

QCD description of multi-particle production

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6.2.2001
DESY

1. Theoretical approach

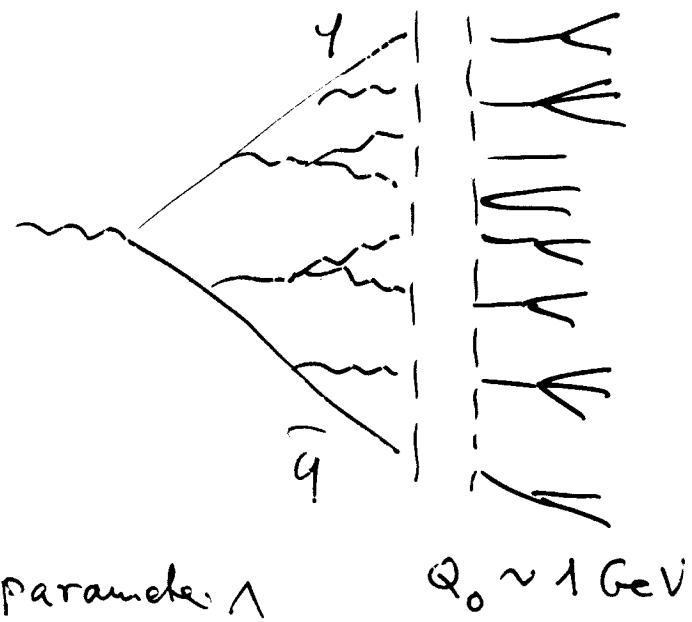
2. Results

- multiplicity : mean and moments
- inclusive distribution
→ soft limit
- multi-jet events
- angular correlation
- rapidity gaps
- outlook

3. Conclusions

Theoretical approach

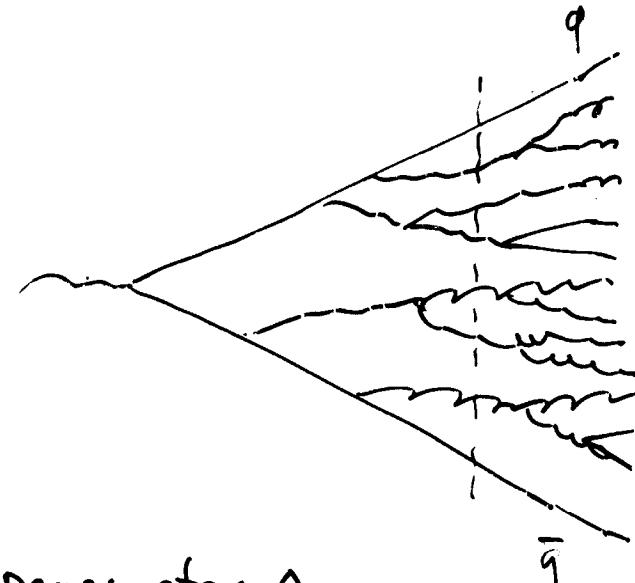
Hadronization module



$$Q_0 \sim 1 \text{ GeV}$$

Hadrons
Resonances

"Parton Hadron Duality"



parameter Λ

cascade evolves towards

is QCD relevant here?

study :

- $\alpha_s \gtrsim 1$
- coherence effects of soft gluons
- color factors C_F, C_C

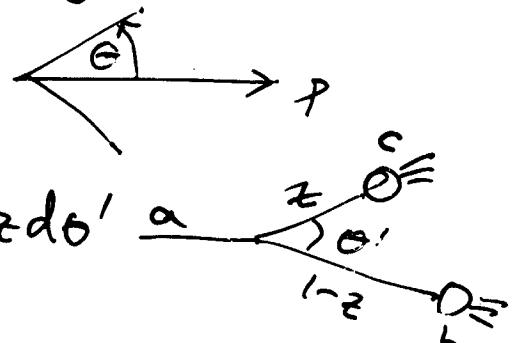
$K_T > Q_0 \gtrsim \Lambda$ (few 100 MeV)

Theoretical approach

Parton cascade

Master equation for generating functional

$$Z(P, \theta, \{u(p_i)\})$$



$$\frac{a}{\partial \theta} = \frac{a}{e^{-w_a(P, \theta)}} u_a + \sum_{a \rightarrow b c} \int dz d\theta' \frac{\alpha_s(k_\perp)}{2\pi} P_a^{bc}(z) Z(\theta, \theta') Z((1-z)p, \theta')$$

- $P_{abc}(z)$ DGLAP splitting function
- angular ordering $\theta' < \theta$ (coherence)
- initial condition: absolute normalization
- running coupling $\alpha_s(k_\perp)$ at 1-loop
- non-perturbative cut-off $K_\perp > Q_0 > \Lambda_{QCD}$
- approximations in leading log: $\frac{d\theta}{\theta} \sim \frac{d\ln Q}{\ln Q}$

LC (PLA): $\alpha_s^n \ln^2 y \rightarrow e^{Y \sqrt{\alpha_s}} \sim e^{\sqrt{\ln Q}}$

NLO (MLLA): $" + \alpha_s^n \ln^{2n-1} y \rightarrow e^{Y \sqrt{\alpha_s} + P \ln \alpha_s}$

numerical solution of eq. \rightarrow yet higher order
but not complete

- match with matrix element $C(x_s), C(x_s^2)$

• numerical method MC:

ARIADNE parton level $K_\perp > Q_0, \alpha_s(K_\perp)$

Parameters: Λ_{QCD}

$Q_0 \sim R'$ non-perturbative
(K normalization)

Hadrons

Parton - Hadron - Duality

$$O(x_1 \dots x_n) |_{\text{hadron}} = K O(x_1 \dots x_n, Q_0, \Lambda) |_{\text{parton}}$$

$$K \approx 1 \dots 2$$

Azimov

Dokshitzer

Khoze

Troyan

for nucle. > 1 GeV

spectra

in MLLA

mismatch near kinematic border

minimal K_T :

$$K_T |_{\text{hadron}} \rightarrow 0$$

$$K_T |_{\text{parton}} \rightarrow Q_0$$

simple model:

a hadron is like a parton with mass Q_0

partons

$$E_p^2 = k_p^2 > Q_0^2$$

hadrons

$$E_h^2 = k_h^2 + Q_0^2 > Q_0^2$$

compare partons and hadrons at $E_p = E_h$

Results:

Mean multiplicities

a) high energy approximation (MLLA)

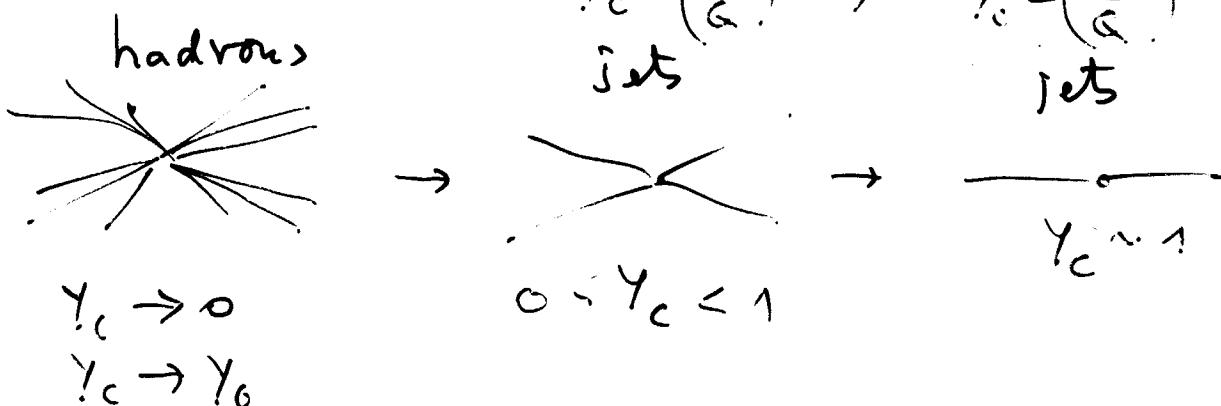
$$N \sim \exp \left(\frac{C_1}{\sqrt{\alpha_s(Y)}} + C_2 \ln \alpha_s(Y) + \dots \right)$$

$$C_1 = \sqrt{96\pi}/b, \quad C_2 = \frac{1}{4} + \frac{10}{27} \frac{n_f}{b}, \quad b = \frac{11}{3} N_c - \frac{2}{3}$$

b) full solution with initial condition $N=1$
 transition jets \rightarrow hadrons

$$\text{Durham} \quad K_T; j > Q_c \quad K_T > Q_0$$

$$Y_c = \left(\frac{Q_c}{Q} \right)^2 \rightarrow \quad Y_c = \left(\frac{Q_c}{Q} \right)^2$$



c) gluon jet multiplicity

$$r = \frac{N_g}{N_q} = \frac{C_A}{C_F} \left(1 - r_1 \gamma_0 - r_2 \gamma_0^2 - r_3 \gamma_0^3 \right)$$

$$\gamma_0 = \left(\frac{2 C_A \alpha_s}{\pi} \right)^{1/2}$$

Multiplicity moments

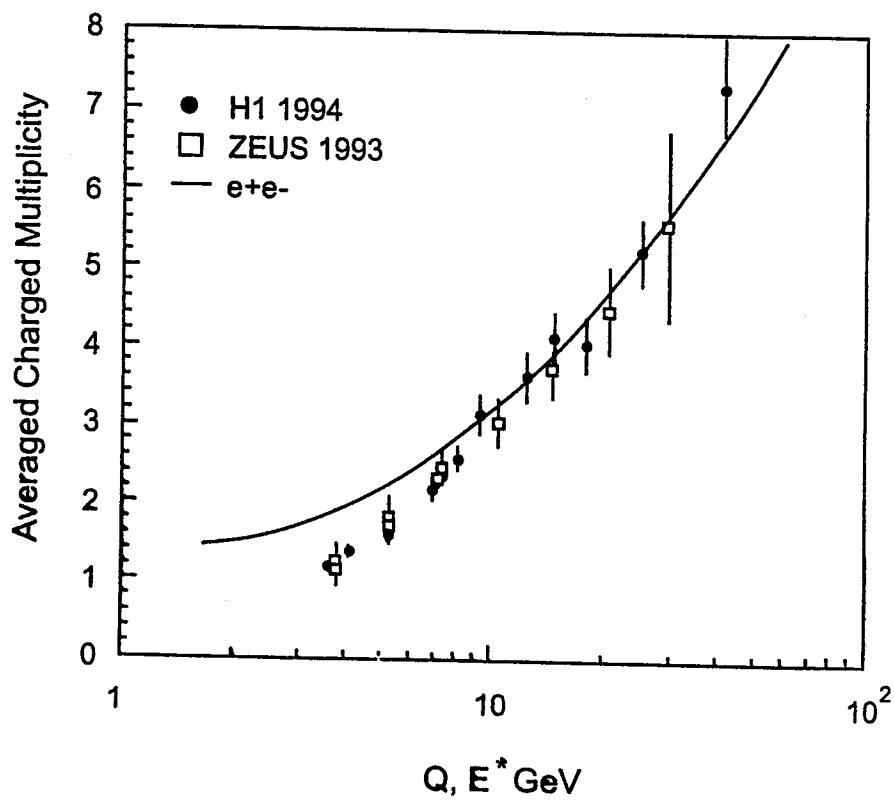
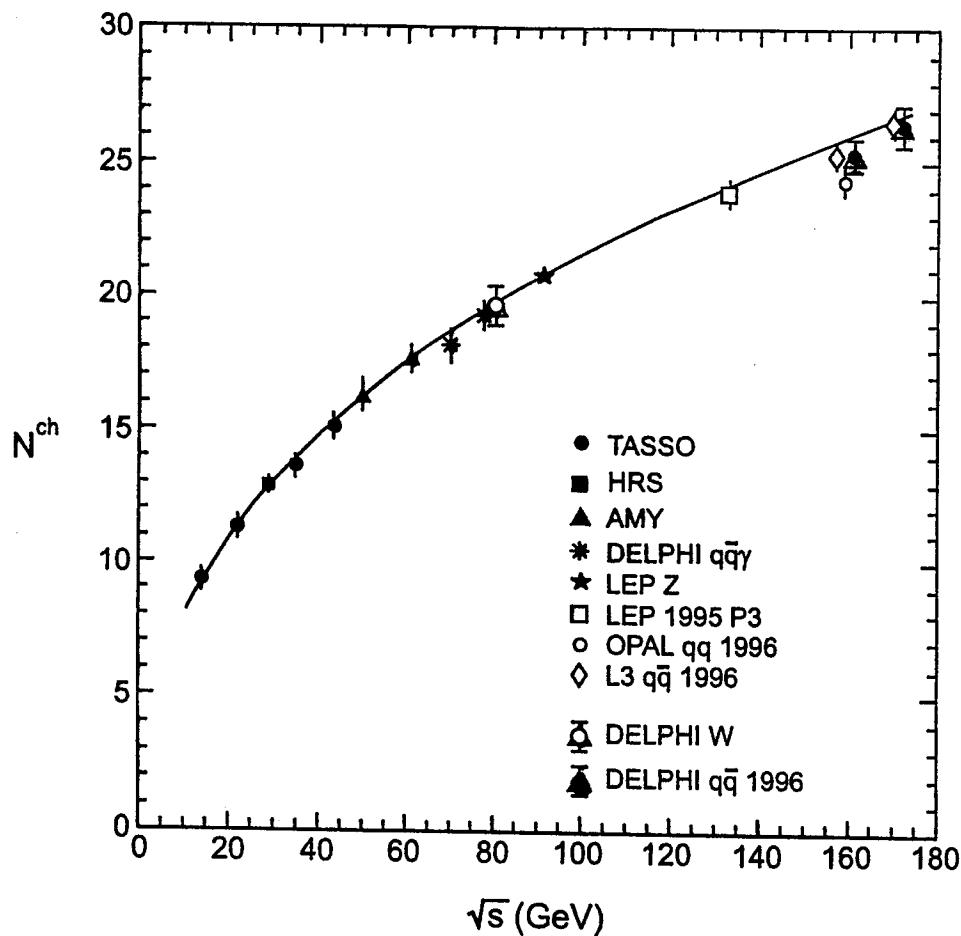
$$\text{factorial } F_q = \langle h(h-1) \dots (h-q+1) \rangle = \int p(p_1 \dots p_q) d\mathbf{p}_1 \dots d\mathbf{p}_q$$

