

Strings for Everybody

DESY, Hamburg

November 05

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Introduction: S/W-coupling dualities

Dualities (strong-weak coupling) have been
main theme of string theory in last decade.

Generalization of Kramers-Wannier duality
in 2D Ising model

non-perturbatively defined !

$$E = -\epsilon \sum s_i s_j$$

$$Z(\beta) \sim \sum_{\{s\}} e^{-\beta E} \sim \sum_{\{s\}} e^{-\beta^* E} \sim Z(\beta^*)$$

$$\beta^* = -\frac{1}{2} \log \tanh \beta$$

β small

$\leftrightarrow \beta^*$ large

Introduction: S/W-coupling dualities

2D Sine-Gordon model

Massive Thirring model

Boson Φ ; $V(\Phi) \sim \cos b \Phi$

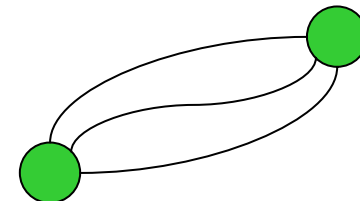
Fermion Ψ ; $V(\Psi) \sim g \Psi^4$

$$b^2 \sim 1/(1+g)$$

Is there dual to strongly coupled gauge theory ?

Proposal: Strongly coupled gauge theory is dual to weakly coupled strings.

Recall: QCD flux tubes \sim strings



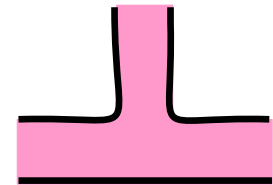
Intro: 't Hooft's analysis of large N

$$S(A) \sim 1/g^2 * \int \text{Tr} (\partial A \partial A + A^3 + \dots)$$

propagator $\sim g^2 \sim$

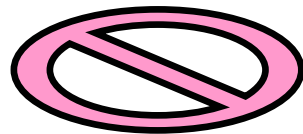


3-vertex $\sim 1/g^2 \sim$



$$Z^{\text{YM}} \sim \sum N^L g^{2E-2V} f(\Gamma) \sim \sum N^{V-E+L} \lambda^{E-V} f(\Gamma)$$

double line
networks Γ



$$\sim g^2 N^3 \sim \lambda N^2$$

$$\lambda = g^2 N$$

't Hooft coupling

$$\sim \sum N^{2-2h} \lambda^{E-V} f(\Gamma) \sim \sum_h (g_s)^{2h-2} F(X_\lambda) \sim Z^{\text{string}}$$

Network $\Gamma \rightarrow$ 2D surface S_Γ

$h = \#$ handles of S_Γ

$$g_s = 1/N$$

String
coupl.

Plan of talk

Part 0: Introduction (done)

Part I: Elements of modern string theory

**closed strings & SUGRA, p-branes,
branes & gauge theories, dualities.**

Part II: AdS/CFT duality – claim & tests

**N=4 SYM: scaling dimension & dilaton
operator, string in AdS, test of duality.**

Part III: Conclusions and Outlook

Part I: Closed strings and gravity

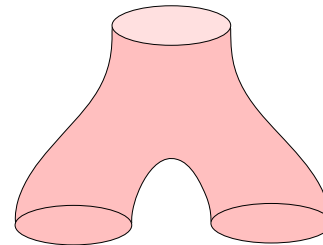
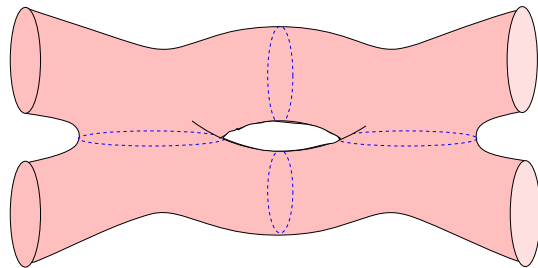
Closed string theory in background X

Infinite tower of vibrational modes $M \sim n/\ell_s$

String
length

String interaction through 3-vertex g_s

String
coupling



At low energies ($E \ll \ell_s$): massless closed str.

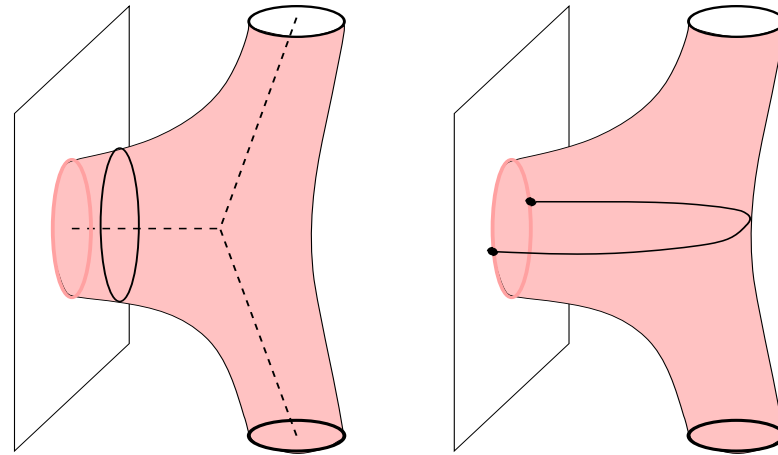
modes behave like gravitons \rightarrow 10D SUGRA

Part I: solitonic p- and D-branes

10-dim SUGRA has
solutions describing
massive & charged
objects localized on
 $p+1$ - dim. surfaces.

