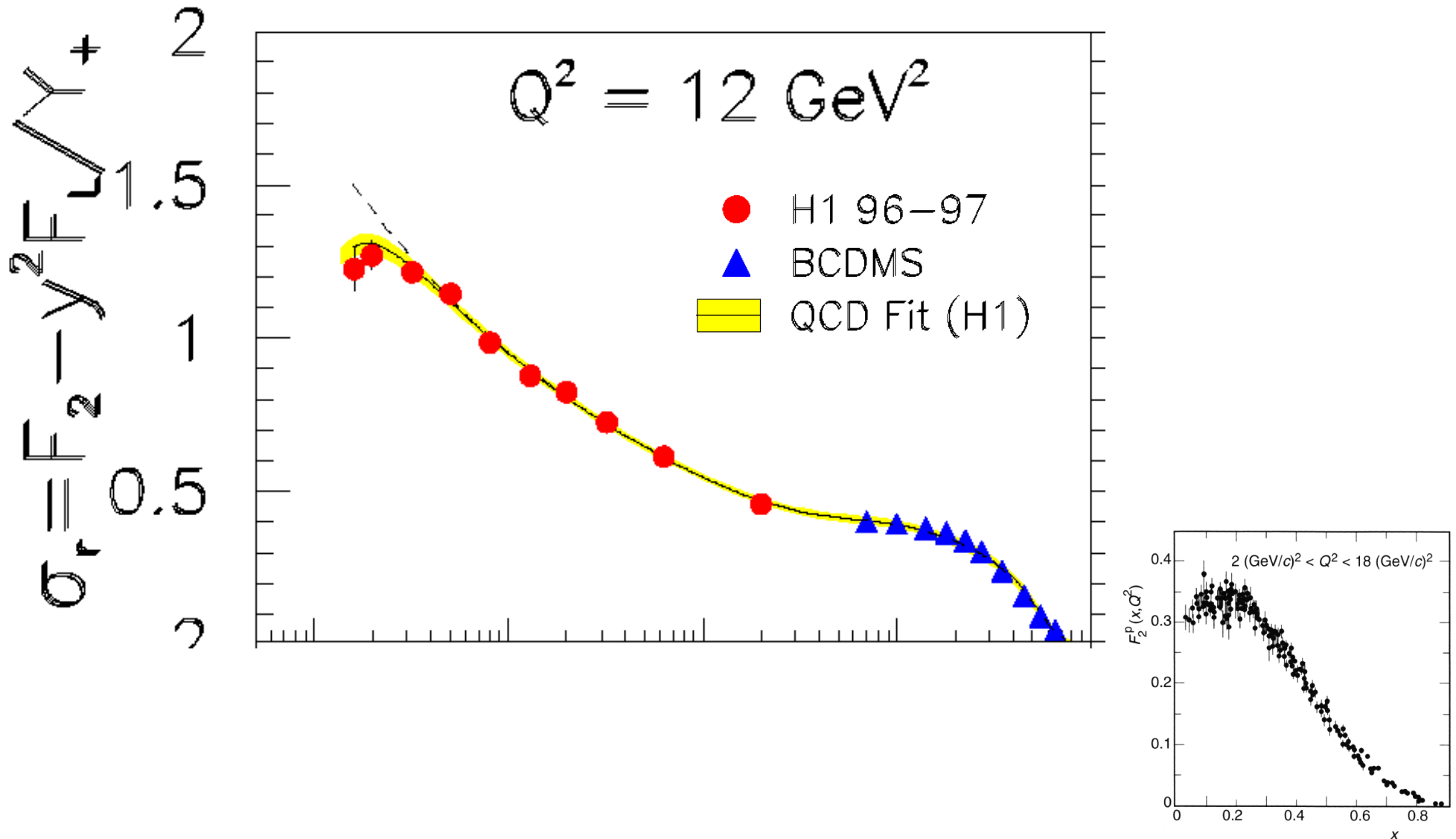
H1 and ZEUS Combined PDF Fit



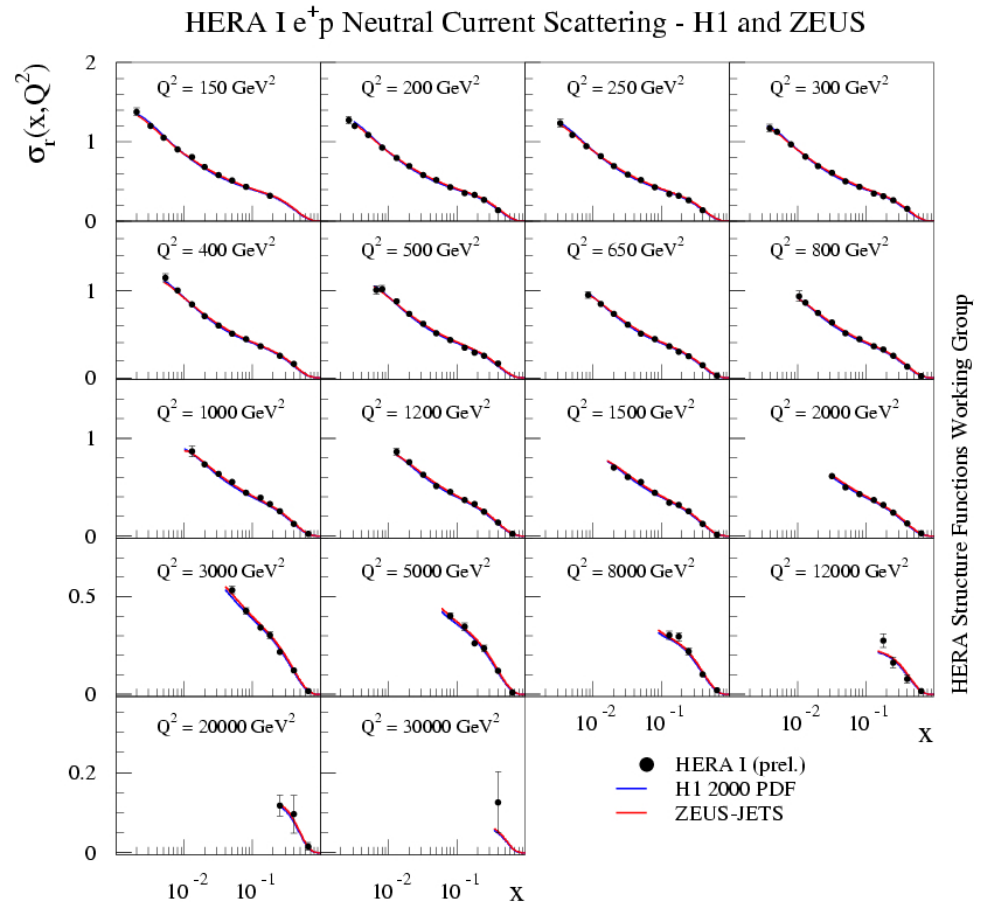
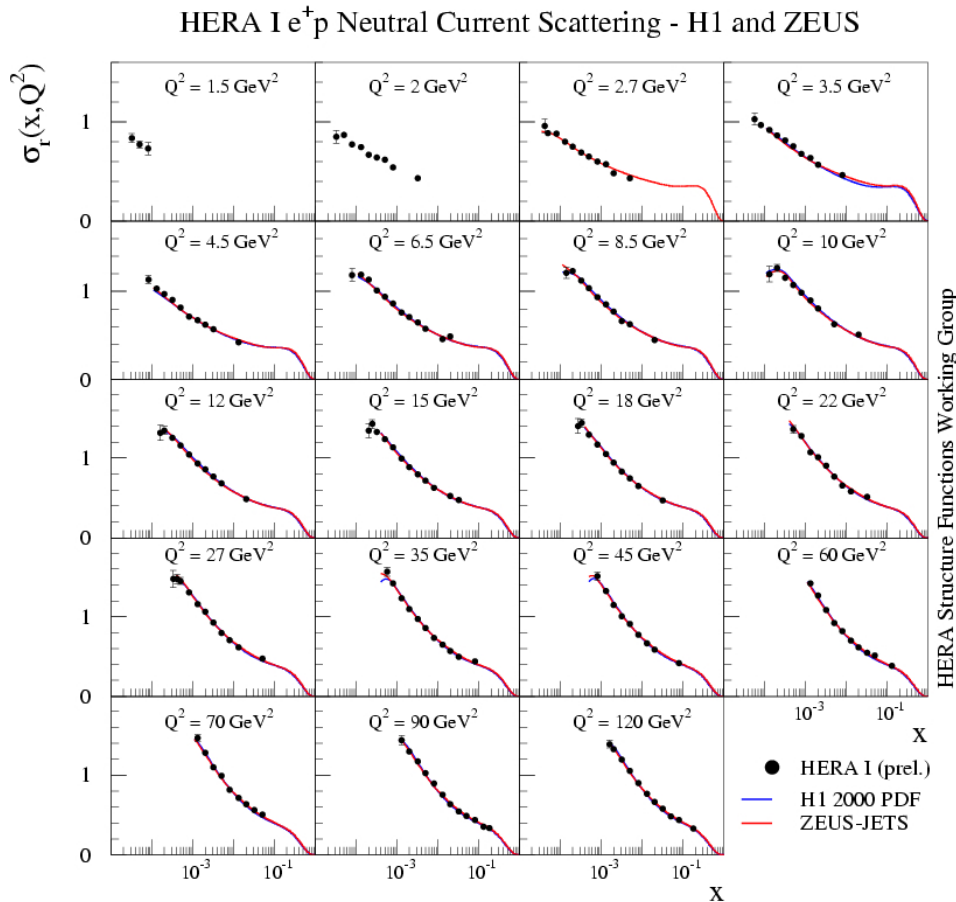
April 2008

HERA Structure Functions Working Group

How does $F_2(x)$ look like?