Kumulant $K_2 = F_2 - 1, \dots$

genuine correlations

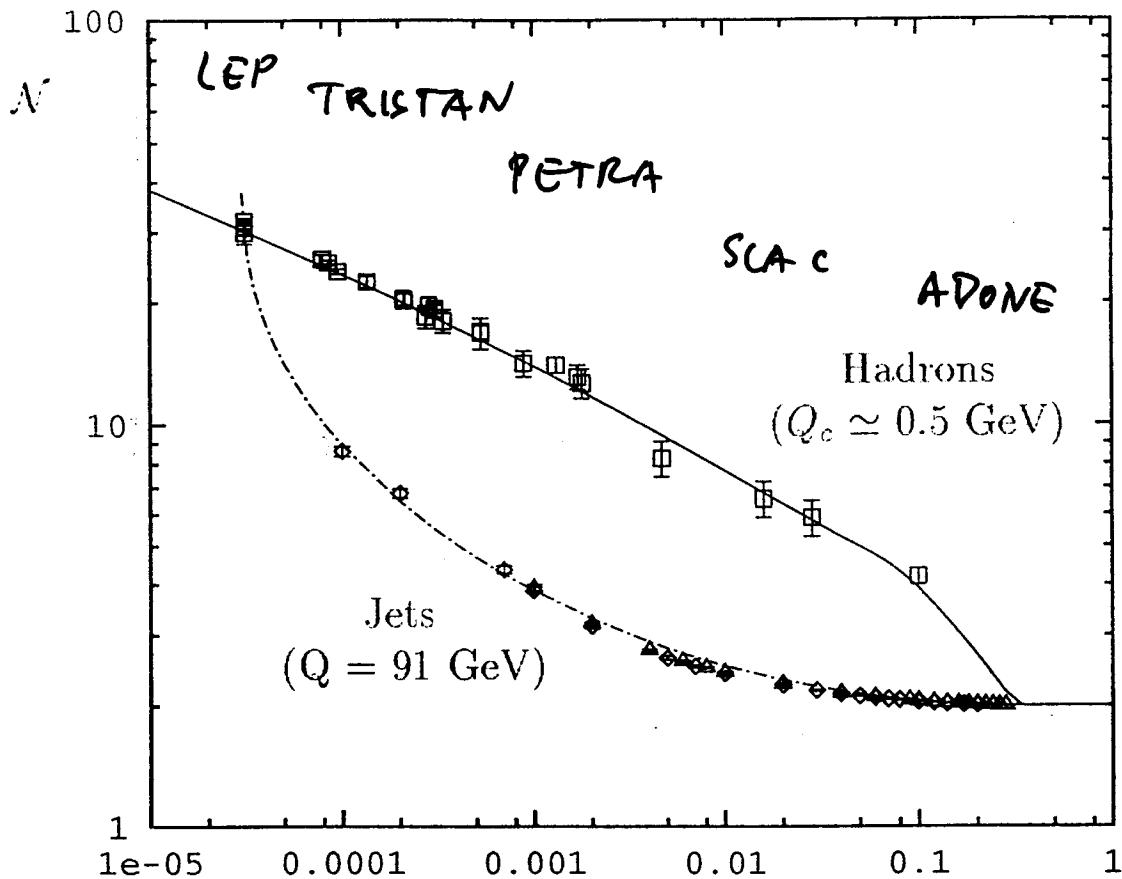
mean particle multiplicity

MLLA
K free



Hadron and Jet multiplicities
numerical solution of
MLLA evolution eqn.
+ $O(\alpha_s)$ ME

Lupia
W.O.
'98



$$y_c = \left(\frac{Q_c}{Q} \right)^2$$

Durham algorithm

$$K_T > Q_c$$

⇒ strong particle production if $\alpha_s(K_T)$ large.

Parameters $\Rightarrow K=1$

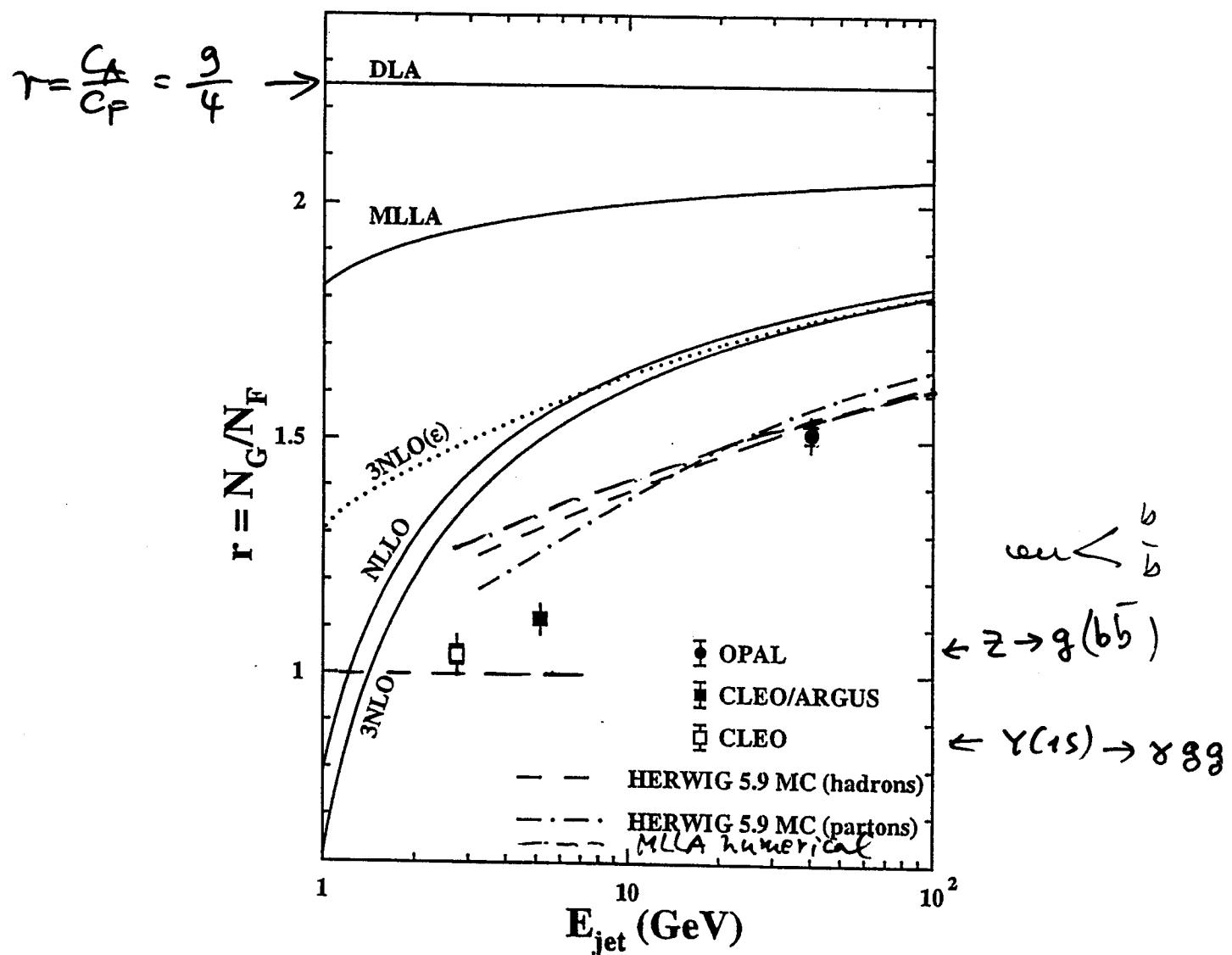
$$\Lambda = 500 \text{ MeV}, \lambda = 0.015 \quad (\text{Durham } K_T)$$