F_2 vs. x



strong rise towards low x , steepness rising with Q^2

DGLAP Evolution Equations

$$\frac{\partial}{\partial \log Q^2} \begin{bmatrix} q(x, Q^2) \\ g(x, Q^2) \end{bmatrix} = \frac{\alpha_s}{2\pi} \begin{bmatrix} \mathcal{P}_{q/q} \left[\begin{array}{c} \gamma \\ \swarrow \downarrow \searrow \\ x \end{array} \right] & \mathcal{P}_{q/g} \left[\begin{array}{c} \gamma \\ \swarrow \downarrow \searrow \\ x \end{array} \right] \\ \mathcal{P}_{g/q} \left[\begin{array}{c} \gamma \\ \swarrow \downarrow \searrow \\ x \end{array} \right] & \mathcal{P}_{g/g} \left[\begin{array}{c} \gamma \\ \swarrow \downarrow \searrow \\ x \end{array} \right] \end{bmatrix} \otimes \begin{bmatrix} q(x, Q^2) \\ g(x, Q^2) \end{bmatrix}$$

$$\mathcal{P} \otimes f(x, Q^2) = \int_x^1 \frac{dy}{y} \mathcal{P}(x/y) f(y, Q^2)$$

- Q^2 dependence of quark densities $q(x, Q^2)$ and gluon density $g(x, Q^2)$ is predicted

Parton Density Fits

DGLAP predicts only Q^2 dependence

→ assume parametrisation of the parton density functions (PDFs) as a function of x at a starting scale Q_0^2 (typically around 4 - 7 GeV^2):

$$x q(x, Q_0^2) = A x^B (1-x)^C [1 + D x + E x^2 + F x^3]$$

→ evolve the PDFs to all measured Q^2 , calculate F_2 , and fit the parameters to match the data

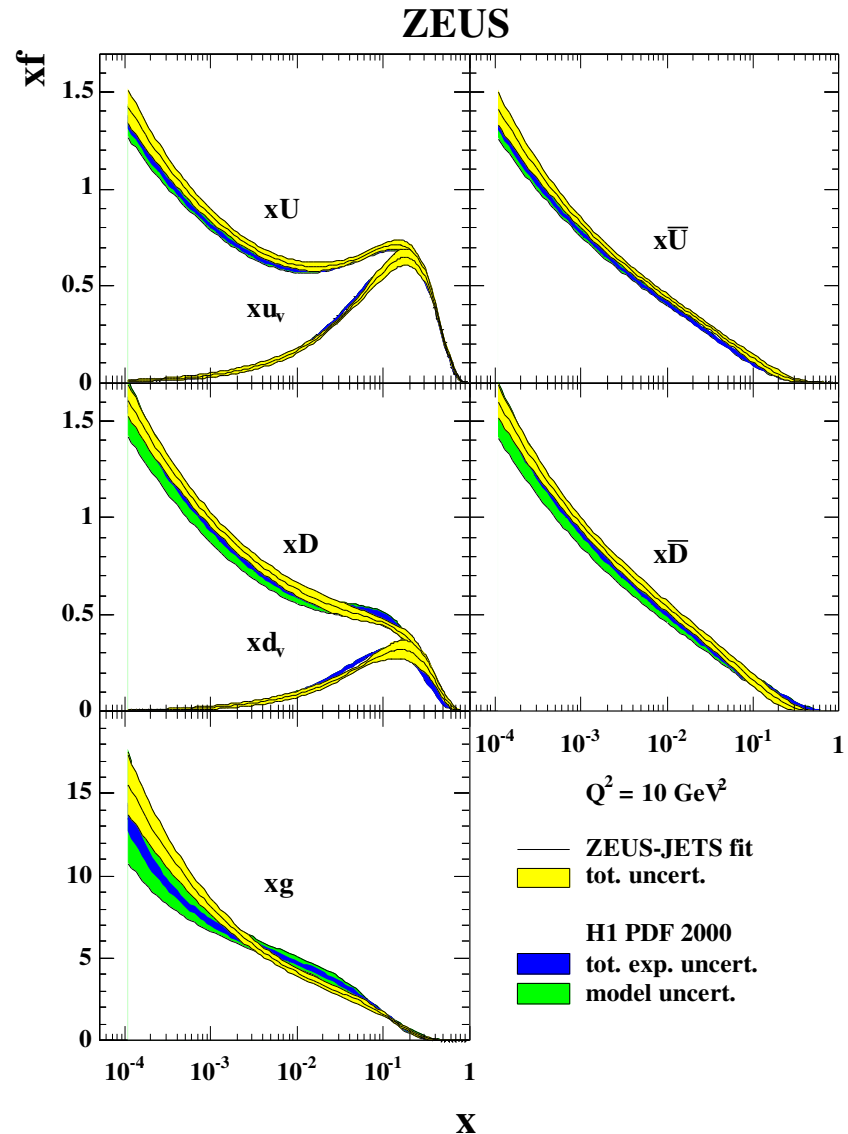
● some freedom in the procedure!

- how many parameters, which Q_0^2 ?
- how to combine quark and antiquark densities?

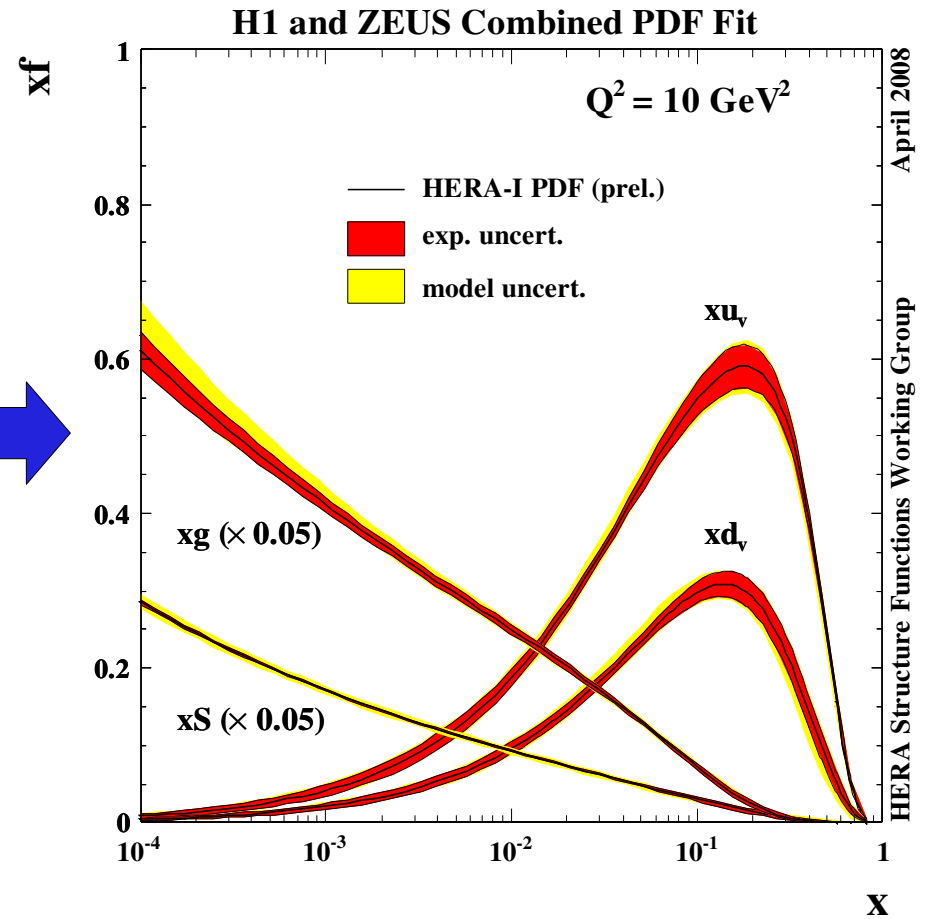
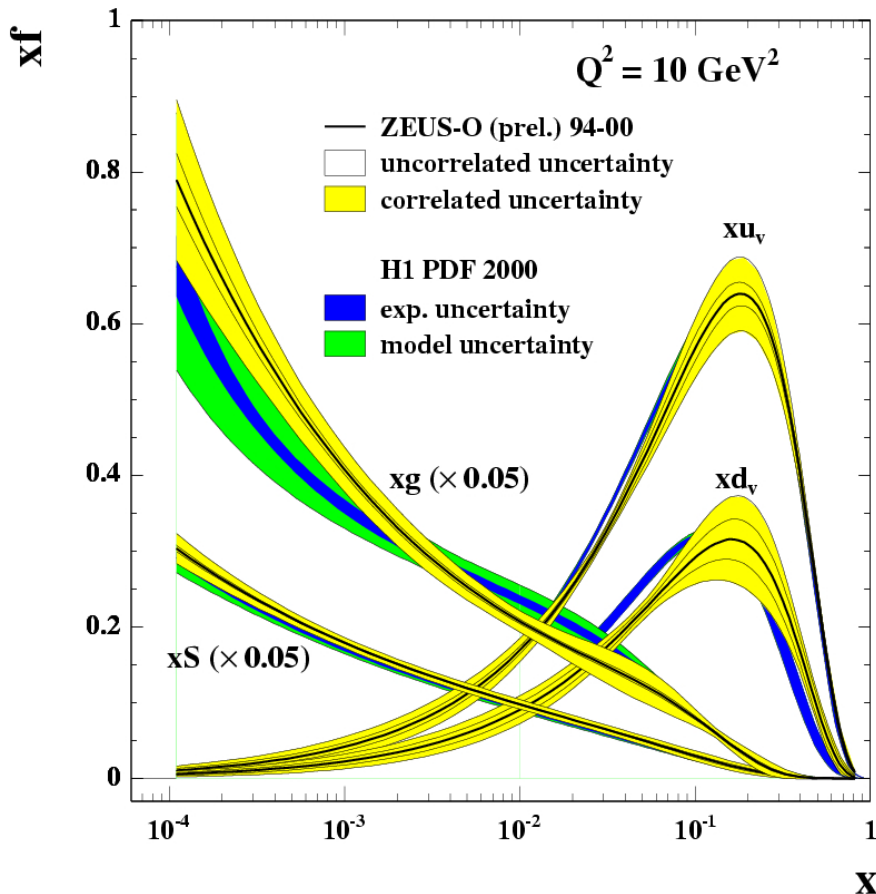
Parton Density Fits

quark and antiquark densities:

- most general: $u, \bar{u}, d, \bar{d}, s, \bar{s}, c, \bar{c}, (b, \bar{b})$
- distinguish valence and sea quarks (ZEUS):
 $u_v, d_v, sea, \bar{d} - \bar{u}$
- distinguish *up*-type and *down*-type quarks (H1):
 $U = u + c, D = d + s (+b)$
 $\bar{U} = \bar{u} + \bar{c}, \bar{D} = \bar{d} + \bar{s} (+\bar{b})$
 $\rightarrow u_v = U - \bar{U}, d_v = D - \bar{D}$



Combined H1 & ZEUS Parton Density



combination of data from H1 and ZEUS
gives big improvements!

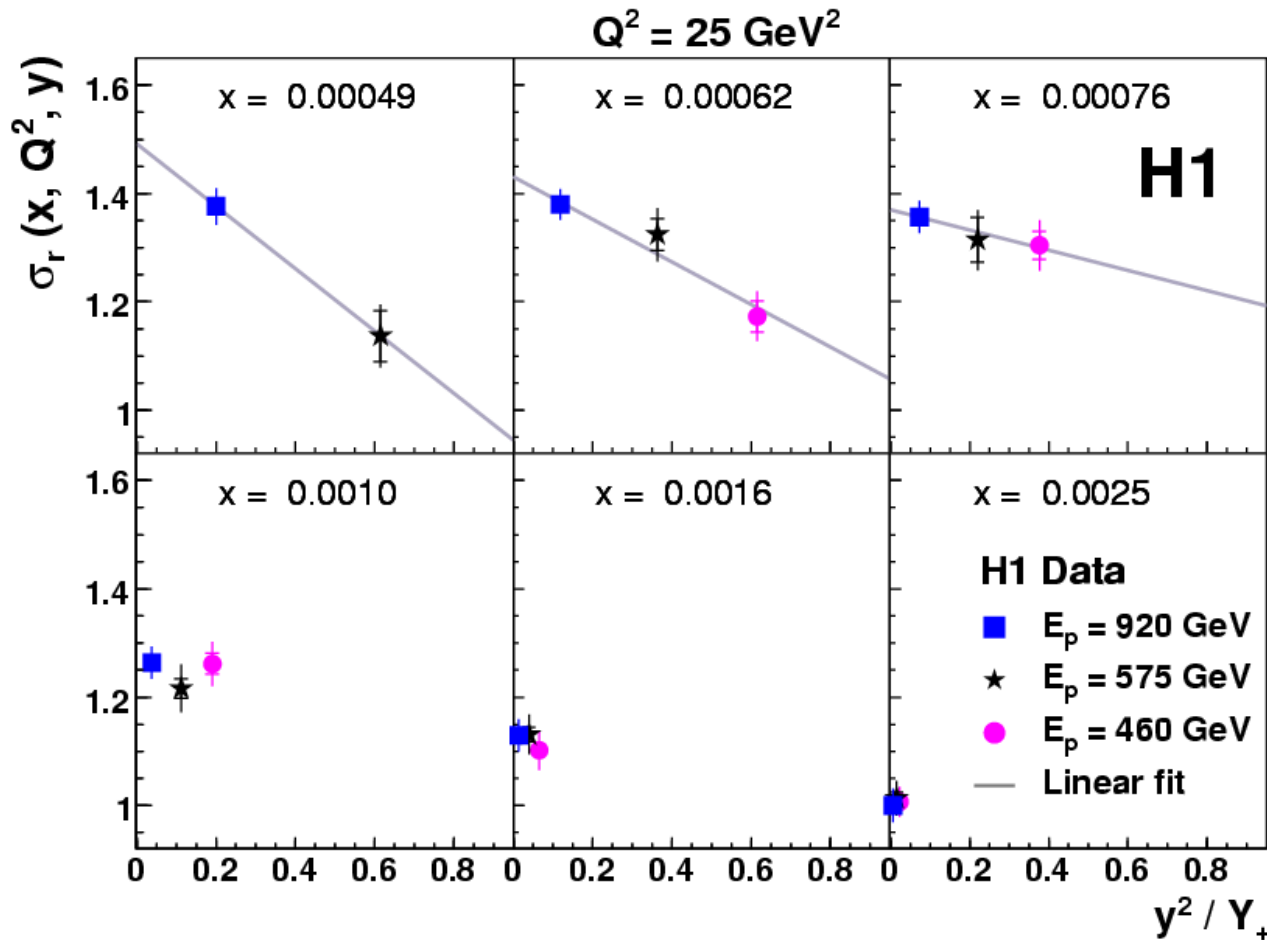
Longitudinal Structure Function F_L

- Callan-Gross relation $2 \times F_1 = F_2$ only true in naive Quark-Parton-Model
- the longitudinal structure function F_L is defined as $F_L = F_2 - 2 \times F_1$
- F_L is directly proportional to the gluon density
- for a measurement of F_L one needs data at the same x and Q^2 , but different y

$$\frac{d^2 \sigma}{dx dQ^2} = \frac{4 \pi \alpha^2}{Q^4} \frac{1}{x} \left(1 - y + \frac{y^2}{2}\right) \left[F_2(x, Q^2) - \frac{y^2/2}{1 - y + y^2/2} F_L(x, Q^2) \right]$$

- only possible with different s because $Q^2 = xys$
- measure at different beam energies!

• Longitudinal Structure Function F_L



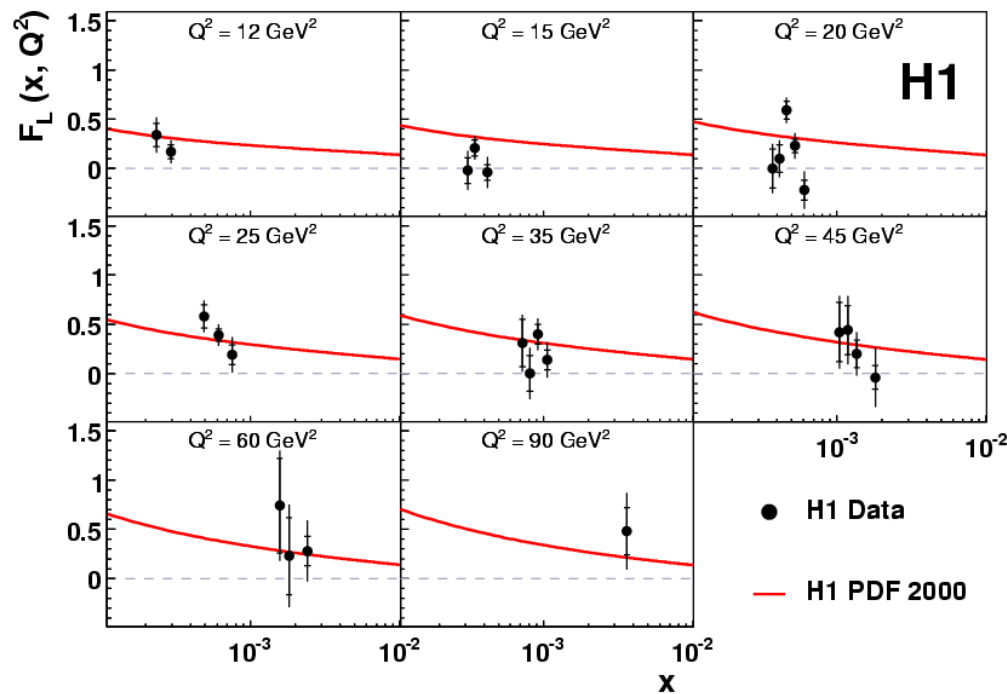
$$\sigma_r = \frac{x Q^4}{2\pi\alpha^2} \frac{1}{Y_+} \frac{d^2\sigma}{dx dQ^2}$$

$$= F_2(x, Q^2) - \frac{y^2}{Y_+} F_L(x, Q^2)$$

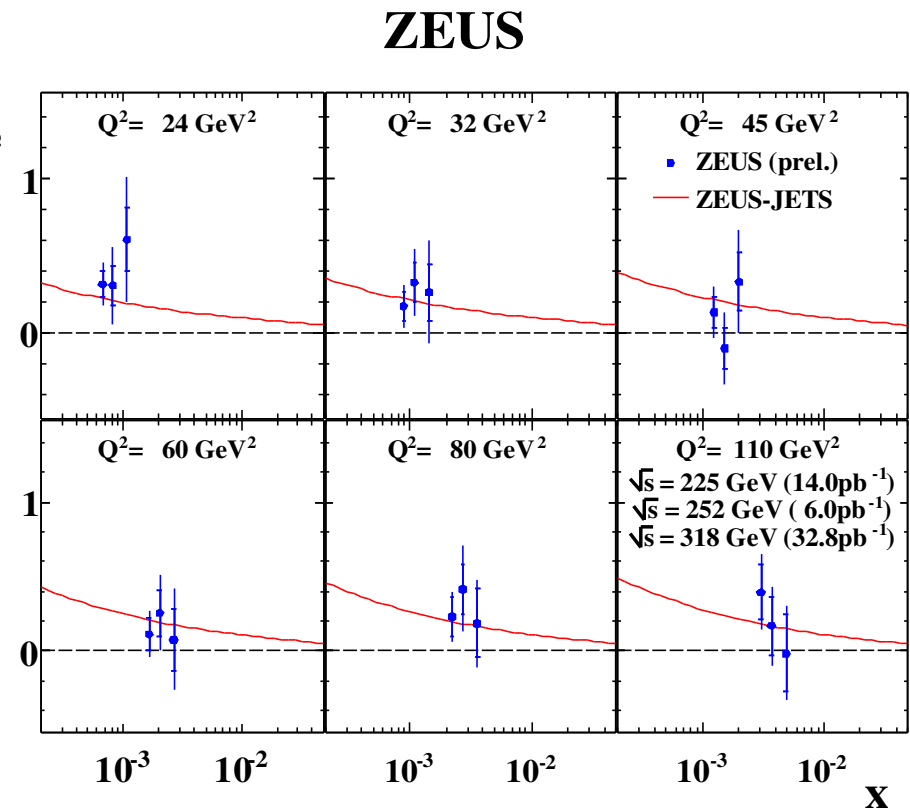
with $Y_+ = 1 + (1-y)^2$

- linear expression in y^2/Y_+
- use linear fits in y^2/Y_+ and determine F_L from slope

Longitudinal Structure Function F_L



- consistent with PDF fit to F_2
- most precise information on gluon still from scaling violations



High Q^2 & Electroweak Physics

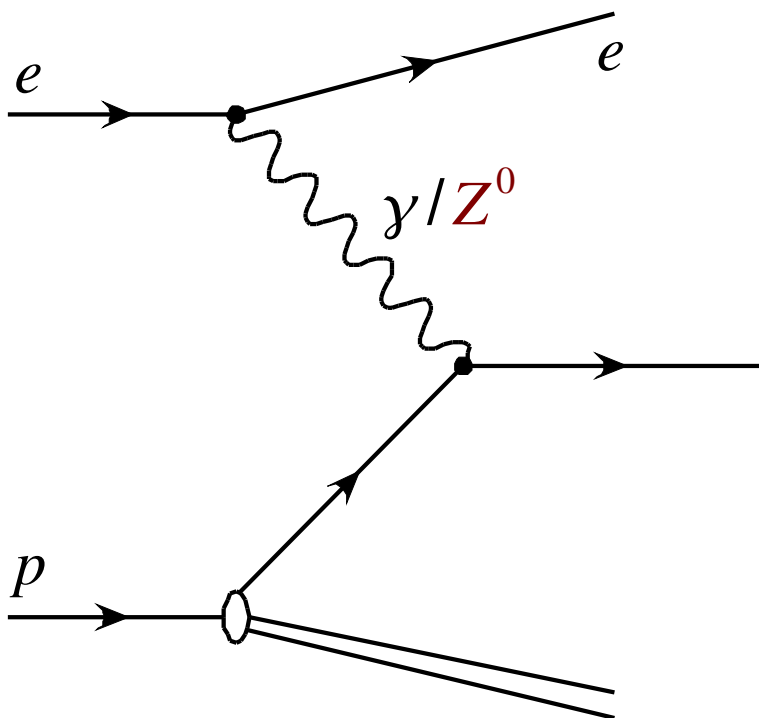
More Structure Functions

$$F_L = F_2 - 2xF_1 = 0 \text{ in the QPM}$$

$$\frac{d^2 \sigma_{NC}^{\pm}}{dx dQ^2} = \frac{2\pi\alpha^2}{Q^4} \frac{1}{x} Y_{\pm} \left[F_2(x, Q^2) - \frac{y^2}{Y_{\pm}} F_L(x, Q^2) \mp \frac{Y_{\mp}}{Y_{\pm}} x F_3(x, Q^2) \right]$$