$$\Lambda = 350 \text{ MeV}, \lambda = 0.015 \quad (\text{standard } K_T)$$

$$K_T = z(1-z)Q$$

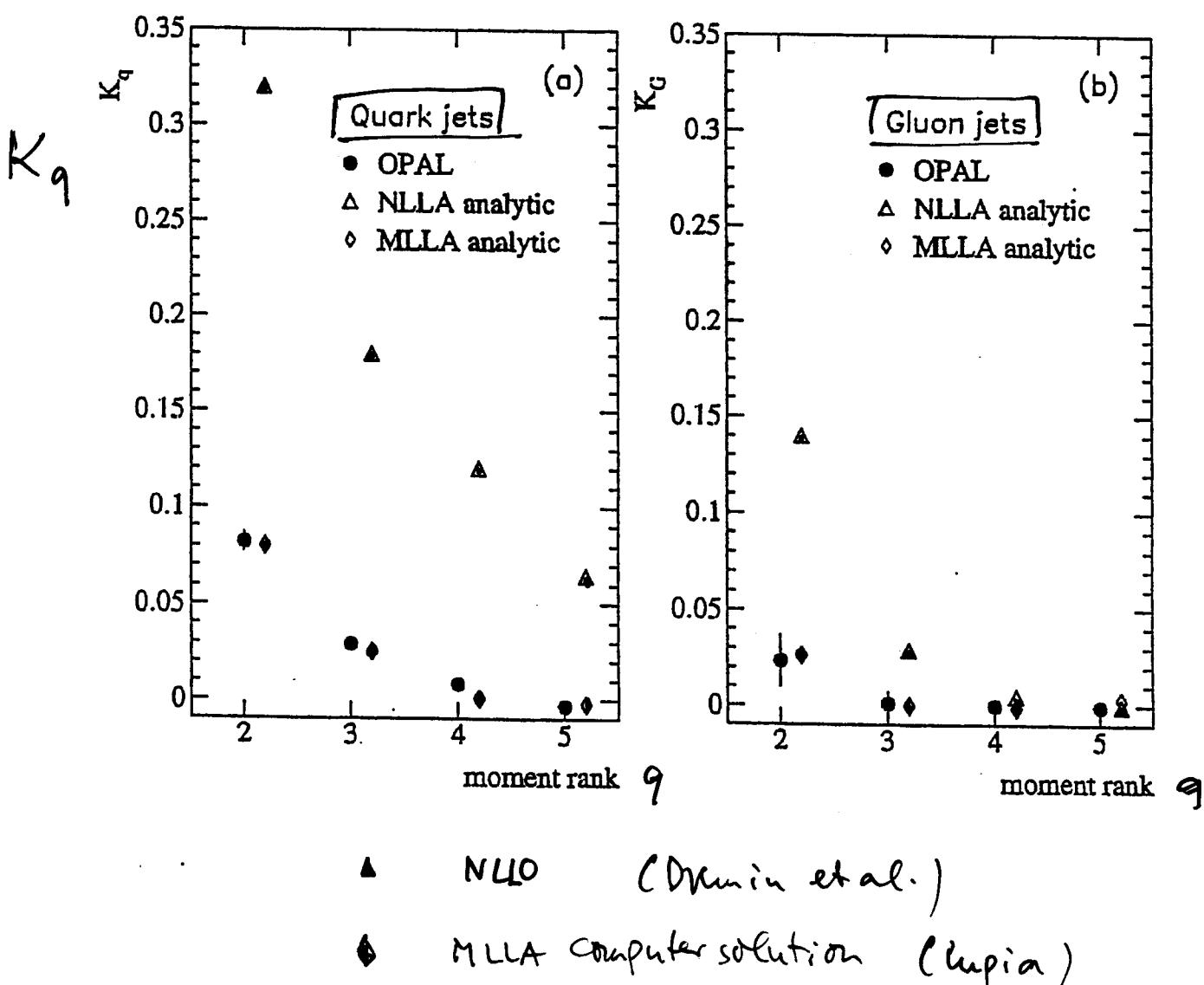
$$\lambda = \ln \frac{Q_0}{\Lambda}$$

$$r = \frac{\text{multiplicity in g-jet}}{\text{multiplicity in q-jet}}$$

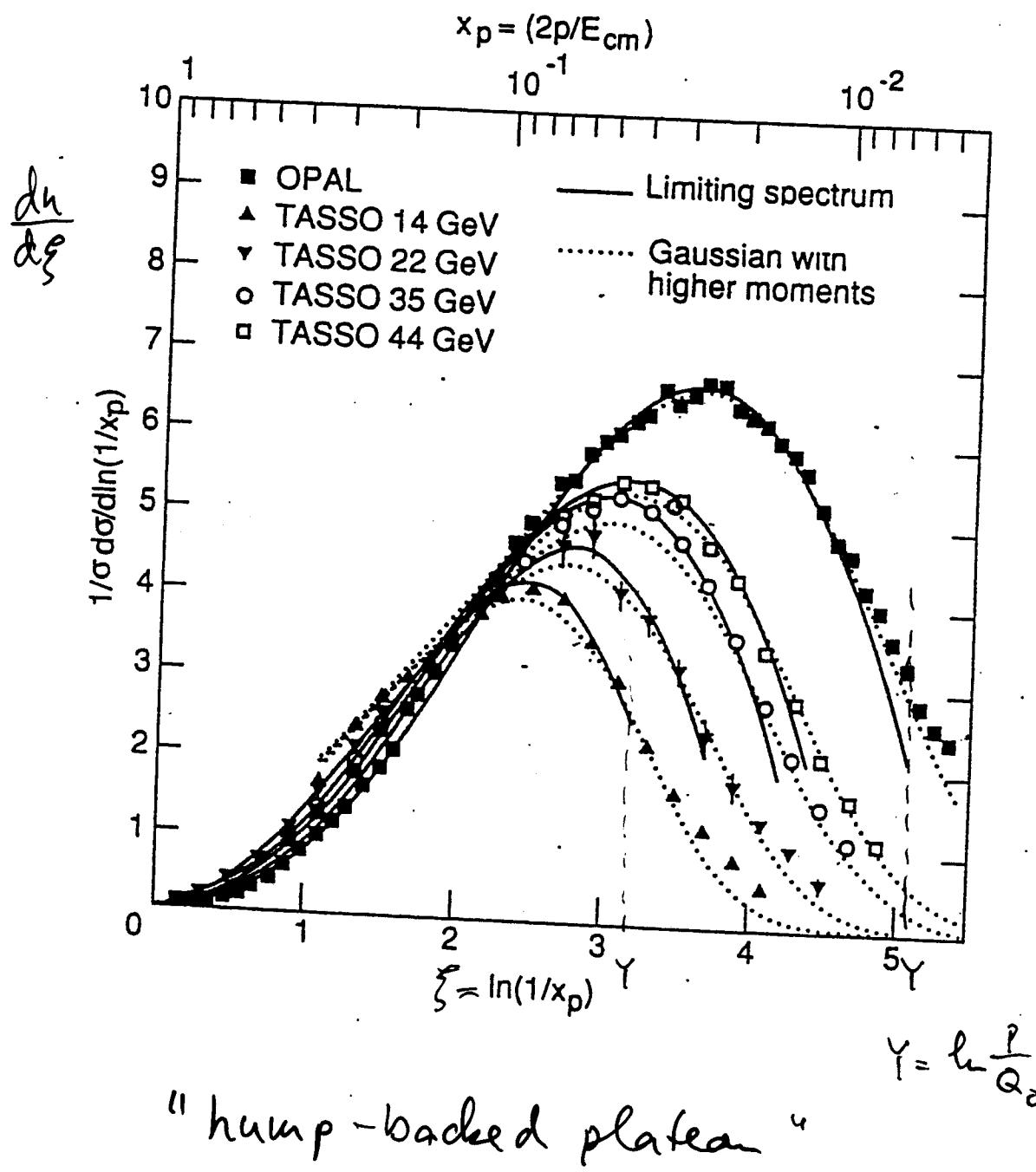


Kumulant moments of multiplicity distribution

O. Biebel / Physics Reports 340 (2001) 165–289



Inclusive spectra in $\xi = \ln \frac{1}{x}$



Shape parameters of ξ -distribution

Moments in MCCA

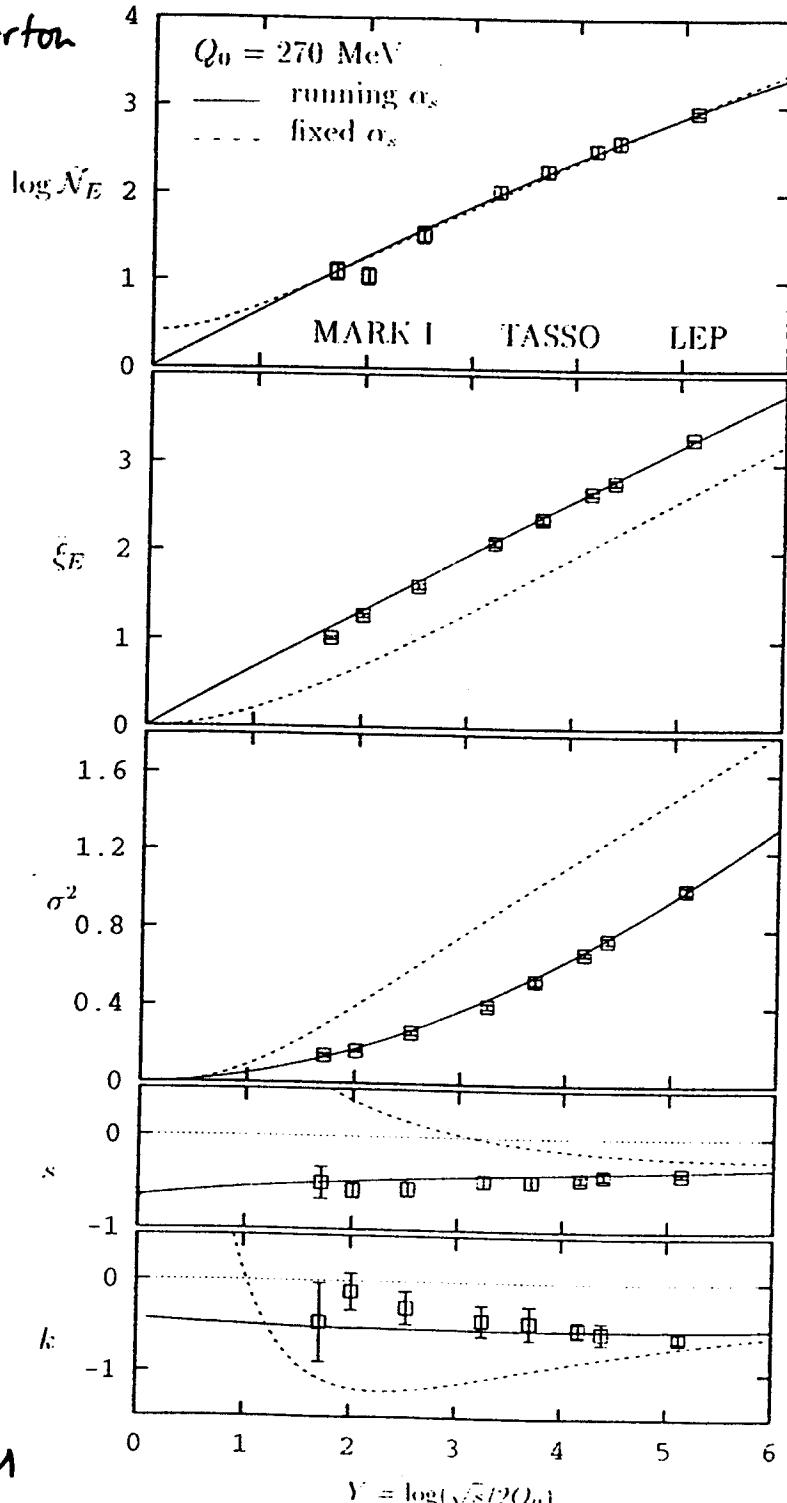
Dokshitzer
Khoze, Troyan
1992

$$\frac{\langle \xi^q \rangle}{Y^q} = P_0^{(q)}(B, z) + \frac{2}{z} \frac{I_{3+2}(z)}{I_{3+1}(z)} P_1^{(q)}(B, z)$$

$$z \sim \sqrt{Y}$$

$$N_{\text{hadron}} = K N_{\text{parton}}$$

$$K \sim 2$$



Skewness

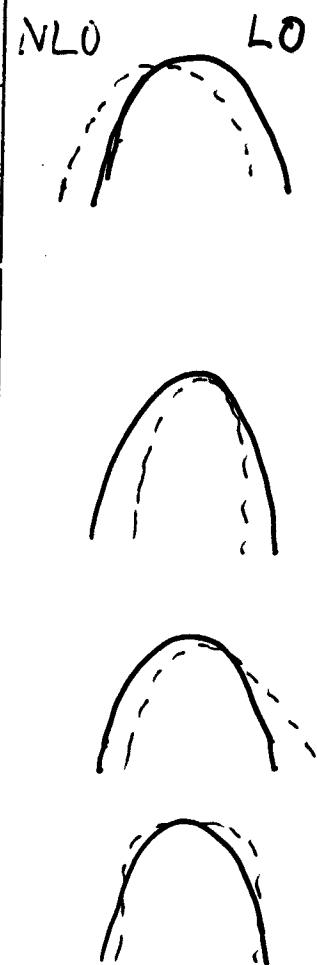
Kurtosis

$$1) Q_0/\Lambda \lesssim 1.1$$

2) Threshold limit

3) running α_s

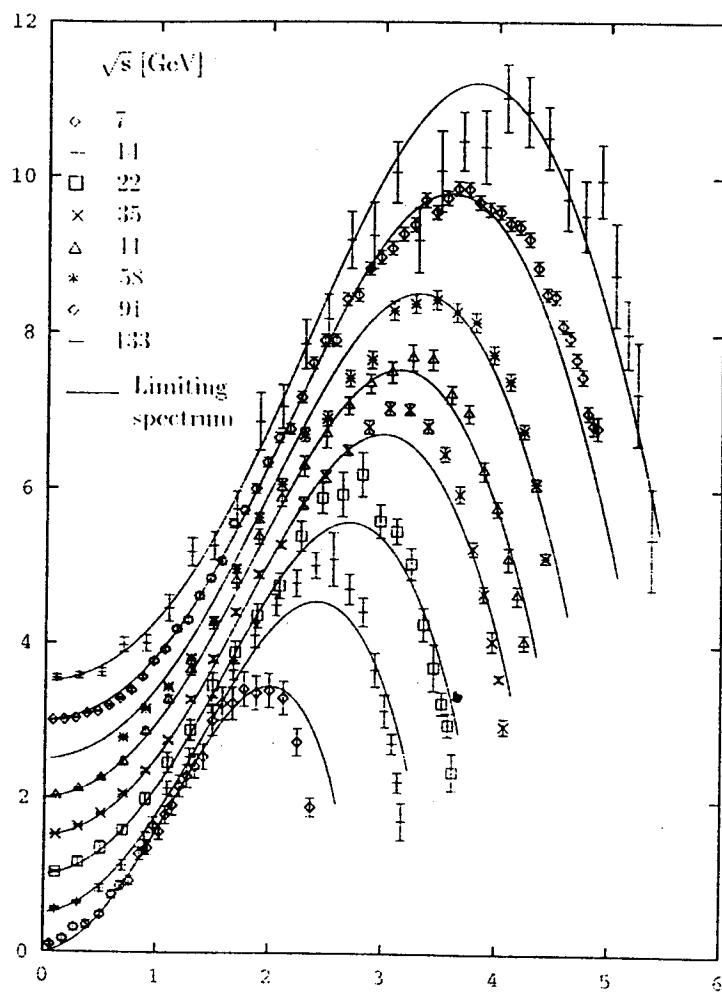
$$Y = \ln\left(\frac{\sqrt{s}}{2Q_0}\right)$$



Inclusive Momentum spectrum

$$D(\xi, \gamma) \simeq \frac{N(\gamma)}{\sqrt{2\pi}\sigma^2} e^{-\frac{(\xi-\bar{\xi})^2}{2\sigma^2}}, \quad \xi = \ln \frac{E}{Q_0}, \quad \gamma = \ln \frac{E}{Q_0}$$

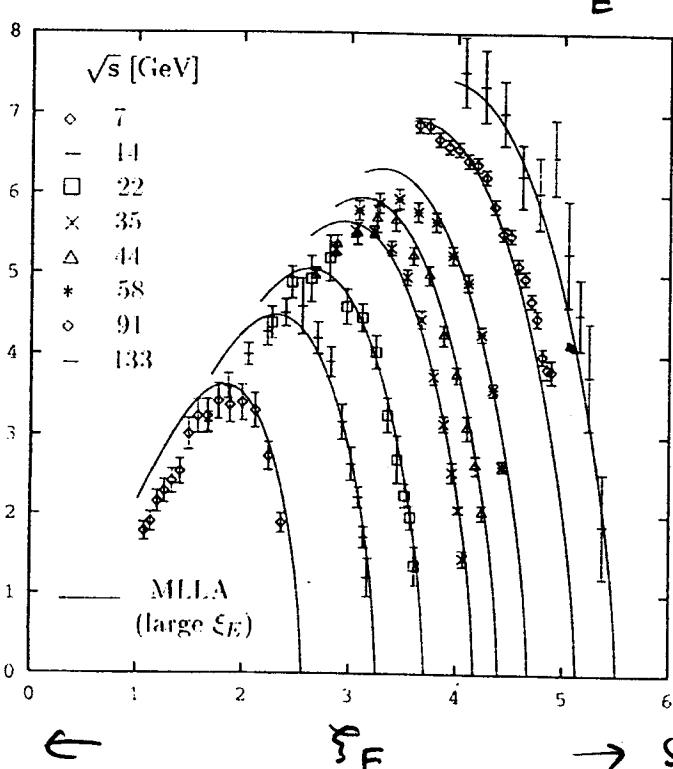
Moments: $\bar{\xi} = \frac{1}{2} \gamma + \dots$, $\sigma^2 = \frac{1}{3} \sqrt{\frac{b \gamma^3}{16 N_c}} + \dots$



"hump-backed plateau"
 Dokshitzer et al.
 Bassotto et al.
 Mueller (82)

Limiting spectrum

($Q_0 \sim \Lambda$)
 Dokshitzer, Kuraev,
 Troyan



MLLA, large ξ
 ($Q_0 > \Lambda$)

expansion $O(\alpha_s^2)$
 around soft limit
 $\xi \approx \gamma$

hard ← ξ_E → soft

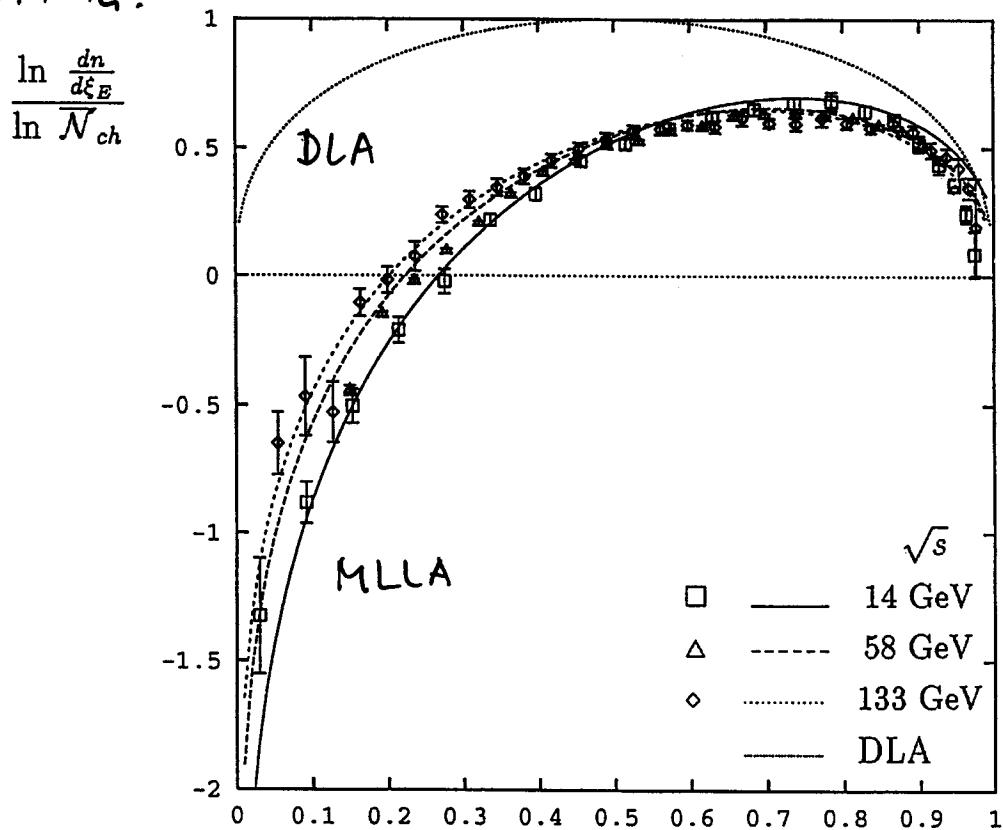
Lupia
 w.c.
 EPJ-C (98)

Asymptotic (DLA) limit of particle spectra

$$s \rightarrow \infty \quad \frac{dn}{d\xi} \sim \bar{N} e^{-\frac{(\xi - \bar{\xi})^2}{2\sigma^2}}$$

$$\bar{N} \sim e^{c\sqrt{Y}}, \quad \bar{\xi} = \frac{1}{2} Y, \quad \sigma^2 \sim Y^{3/2}.$$

Scaling limit is:



$$\zeta_E \equiv \frac{\xi_E}{Y} \equiv \frac{\ln \frac{E_{\text{jet}}}{E}}{\ln \frac{E_{\text{jet}}}{Q_0}}$$

hard \leftarrow \rightarrow soft

\Rightarrow slow approach of asymptotic limit
but: soft particles already asymptotic
($\xi \sim Y$)

Soft limit of particle spectrum

soft gluon emission in $q\bar{q}$ jet system



gluon with large wavelength "sees" only total charge
(coherent emission from all sources)

→ Born term dominates $\mathcal{O}(\alpha_s)$
like QED Bremsstrahlung

$$\frac{dn}{dy dP_T} \sim C_{A,F} \frac{\alpha_s(P_T)}{P_T} \left(1 + \mathcal{O}\left(\ln \frac{\ln P_T/\Lambda}{\ln Q/\Lambda} \ln \frac{\ln P_T/\Lambda}{\ln Q/\Lambda}\right) \right)$$

↑
 energy independent ↑
 vanishes for $P_T \rightarrow Q$

Predictions

$C_A = 3$
 $C_F = 4/3$
 $\omega = 0$.

- 1) $\frac{dn}{dy dP_T^2}$ | _{$P_T \rightarrow 0$ or $P_T \rightarrow \infty$} - energy independent.
- flat plateau
- 2) $\frac{dn}{dy dP_T^2} \sim \begin{cases} C_A = 3 & \text{gluon jet} \\ C_F = 4/3 & \text{quark jet} \end{cases} \frac{N_g}{N_q} = \frac{9}{4} = 2.25$
(total multiplicity $\frac{N_g}{N_q} \sim 1.5$ at LEP)
- 3) multiplicity distribution of particles at small P_T :
Poissonian

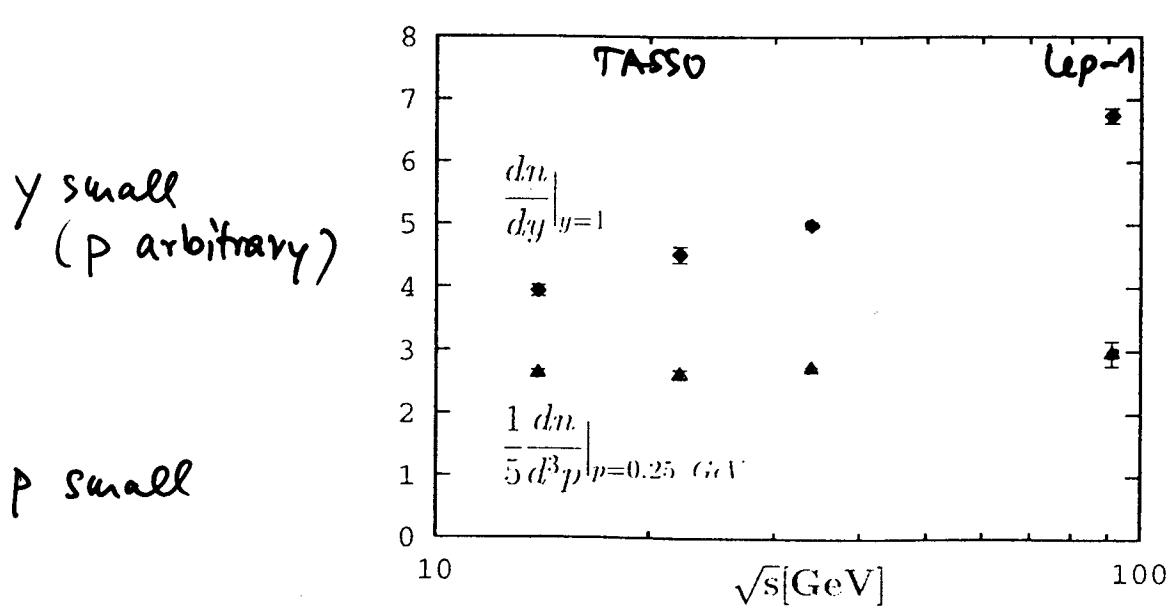
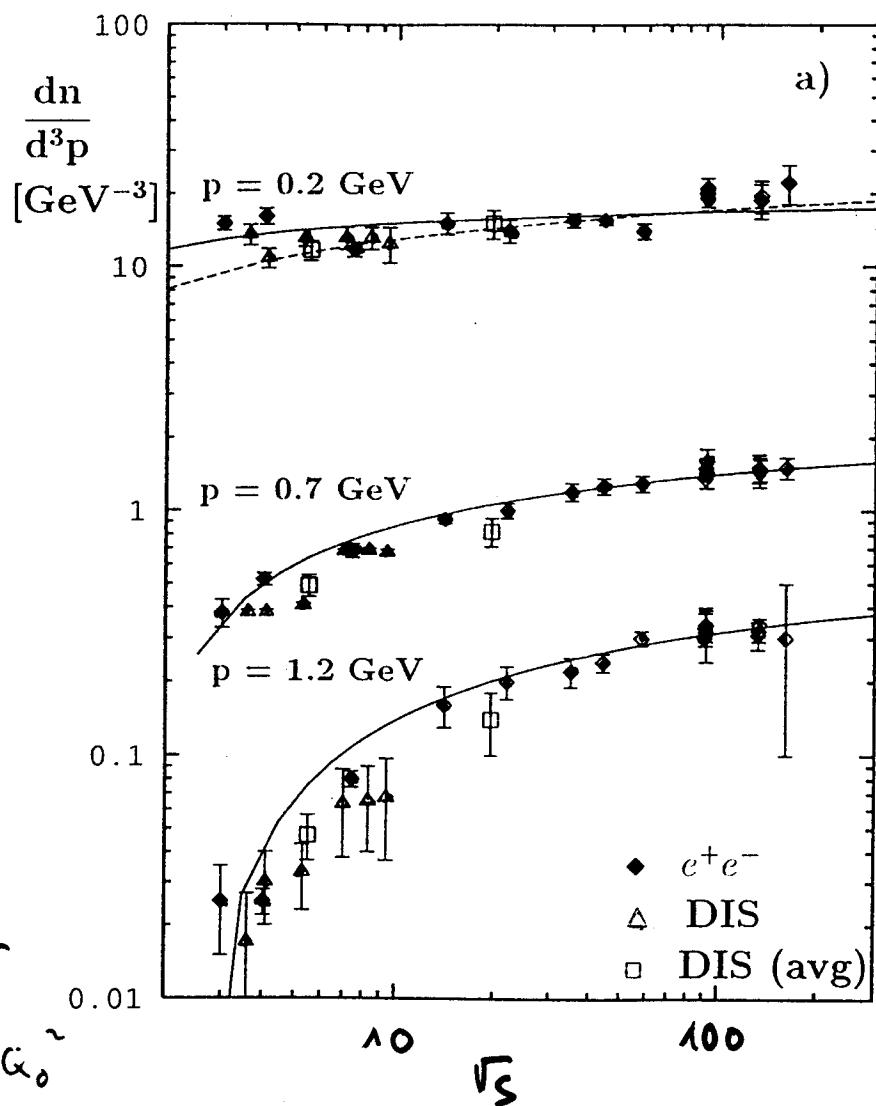
Realistic tests