F_3 : γ - Z^0 -interference

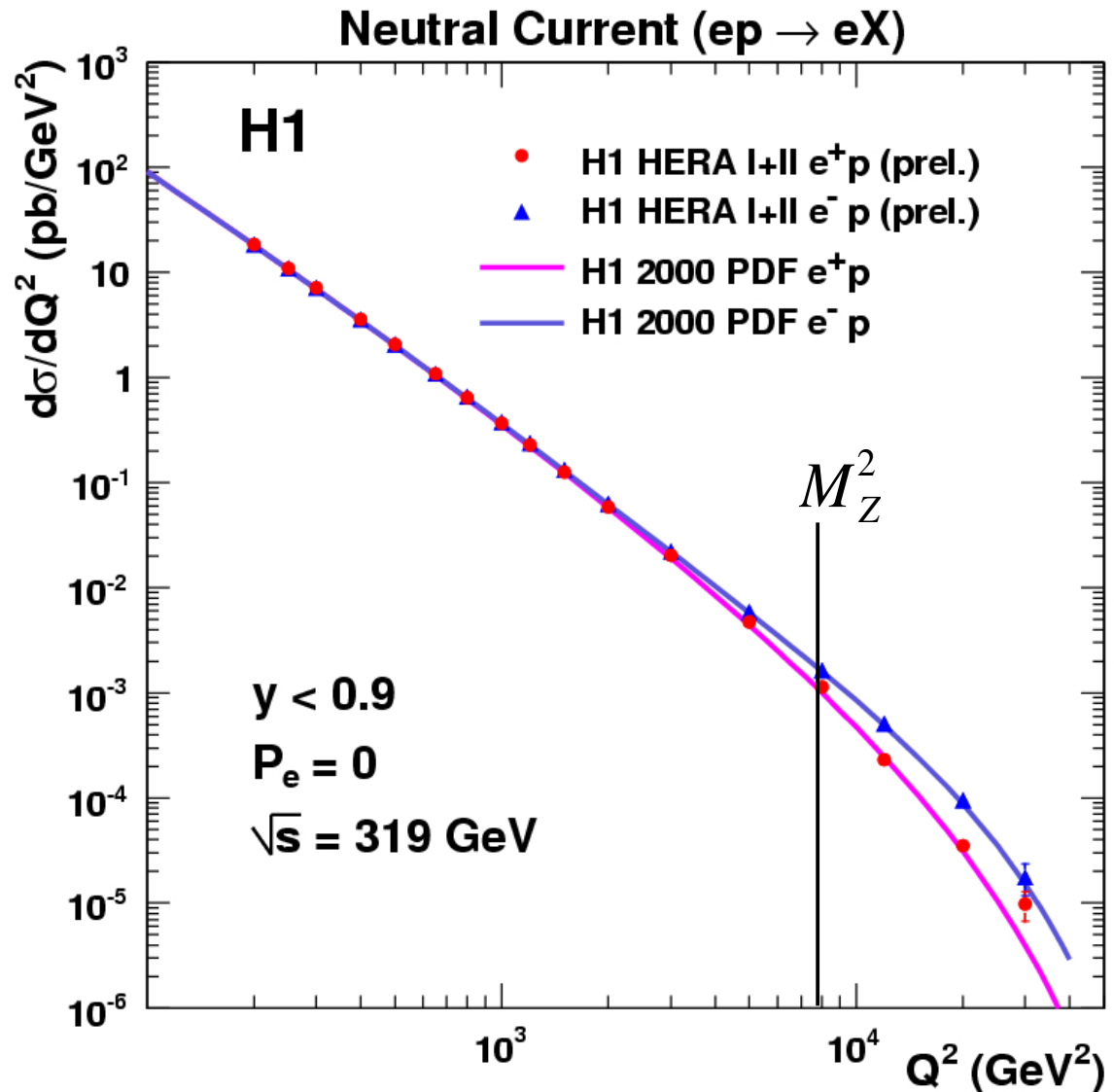
$$Y_{\pm} = 1 \pm (1-y)^2$$



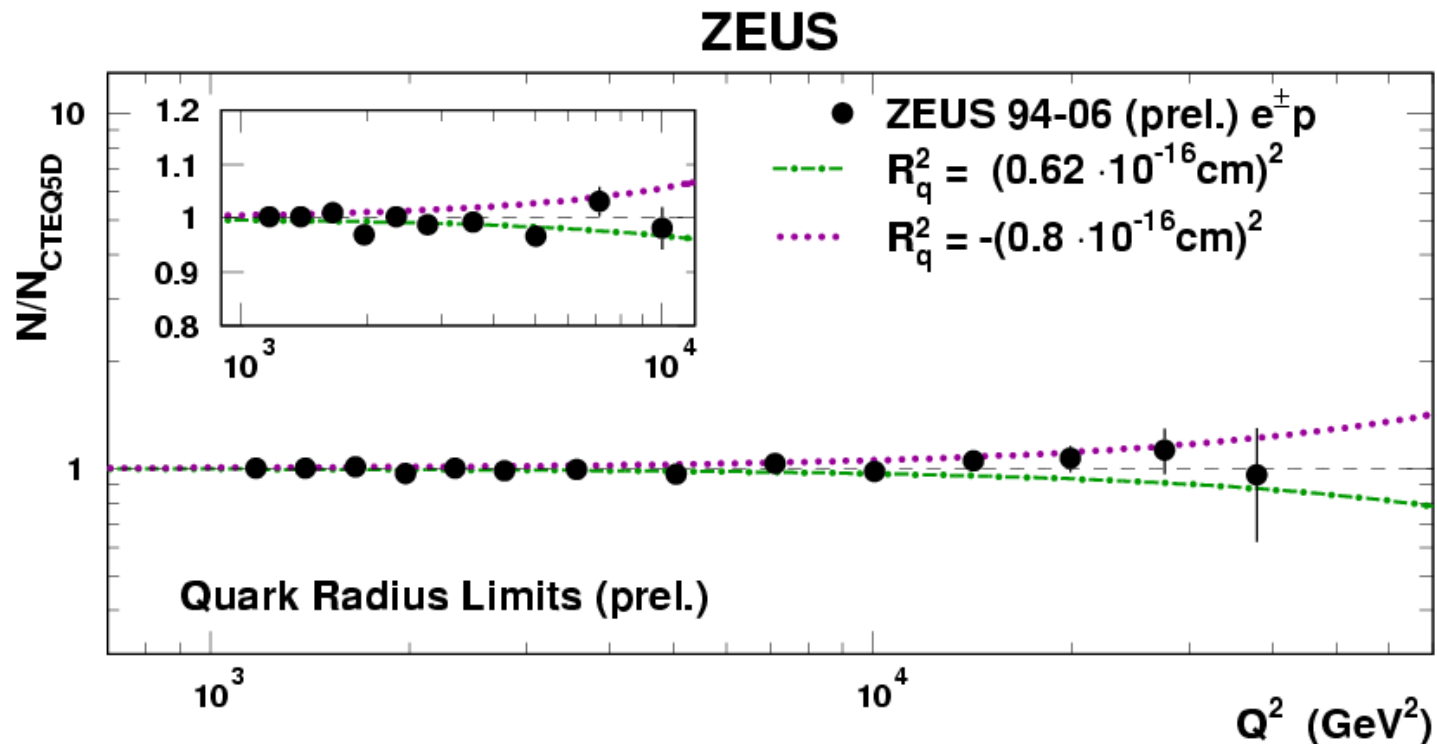
- F_L relevant only at large y
- F_3 relevant only at large Q^2 ,
different sign for e^+ and e^-

High Q^2 Neutral Current

- difference between e^+p and e^-p only at large $Q^2 \approx M_Z^2$
→ $\gamma - Z^0$ interference



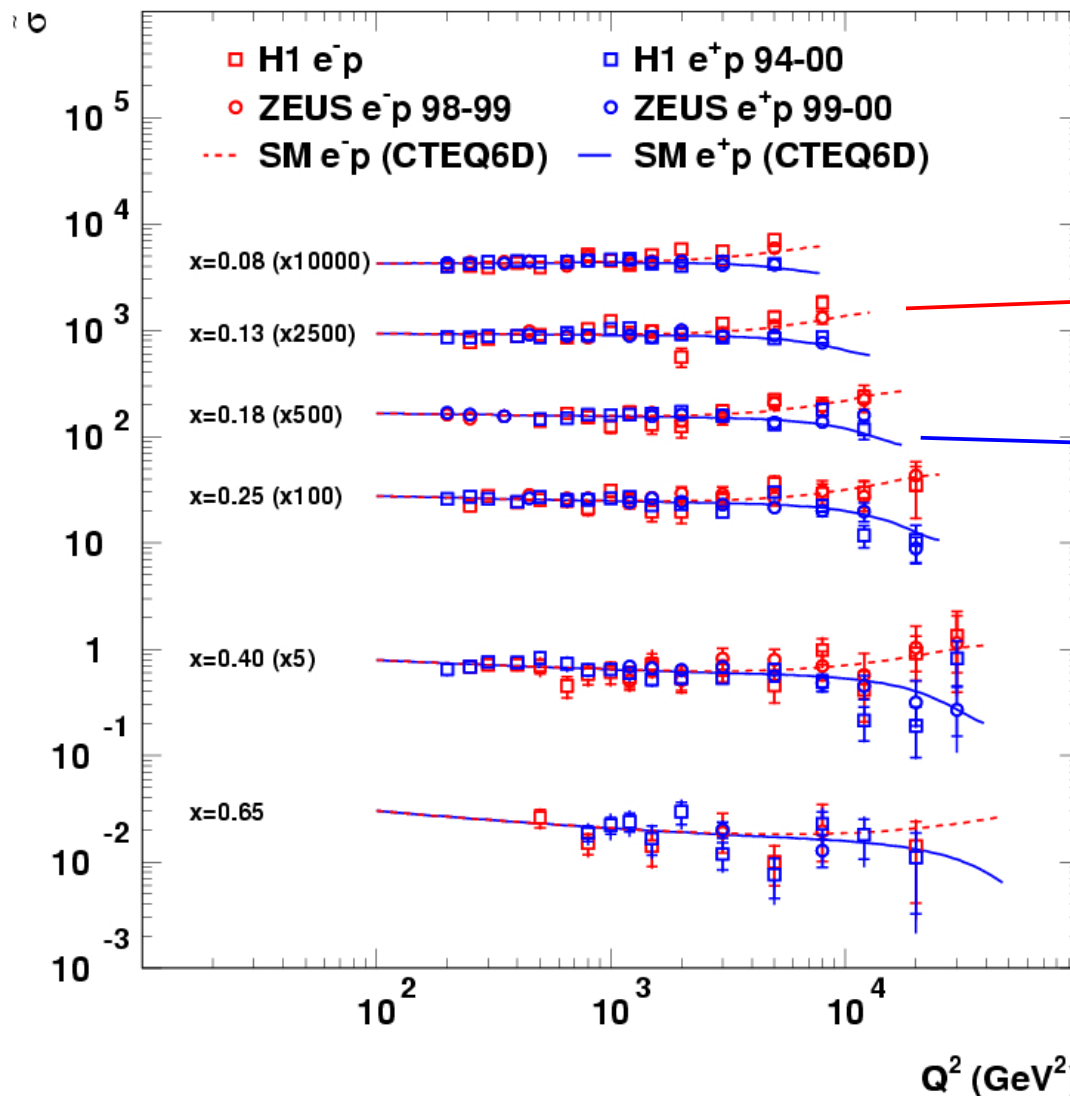
High Q^2 Neutral Current



- no significant deviation from Standard Model Fit at high Q^2
- can be interpreted as limit on quark size