e^+e^- 3jet events

DIS - resolved and unresolved photo production

Inclusive particle spectrum

Kluge
Urgin
Ochs



the spectrum tends to become energy independent
for $p \rightarrow 0$

B. Gary
OPAL

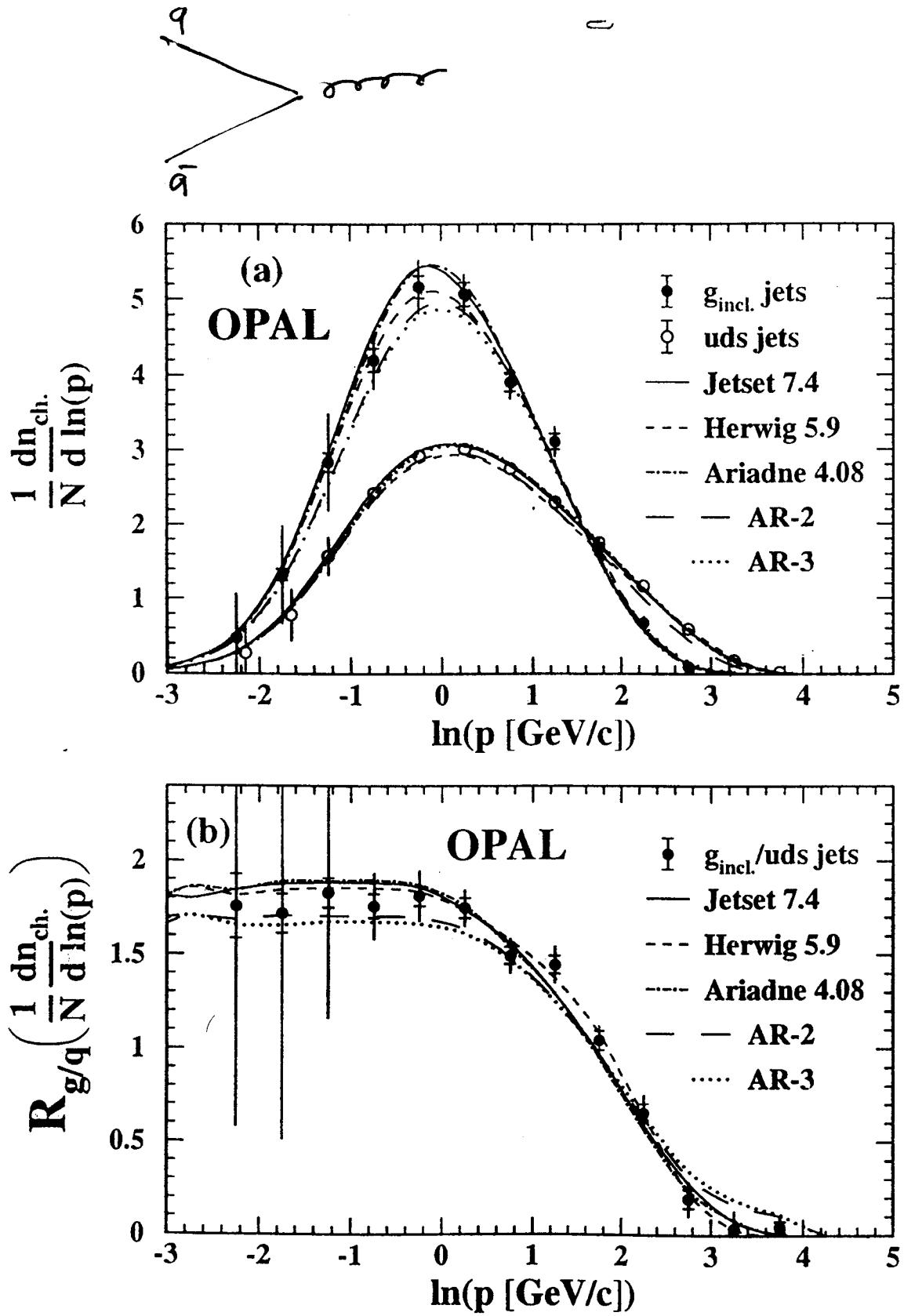
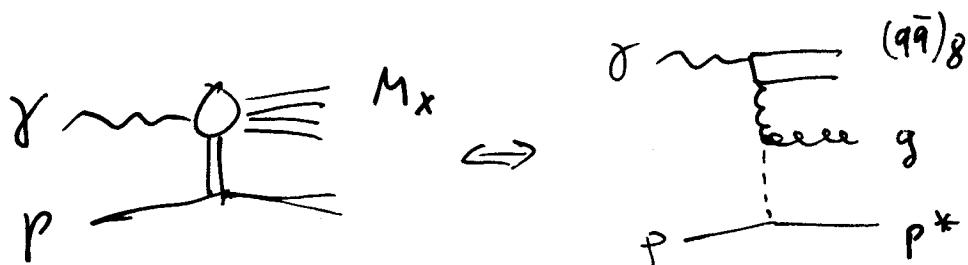


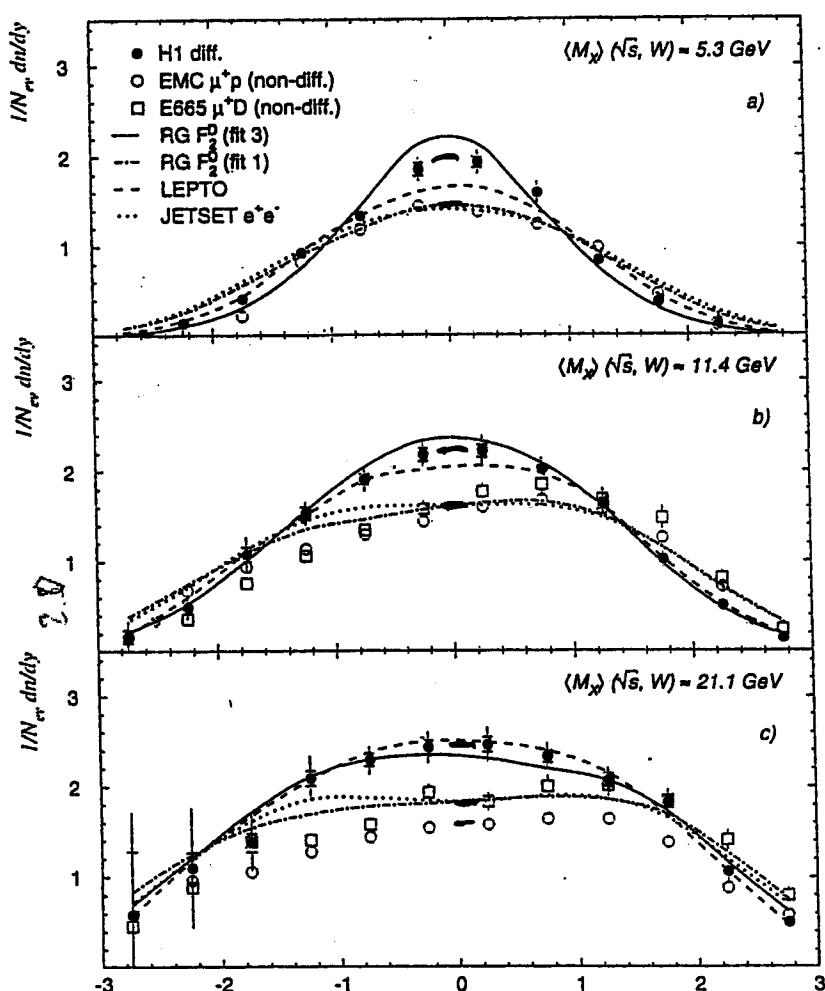
Figure 7: (a) Corrected distributions of the logarithm of charged particle momentum, $\ln(p)$, for 40.1 GeV $g_{\text{incl.}}$ gluon jets and 45.6 GeV uds quark jets. (b) The ratio of the gluon to quark jet $\ln(p)$ distributions for 40.1 GeV jets. The total uncertainties are shown by vertical lines. The experimental statistical uncertainties are indicated by small horizontal bars. (The statistical uncertainties are too small to be seen for the uds jets.) The predictions of various parton shower Monte Carlo event generators are also shown. These data are tabulated in Table 3.

Diffraction in DIS

411



$$\frac{dy}{dy} \\ M_X = 5.3 \text{ GeV}$$



$$= 11.4 \text{ GeV}$$

$$= 21.1 \text{ GeV}$$

$$\frac{N_g}{N_q} \text{ at } y \approx 0$$

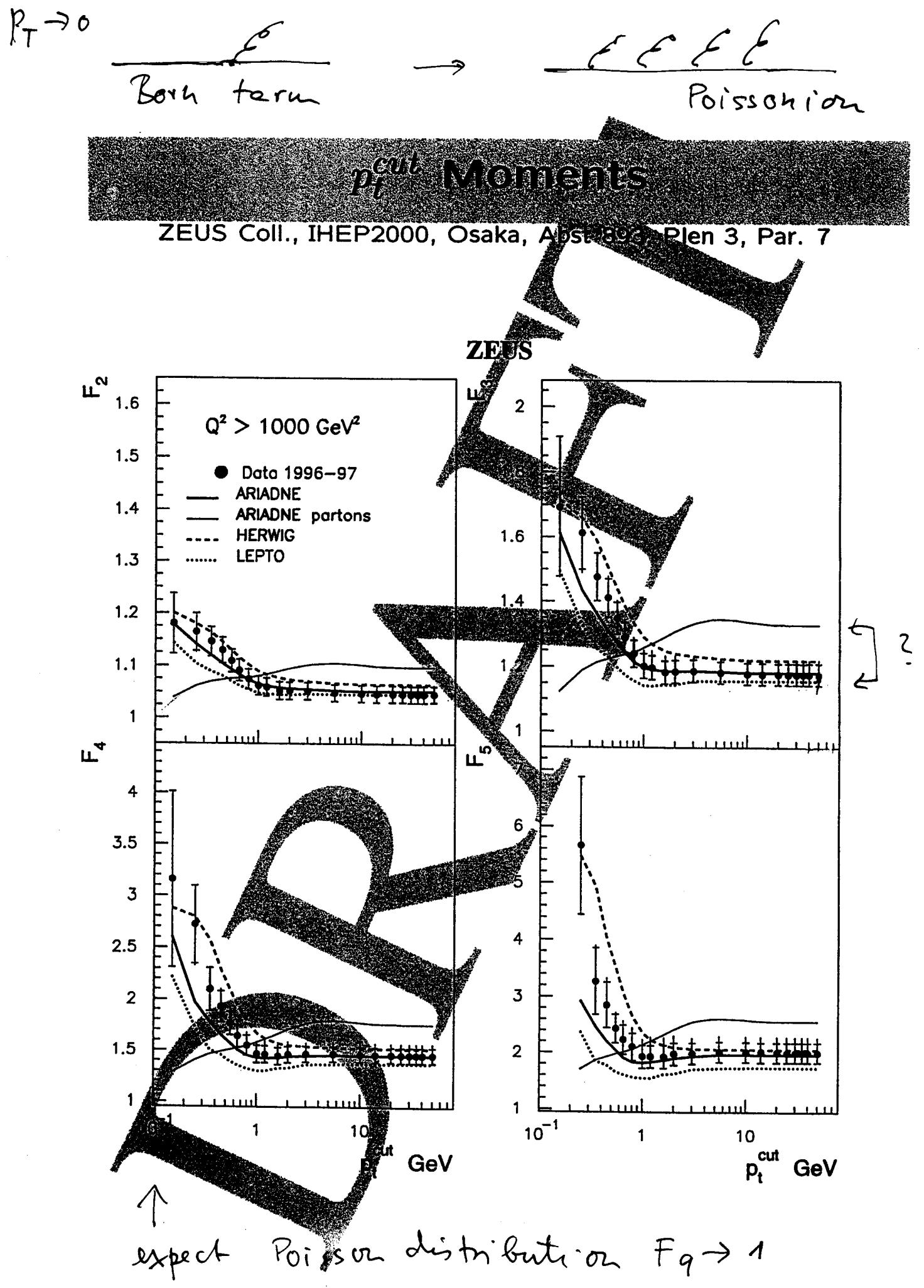
$$\tau \approx 1.4$$

$$1.4$$

$$1.3-1.6$$

$$y \rightarrow$$

further increase for $p \rightarrow 0$?

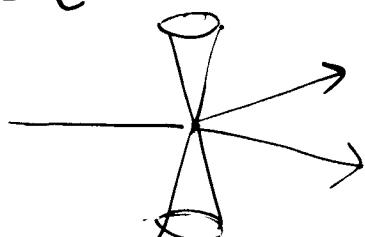


Study soft radiation perpendicular to event plane

Test of duality picture for small scale

Khoze
Lugia
Raths

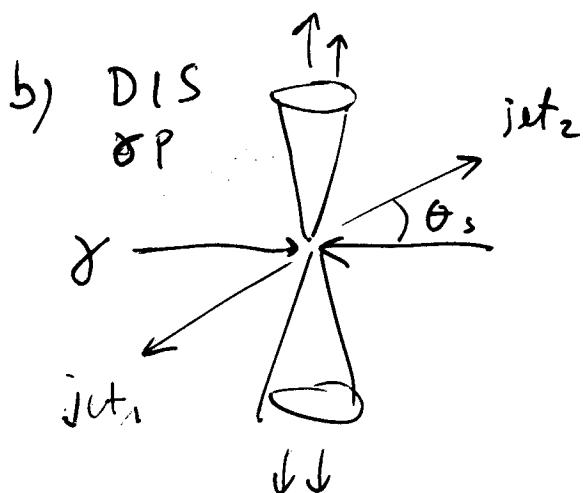
a) e^+e^-



$\sim C_F$

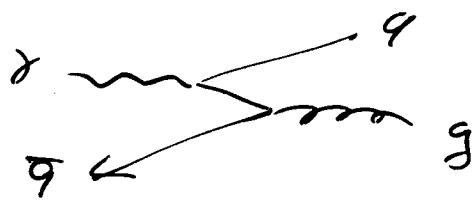


$\sim N_c$



photoproduction of dijets

- direct

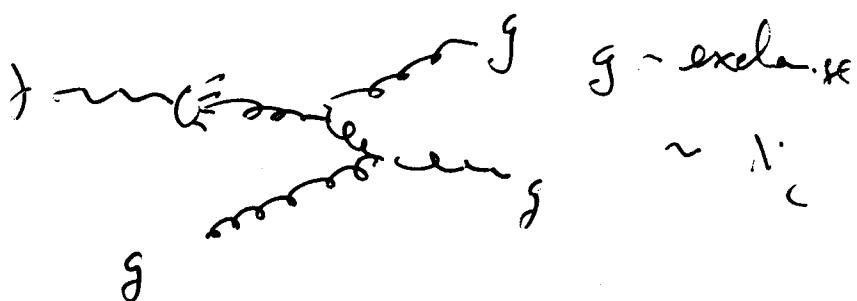


small θ_S

γ -exchange

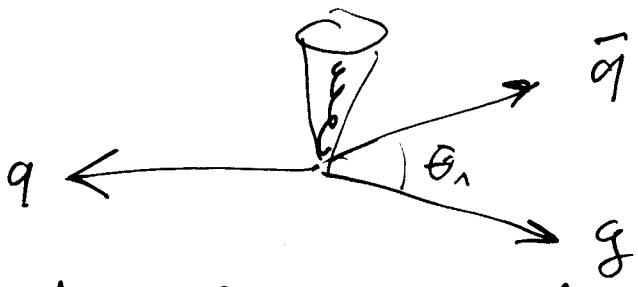
$\sim C_F$

- indirect



g -exchange

$\sim N_c$

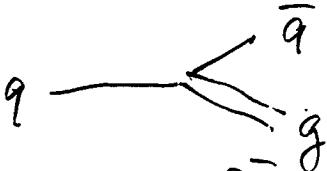
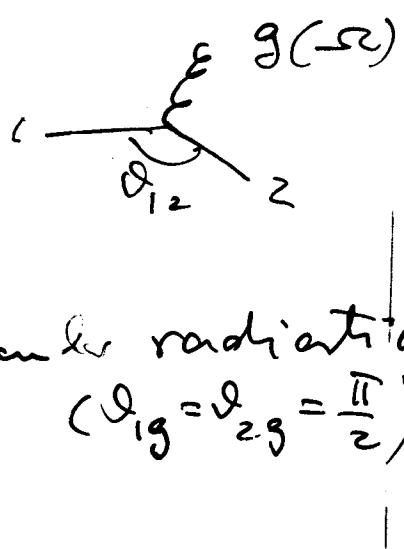


radiation from 1 dipole

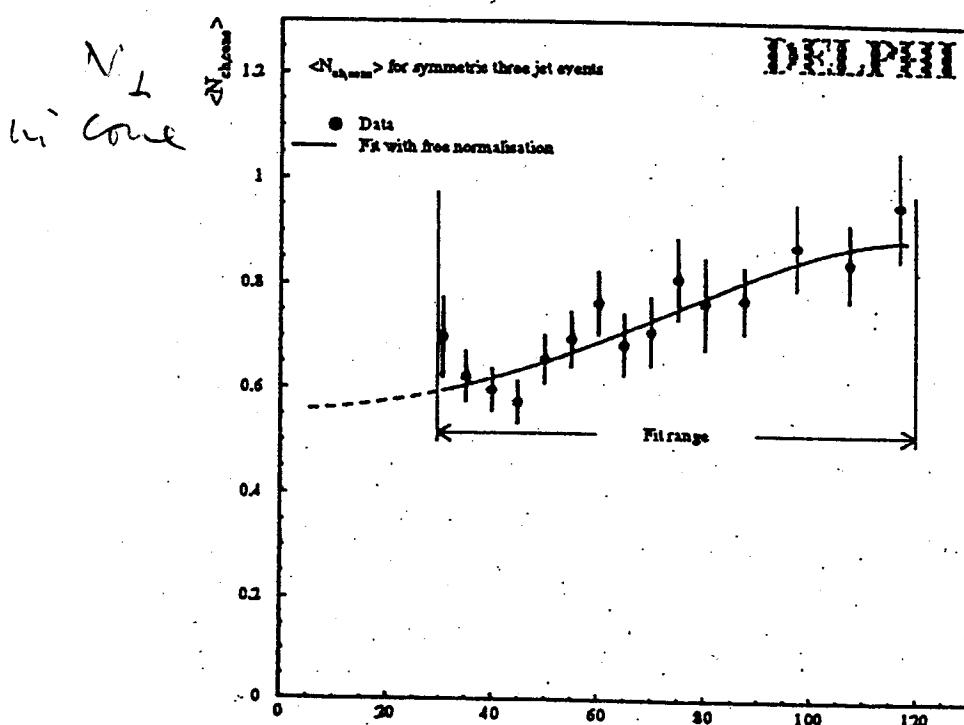
$$\frac{dn}{ds} \approx \frac{1 - \cos \delta_{12}}{(1 - \cos \delta_{1g})(1 - \cos \delta_{eg})}$$

$\rightarrow 1 - \cos \delta_{12}$ for perpendicular radiation
($\delta_{1g} = \delta_{eg} = \frac{\pi}{2}$)

radiation from $(q\bar{q}g)$

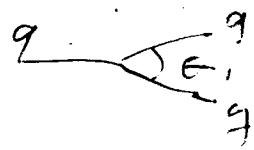


$$\frac{dN^{\text{q}\bar{q}g}}{dN^{\text{q}\bar{q}}} = \frac{N_c}{4C_F} \left[(1 - \cos \delta_{gg}) + (1 - \cos \delta_{g\bar{g}}) - \frac{1}{N_c^2} (1 - \cos \delta_{g\bar{g}}) \right]$$



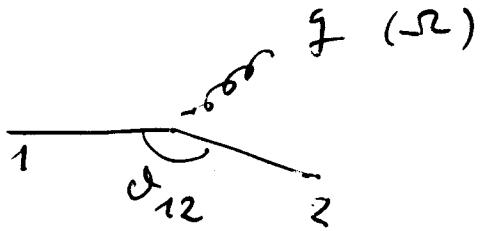
$T \sim 1.64$

- coherent emission from 3 partons!
- multiplicity depends on eff. color charge
- momentum dependence?



Radiation or perpendicular to scattering plane

$q\bar{q}$ dipole



$$\frac{dn}{dr} \approx \frac{1 - \cos \theta_{12}}{(1 - \cos \theta_{1g})(1 - \cos \theta_{2g})} \xrightarrow{\text{perp.}} 1 - \cos \theta_{12} \xrightarrow{\theta_{12} \rightarrow \pi}$$

$$\text{Ratio } R_\perp = \frac{\frac{dn}{dr}}{\left(\frac{dn}{dr}\right) \text{Dipole at } \theta_{12} = \pi}$$

Direct production

$$\gamma g \rightarrow q\bar{q}: \quad R_\perp = 1$$

$$\gamma q \rightarrow q\bar{q}: \quad R_\perp = \frac{N_c}{4C_F} \left[3 - \cos \theta_s - \frac{1}{N_c^2} (1 + \cos \theta_s) \right]$$

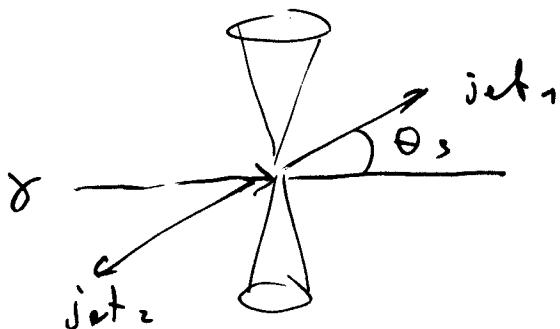
$$R_\perp \rightarrow 1 \quad \text{small angles } \theta_s$$

Indirect production

$$\gamma g \rightarrow gg: \quad R_\perp = \frac{N_c}{4C_F} (5 - \cos \theta_s)$$

$$\gamma g \rightarrow gg: \quad \text{etc.} \quad R_\perp = \frac{N_c}{4C_F} (4 - \frac{1}{N_c^2} (1 - \cos \theta_s))$$

$$R_\perp \rightarrow \frac{N_c}{C_F} \quad \text{small angles } \theta_s$$



Measure Ratio

$$R(P_T, \theta_s) = \frac{\left(\frac{dn}{dP_T} \right) \text{resolved at } \theta_s}{\left(\frac{dn}{dP_T} \right) \text{direct at } \theta_s}$$

$$R(P_T, \theta_s) \rightarrow \frac{9}{4} \quad \text{for } \theta_s \rightarrow 0$$

Test run with HERWIG
(4.5 pb⁻¹)

Bufforworth
Khote
Ochs
hep-ph/9901419
(Durham)

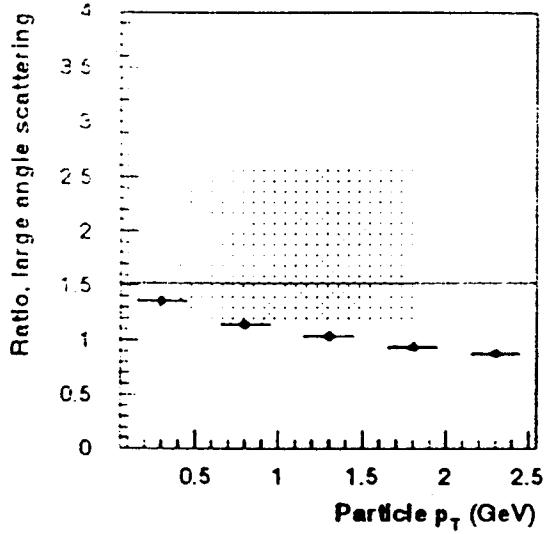
direct : $x_{obs} > 0.9$

resolved : $x_{obs} < 0.75$

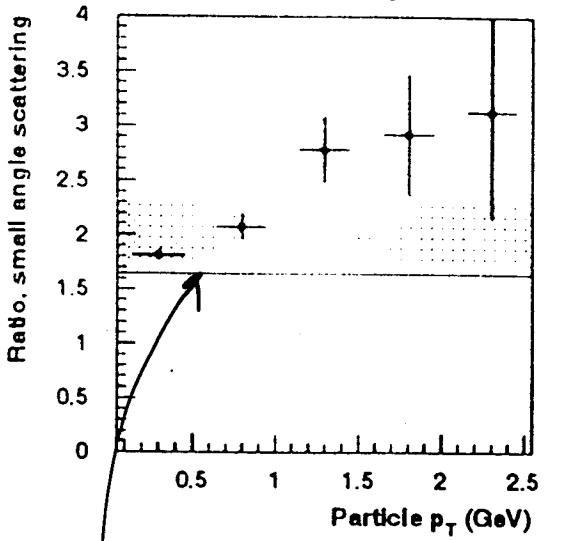
$R(p_T, \theta_S)$

$0 < \cos \theta_S < 0.64$

resolved
direct



$\rightarrow p_T$



$\rightarrow p_T$

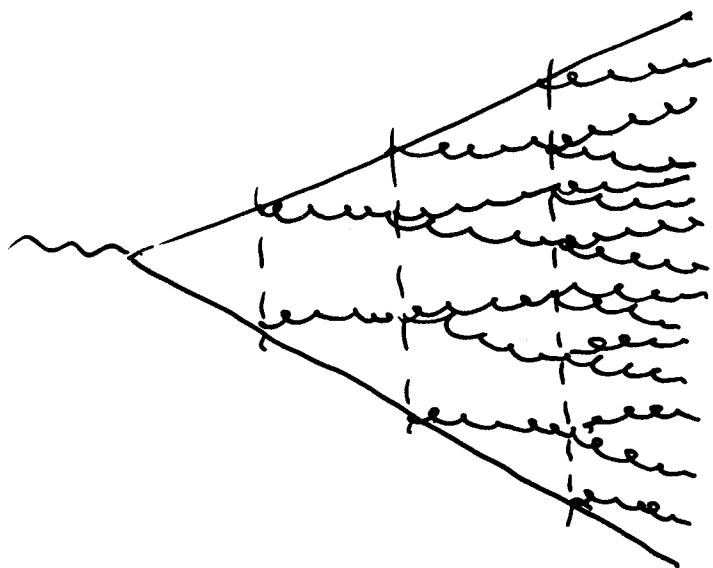
theor. expectation
using HERWIG
structure fcts.