High Q^2 Neutral Current

HERA Neutral Current at high x



$$\tilde{\sigma} = \frac{x Q^4}{2 \pi \alpha^2} \frac{1}{Y_+} \frac{d^2 \sigma_{NC}^{\pm}}{dx dQ^2}$$

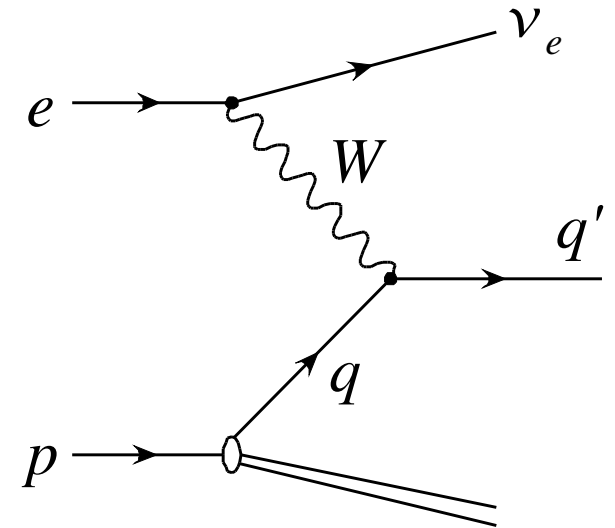
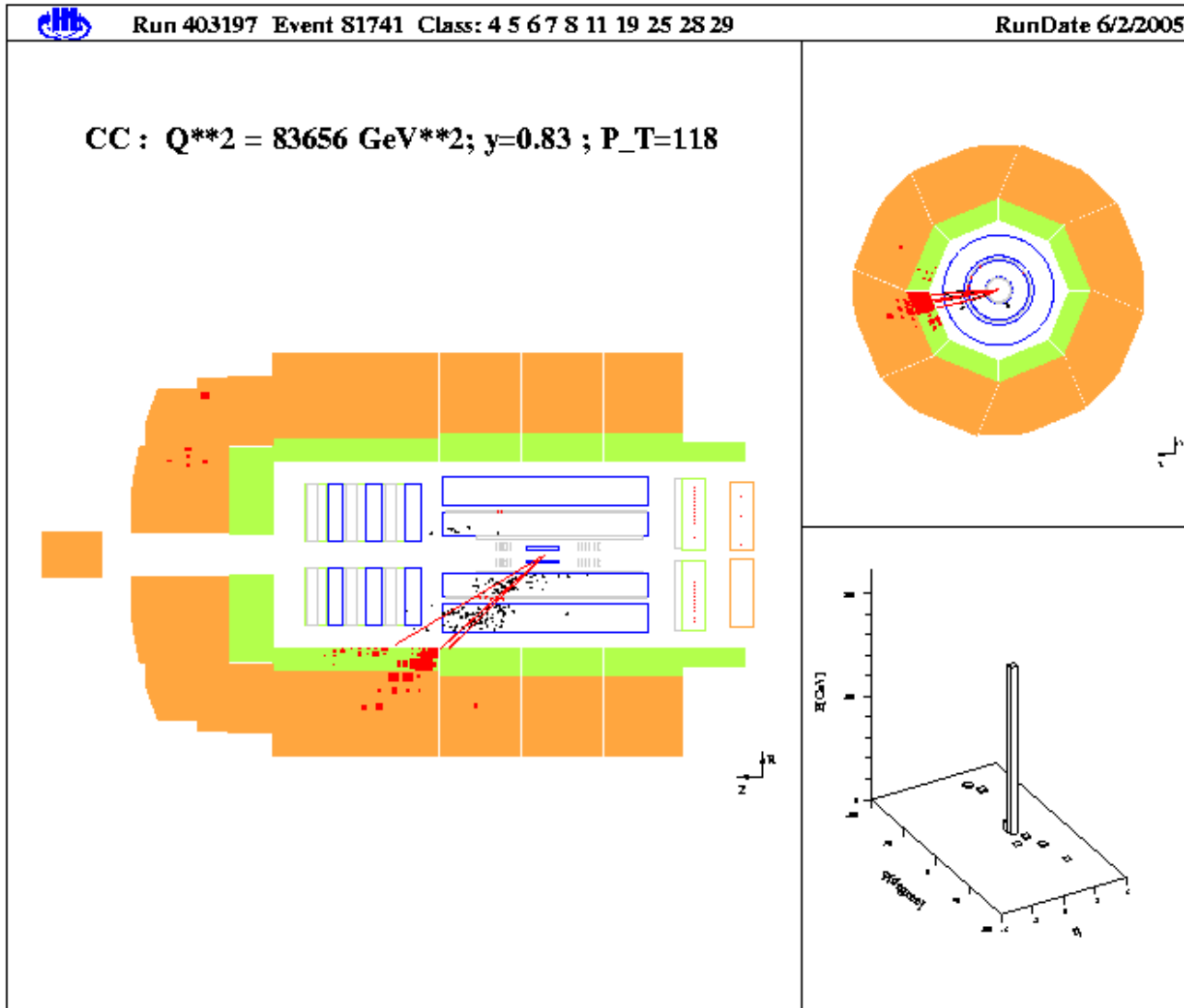
e^- positive interference

e^+ negative interference

$$x F_3 \propto x \sum e_q^2 (q - \bar{q})$$

direct handle on
valence quark
distribution!

Charged Current Interactions



neutrino not visible
in detector

→ imbalance in
transverse plane

Charged Current Cross Section

$$\frac{d^2 \sigma_{CC}^{\pm}}{dx dQ^2} = \frac{G_F^2}{4\pi x} \left(\frac{M_W^2}{M_W^2 + Q^2} \right)^2 Y_{\pm} \left[W_2^{\pm} - \frac{y^2}{Y_{\pm}} W_L^{\pm} \mp \frac{Y_{\mp}}{Y_{\pm}} x W_3^{\pm} \right]$$

- W bosons couple differently to *up*- and *down*-type quarks

- in the QPM:

$$W_2^- = x(U + \bar{D}), \quad x W_3^- = x(U - \bar{D})$$

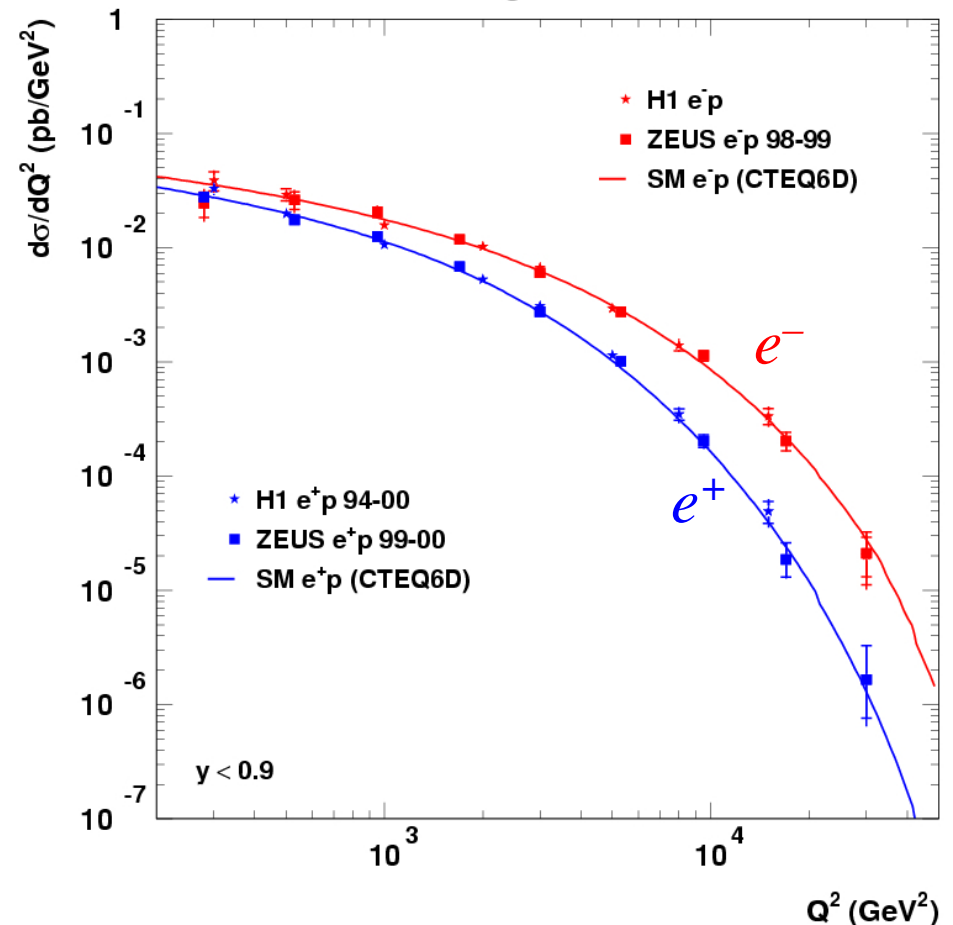
$$W_2^+ = x(\bar{U} + D), \quad x W_3^+ = x(D - \bar{U})$$

$$W_L^{\pm} = 0$$

$$\rightarrow \sigma_{CC}^- \propto x \left[U + (1-y)^2 \bar{D} \right]$$

$$\sigma_{CC}^+ \propto x \left[\bar{U} + (1-y)^2 D \right]$$

HERA Charged Current



Comparison NC vs. CC

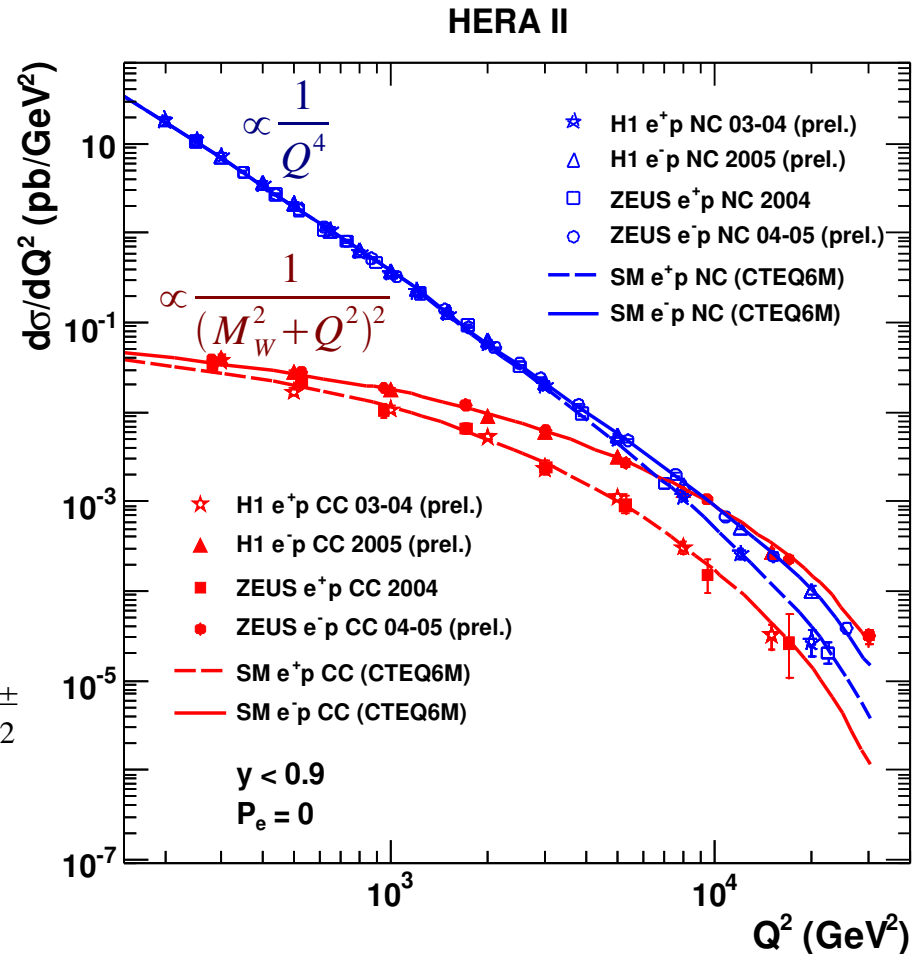
- at low Q^2 : different dependences because of photon in NC
- at high Q^2 : „electroweak unification“

but:

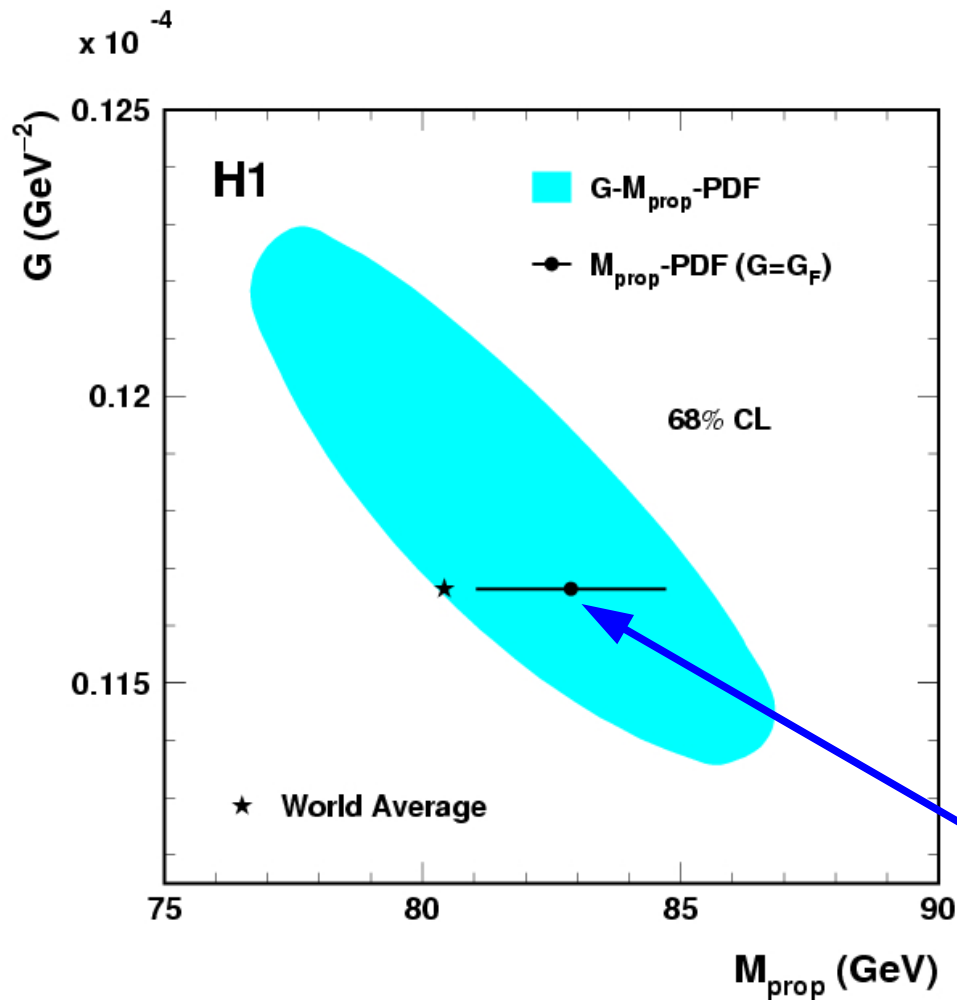
$$\frac{d^2 \sigma_{CC}^{\pm}}{dx dQ^2} \approx \frac{G_F^2}{4\pi x} \cdot \left(\frac{M_W^2}{M_W^2 + Q^2} \right)^2 \cdot Y_{\pm} W_2^{\pm}$$

$$\frac{d^2 \sigma_{NC}^{\pm}}{dx dQ^2} \approx \frac{2\pi\alpha^2}{x} \cdot \frac{1}{Q^4} \cdot Y_{\pm} F_2$$

similar because $G_F \approx \frac{4\pi\alpha}{\sqrt{2}M_W^2}$



Electroweak Parameters: W Mass



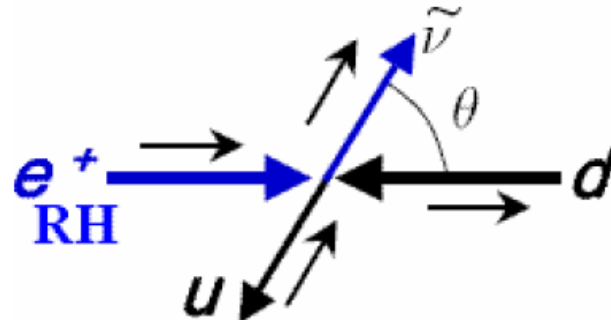
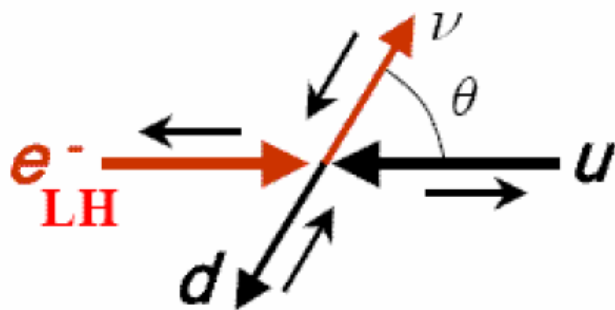
- $G = G_F$
determined by normalization of the CC cross section
 - $M_{\text{prop}} = M_W$
determined by the Q^2 dependence of the CC cross section
- $82.87 \pm 1.82_{\text{exp}} \left(\begin{smallmatrix} +0.30 \\ -0.16 \end{smallmatrix} \right)_{\text{model}} \text{ GeV}$

CC & Polarization

- CC cross section depends on longitudinal electron/positron polarization P_e

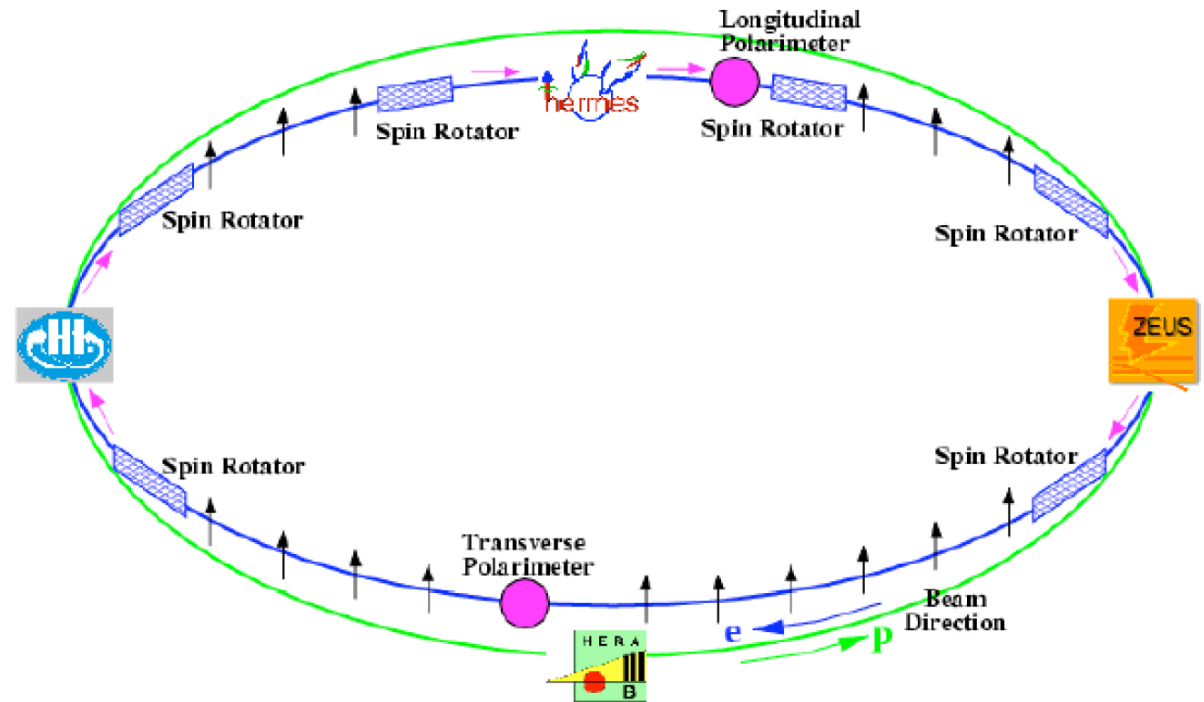
$$\frac{d^2 \sigma_{CC}^{\pm}}{dx dQ^2}(P_e) \approx (1 \pm P_e) \frac{G_F^2}{4 \pi x} \cdot \left(\frac{M_W^2}{M_W^2 + Q^2} \right)^2 \cdot Y_{\pm} W_{\frac{1}{2}}^{\pm}$$