\rightarrow "data" approach expectation for small p_T

\rightarrow studies towards small angles θ_S
appear feasible

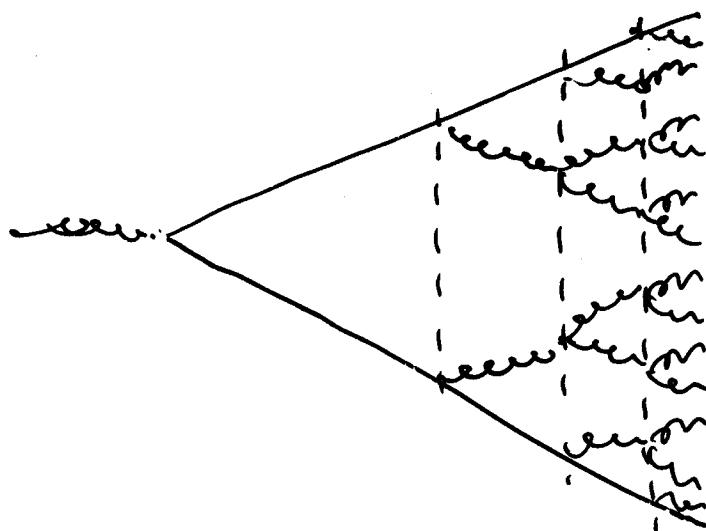
Multi-particle correlations

selfsimilar cascade : fractal



$$\Omega(\delta) \sim \delta^q$$

QCD cascade : selfsimilarity broken by $\alpha_s(Q^2)$

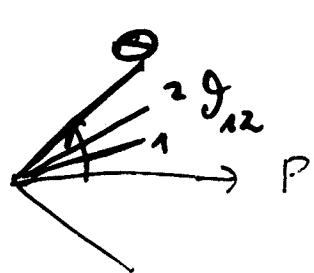


→ resolution power

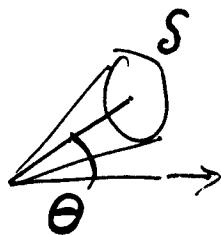
$\alpha_s(Q^2) \rightarrow$ increasing

angular correlations

W.-O.-E Wosiek
 Dokshitzer Dres.;
 Brack, Meunier
 Peschanski;

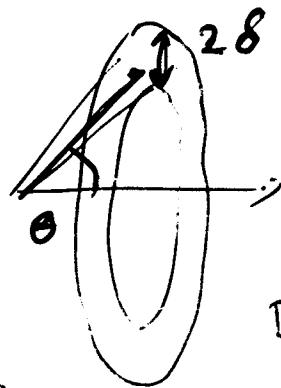


$$r(\delta_{12})$$



$$Dik_1 = 2$$

$$\text{moments } F_q$$



$$Dik_1 = 1$$

Analysis in DLA of QCD

$$h_q(\delta, \theta, p) \sim \left(\frac{\theta}{\delta}\right)^{\phi_q}$$

running α_s

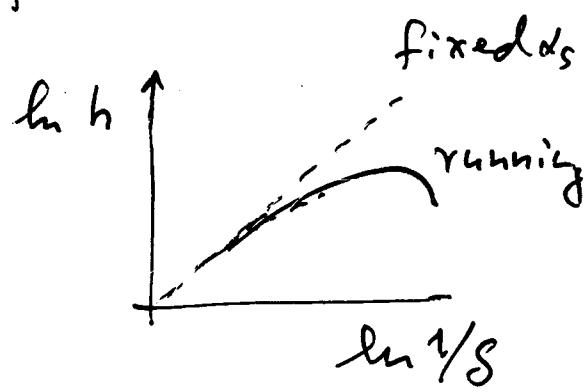
- bending at small angle
- running of initial slope

moments:

$$F_q \sim \left(\frac{\theta}{\delta}\right)^{\phi_q}, \quad \phi_q \approx (q-1)D - (q-\frac{1}{q})\gamma_0$$

$$\gamma_0^2 = \frac{6\alpha_s(p\theta)}{\pi}$$

for fixed α_s



- asymptotic scaling law (ε -scaling)

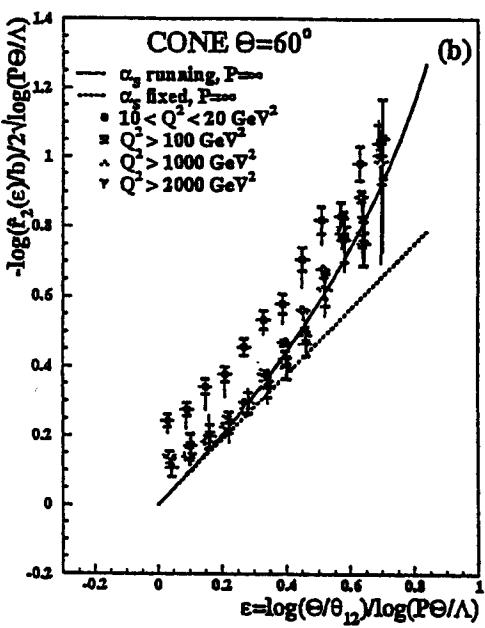
rescaled $\ln h_q$ depends only on ε

$$\varepsilon = \frac{\ln \theta/\delta}{\ln p\theta/\lambda}$$

$$(0 < \varepsilon < 1)$$

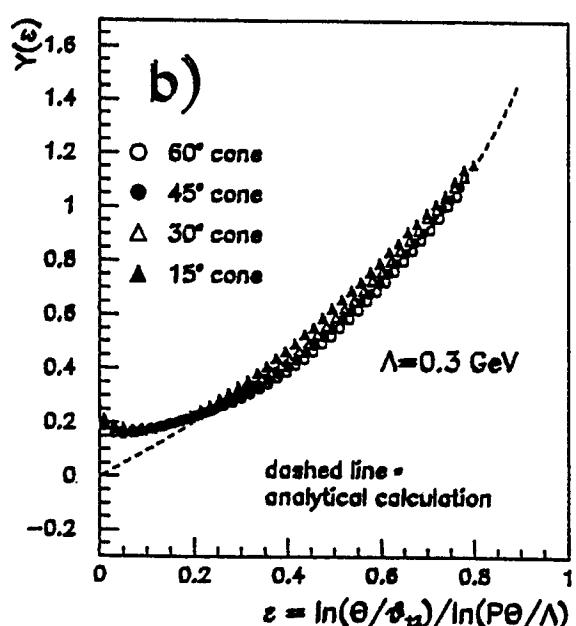
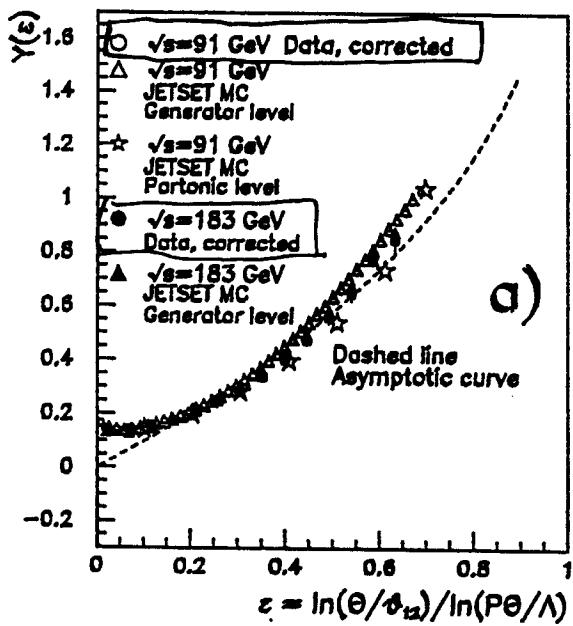
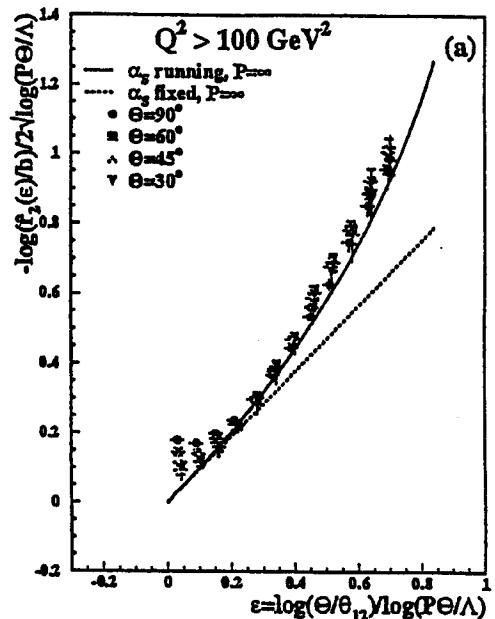
ξ - Scaling

Q varies
 θ fixed



θ varies
 Q fixed

Zens '98 (Vancouver)



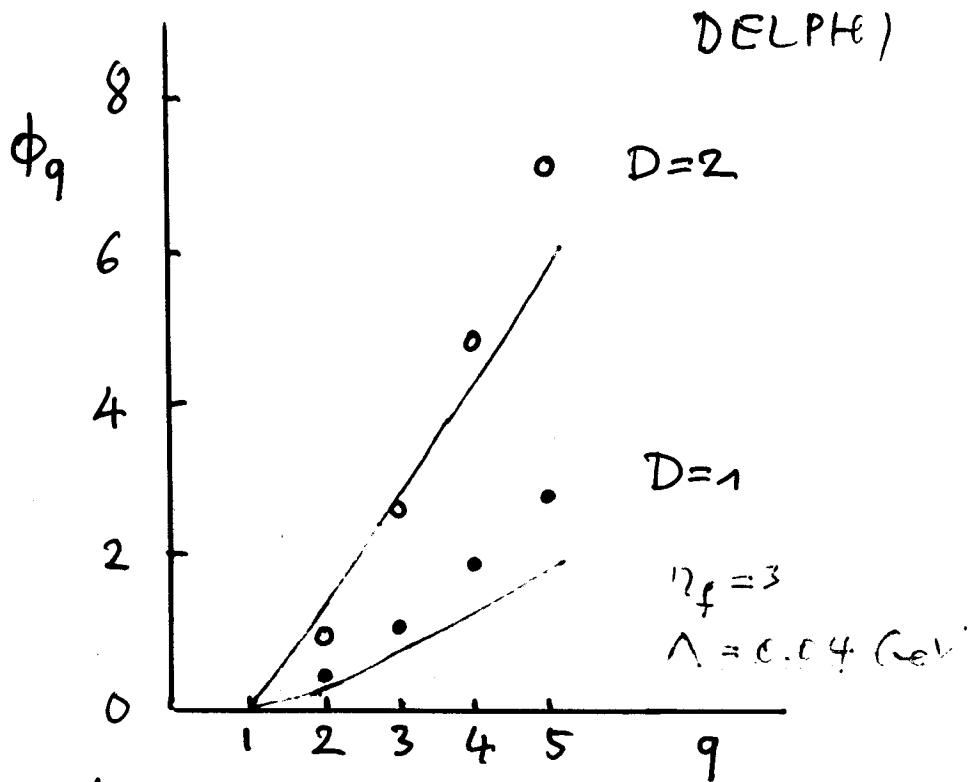
$$\xi = \frac{\ln \theta / \theta_{12}}{\ln P\theta / \Lambda}$$

- early approach to asymptotic
- Dendons.