- reason: W boson couples only to left-handed (LH) particles and right-handed (RH) antiparticles:



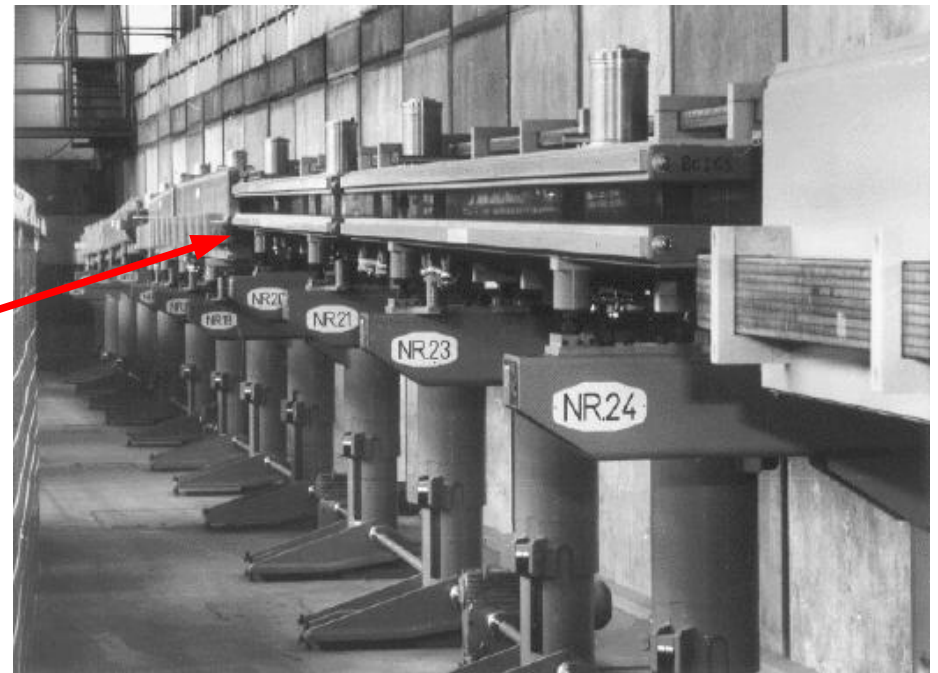
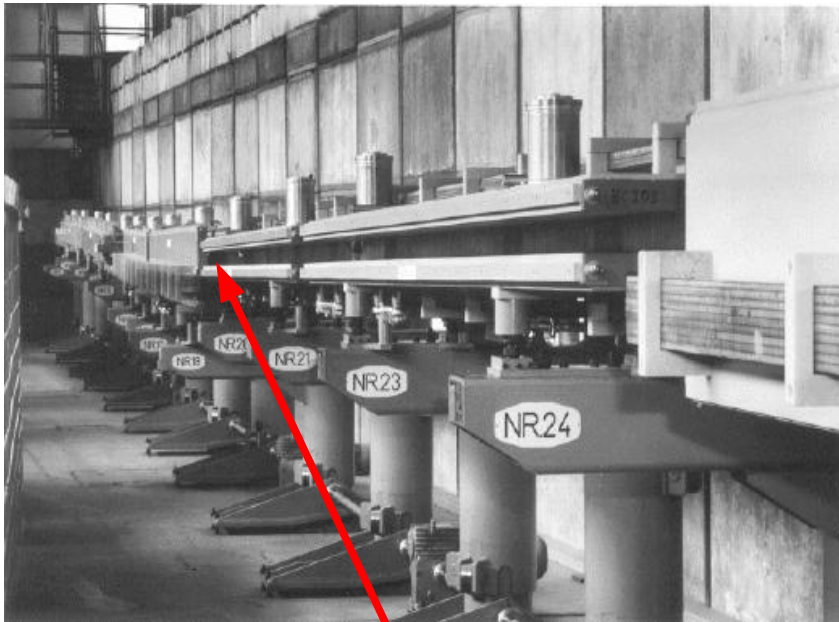
Polarization @ HERA

$$P_e = \frac{N_{RH} - N_{LH}}{N_{RH} + N_{LH}}$$



- transverse polarization builds up in ~40 minutes through synchrotron radiation (Sokolov-Ternov effect)
- spin rotators flip transverse \rightarrow longitudinal before experiments and back after

Polarization @ HERA



spin rotator

CC: Polarization Dependence

- Standard Modell expectation:

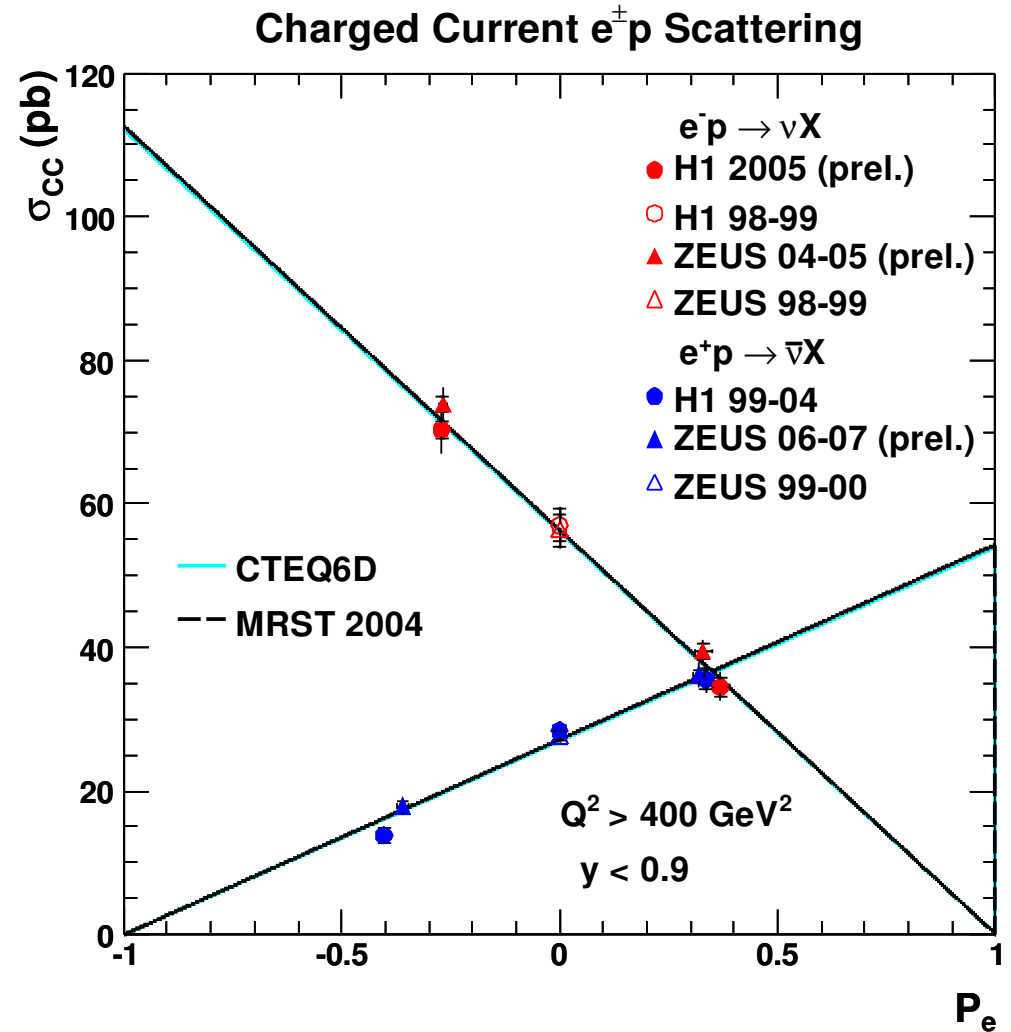
$$\sigma_{CC}^{-}(P_e=+1) = 0$$

$$\sigma_{CC}^{+}(P_e=-1) = 0$$

- experimental result: (H1)