Initial slopes

moments \bar{F}_q : dependence on q and D

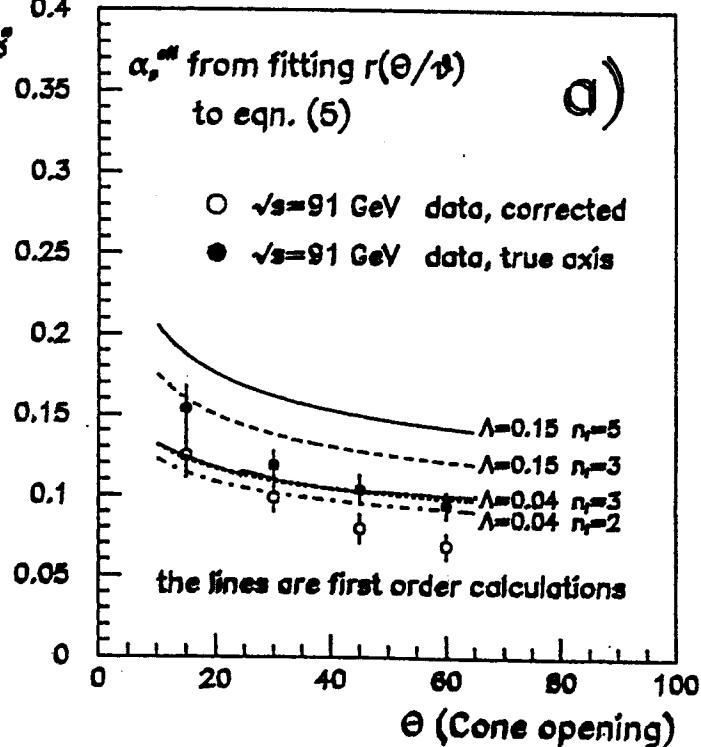
$$\phi_q \approx (q-1) D - (q-\frac{1}{q}) \gamma_0 \quad (\text{DLA})$$



correlation function: dependence on jet opening θ

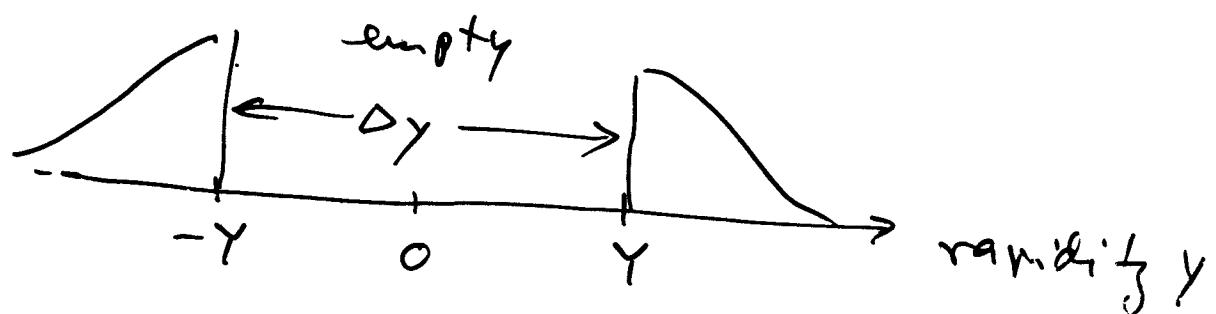
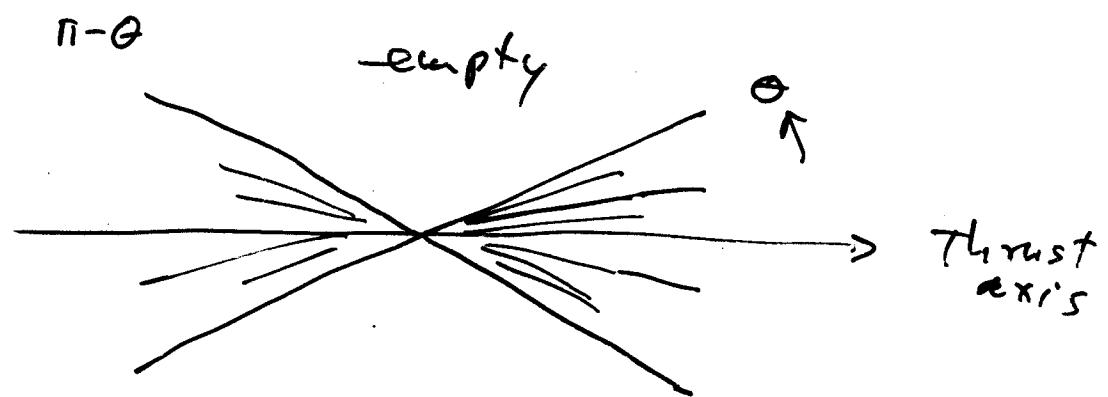
$$r(\theta_{12}) = \left(\frac{\theta}{\theta_0} \right)^{\gamma_0/2}$$

$$\theta_0 = \left(\frac{6 \alpha_s(p_\theta)}{\pi} \right)^{1/2}$$



Rapidity gaps in e^+e^- annihilation

with
T. Shimada
ISMD '99

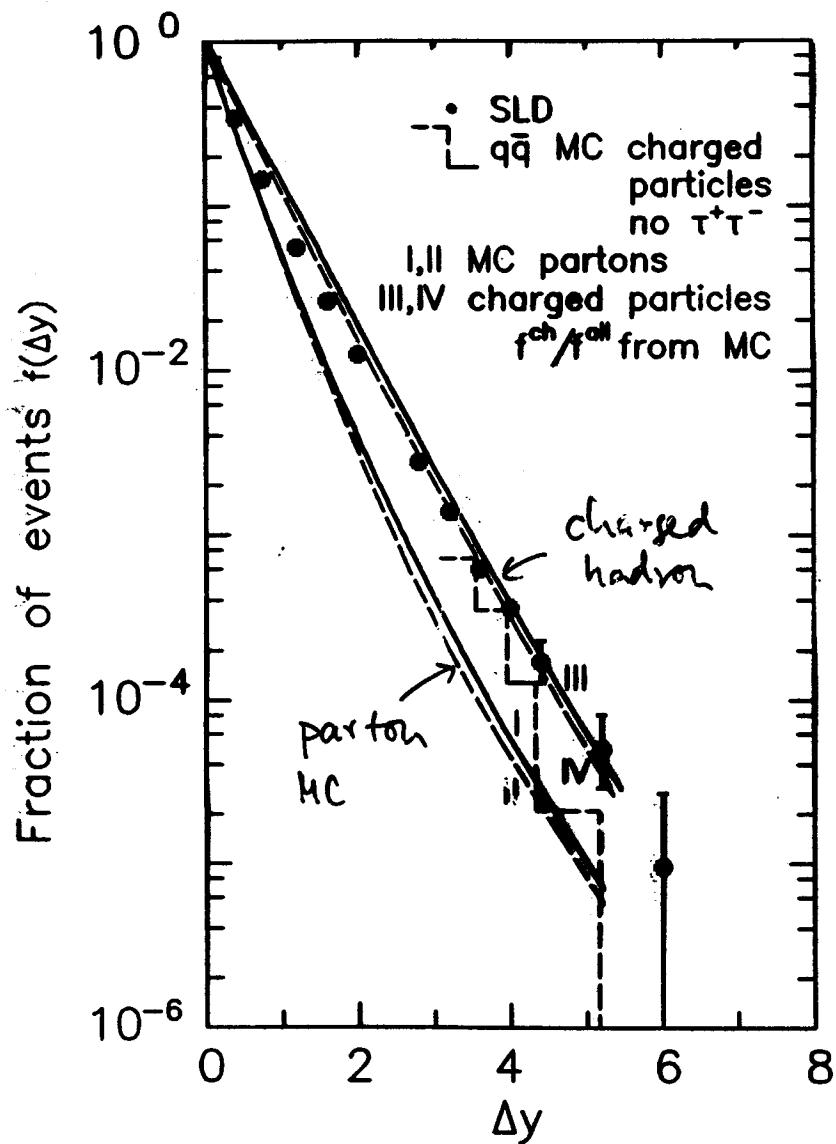


$$\text{rapidity} = -\ln \tan \frac{\theta}{2}$$

study effect of color neutralization
through $\bar{s}p$

Rapidity gaps in e^+e^- annihilation

W.O.
T. Shimada
'99



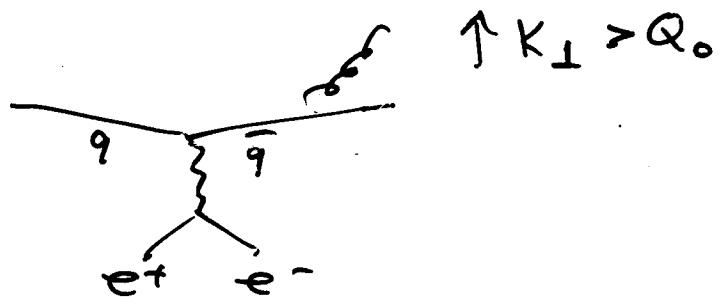
Ariadne MC with parameters Q_0, λ
parton fixed by jet & hadron multiplicity

$$\text{DLA: } f_{\text{gap}} \approx e^{-A_p \Delta y}, \quad A_p = \frac{4C_{MF}}{b} \ln \frac{\gamma}{\lambda}$$

$$\gamma = \ln \frac{P}{Q_0}, \quad \lambda = \ln \frac{Q_0}{\lambda}$$

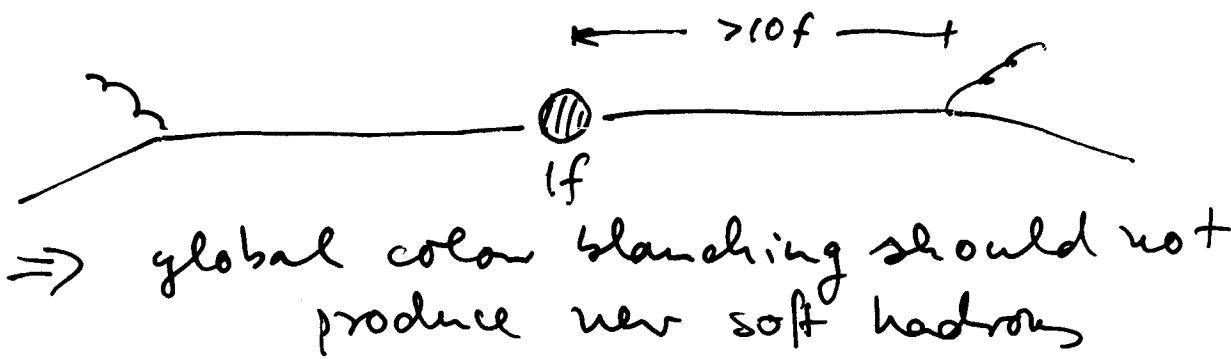
Puzzle

space-time picture



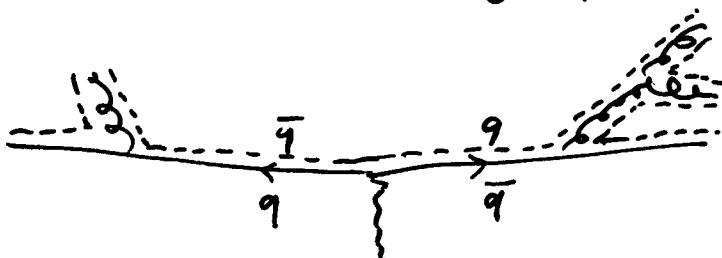
lifetime of quark: $\tau \sim \frac{1}{m_q} \cdot \frac{E}{m_q} \sim \frac{\hbar}{k_{\perp}^2}$
 average lifetime of quark in gap events: $\langle \frac{\hbar}{k_{\perp}^2} \rangle_{DLA} \sim 10-20 \text{ f}$ (for $y=3$)
↓ DLA, k_perp ...

First gluon emission far outside confinement region $R \sim \frac{1}{Q_0} \sim 1 \text{ f}$



Possible mechanisms

soft colour blanching by $q\bar{q}$ pairs with $K_{\perp} < Q_0$

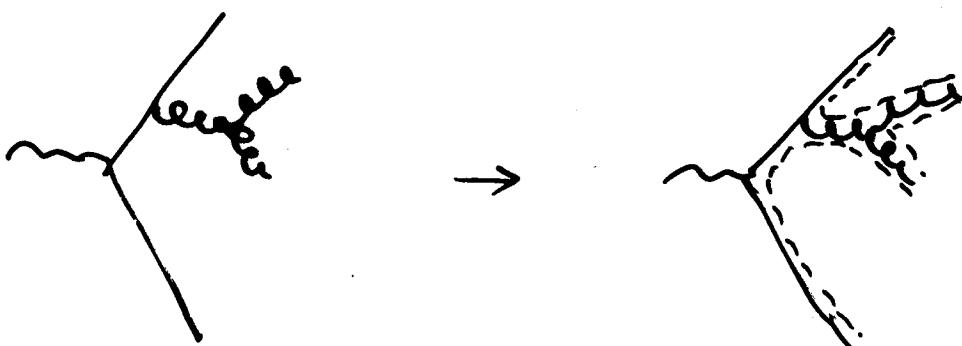


Dual hadron cascade
 Cohen-Tannoudji - w.o. '88

Summary

hadronic jets represented by QCD cascade with low cut off $Q_0 \gtrsim 1$.

- running coupling $\alpha_s \sim 1$
- soft gluon coherence } also for soft particles
- colour factors C_A, C_F
- * simplifying & analytic predictions
- * successes of perturbative picture
 - global multiplicities and moments
MLLA & complete solutions
almost quantitative
 - spectra - take into account
MLLA Kinematic corrections
 - angular distributions & correlations
DLA : at least qualitative
 - rapidity gaps in e^+e^-
quark travels $\sim 10 f$ before g emission
 \rightarrow soft confinement
 - outlook: MC for shapes
- * failure
 - soft multiparticle correlations (no Poisson limit)
- * Picture of soft hadronization



colour
neutralization
by soft
 $q\bar{q}$ pairs
(possibly
also gluons)