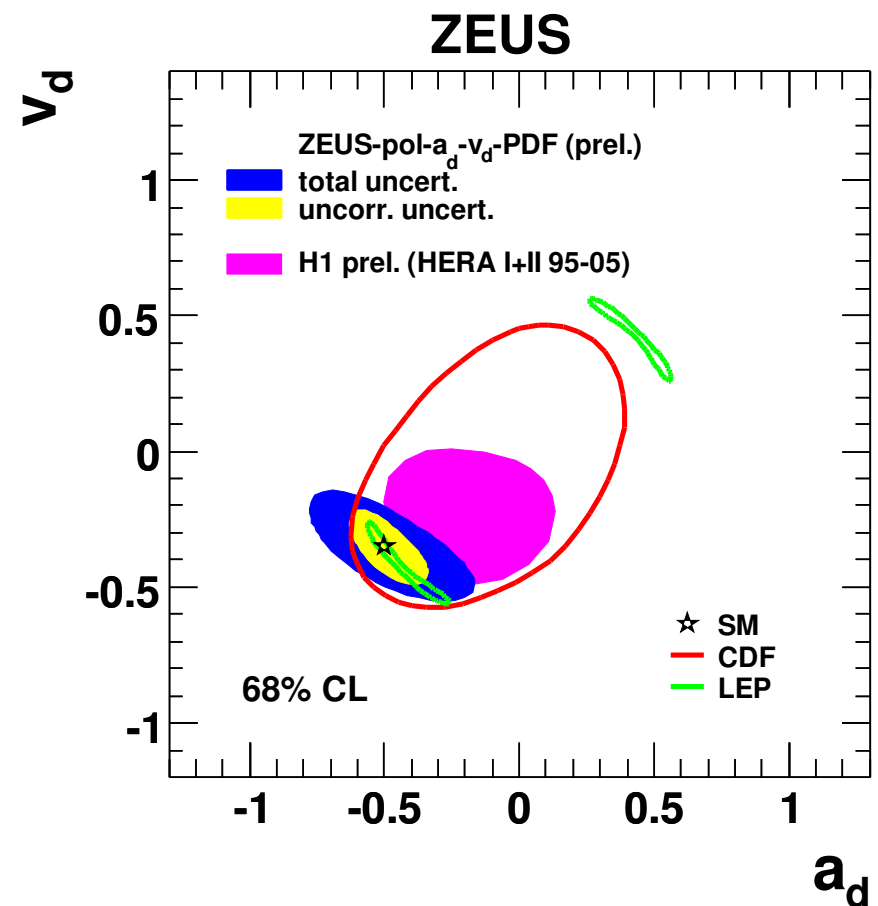
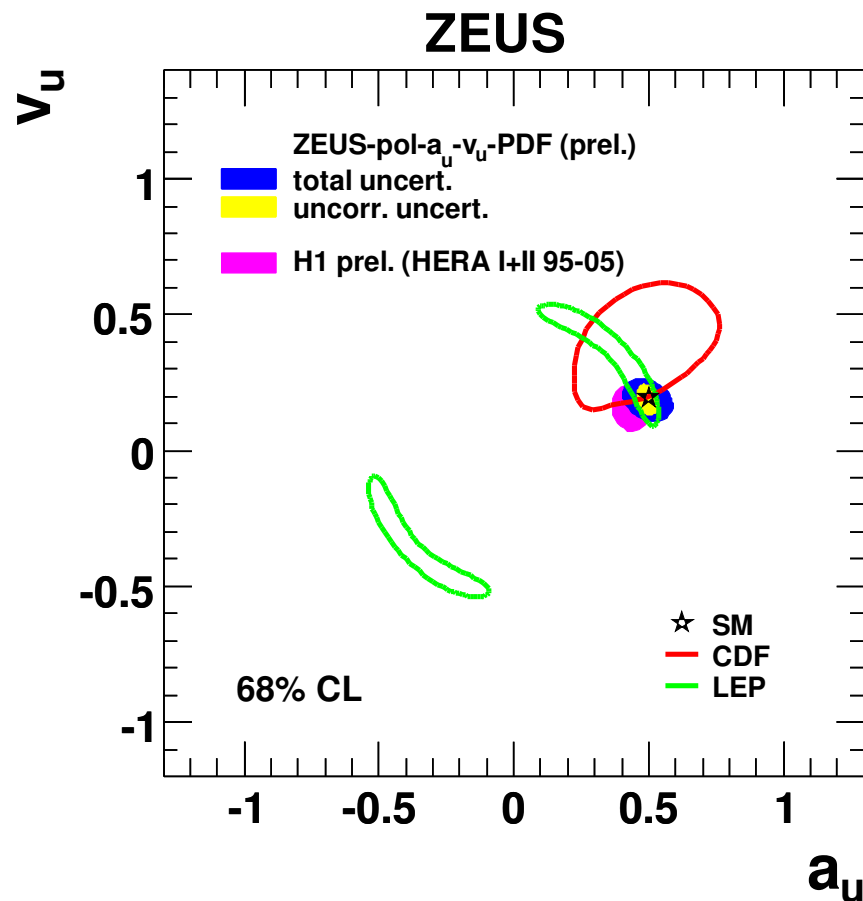
$$\sigma_{CC}^{-}(+1) = -0.9 \pm 2.9_{stat} \pm 1.9_{syst} \pm 1.9_{pol} \text{ pb}$$

$$\sigma_{CC}^{+}(-1) = -3.9 \pm 2.3_{stat} \pm 0.7_{syst} \pm 0.8_{pol} \text{ pb}$$



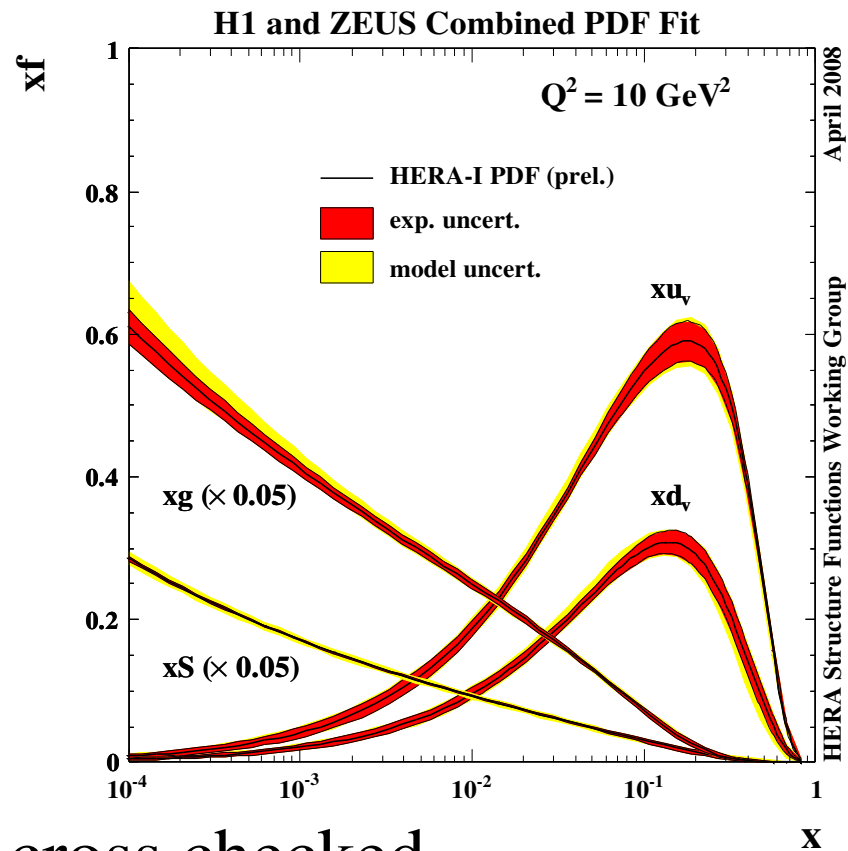
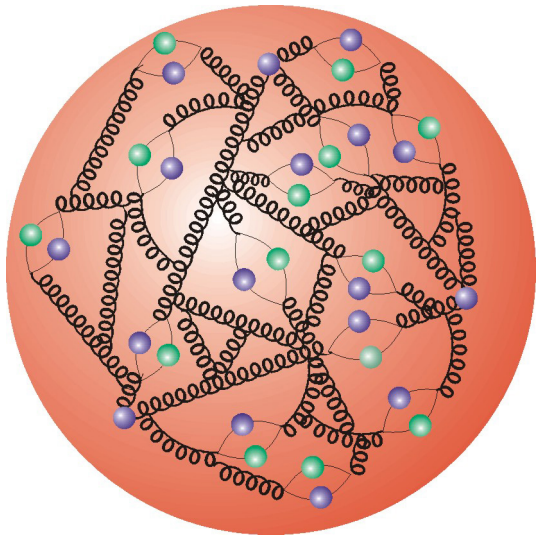
Electroweak Parameters: Z^0 Couplings

polarization also allows better sensitivity to vector and axial-vector couplings of up - and $down$ -type quarks to the Z^0



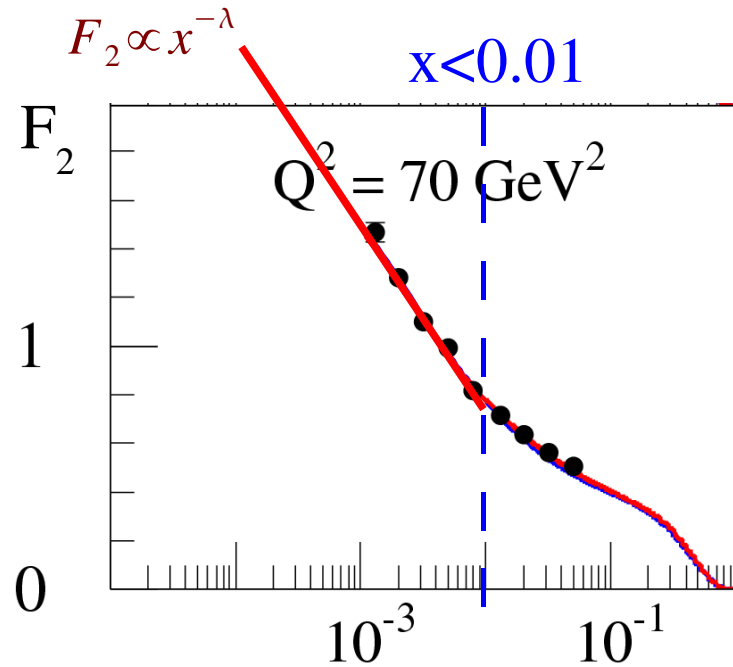
Summary

- inclusive ep scattering reveals structure of the proton
- large amount of gluons in the proton

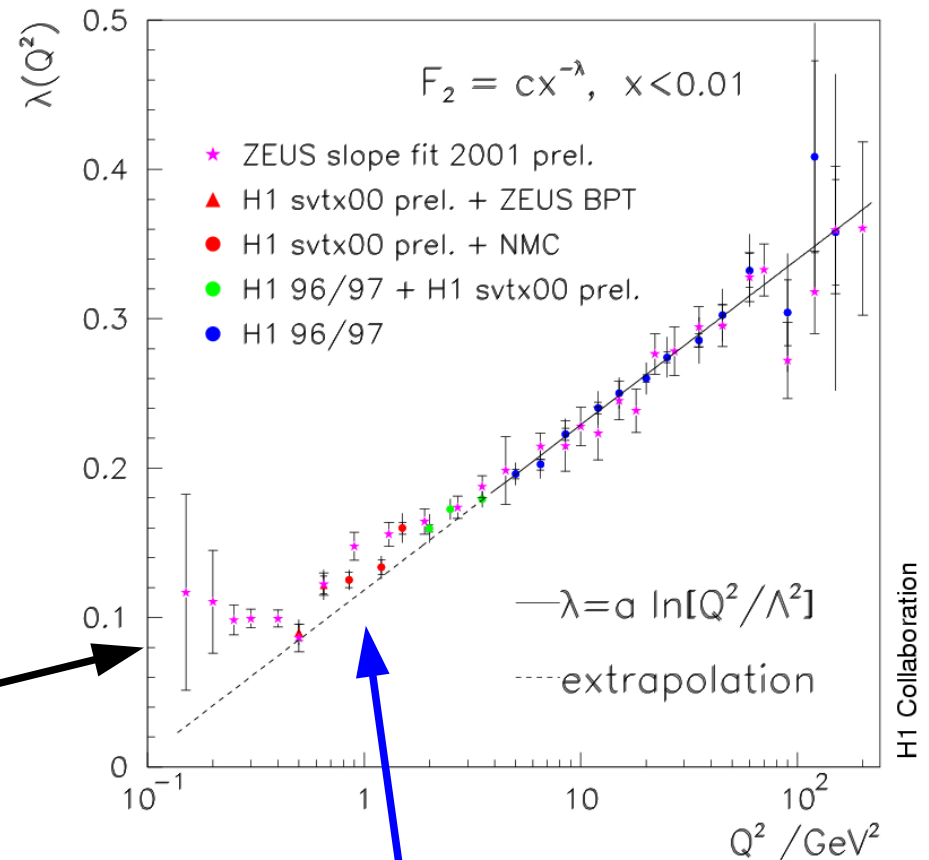


- Standard Modell can be cross checked

Hadrons vs. Partons



0.08 corresponds to
hadron-hadron scattering



transition from hadronic to partonic behaviour