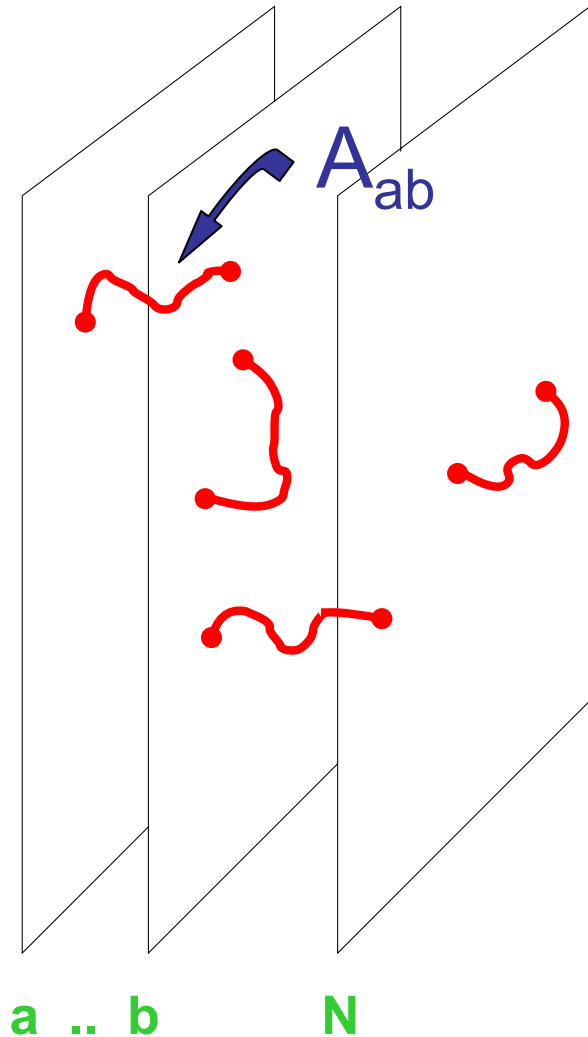
[black branes]

Branes in string theory ?



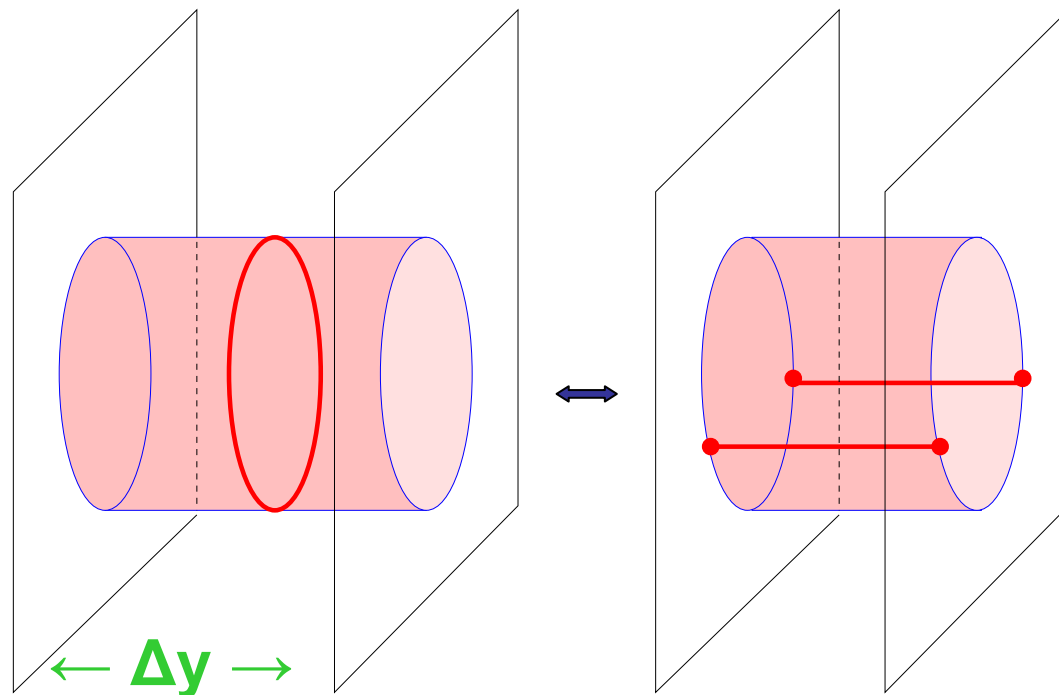
D-branes are objects on
which open strings can
end [Polchinski]

Part I: Open strings & gauge theory



At low energies ($E \ll \ell_s$):
massless modes of the
open string behave like
gauge bosons + matter
on the world-volume of
D-branes. **$p+1$ - dimensional
gauge theory**

Part I: Gauge-String theory dualities



Depending on Δy
one side simpler
 $\Delta y \ll \ell_s$: gauge th
 $\Delta y \gg \ell_s$: SUGRA

- Loop order is not preserved
- Closed string propagates in different dim. !

e.g. classical string \rightarrow
quantum gauge theory

Part II: The $\text{AdS}_5 / \text{CFT}_4$ duality

Gauge theory on a
stack of D3-branes




Strings in the near
horizon geometry

$\mathcal{N}=4$ 4D S Yang-Mills

U(N) gauge field A ; 6 scalars Φ
& their fermionic superpartners

4-dim conformal QFT

Closed str. in $\text{AdS}_5 \times S^5$


$$x_0^2 + x_1^2 - x_2^2 - x_3^2 - x_4^2 - x_5^2 = R^2$$

Curved backgrnd X_R

$$\lambda = g_{\text{YM}}^2 N \quad \sim \quad R^4 / \ell_s^4$$

Part II: Scaling dimension in CFT

Scaling dimension Δ characterizes behavior of fields under scale trasfos: $\Psi(ax) \rightarrow a^{-\Delta} \Psi(x)$




Class. theory: $A(x), \Phi(x), \dots \text{Tr}(\Phi_1(x)\Phi_2(x)\partial\Phi_1(x))$

$$\Delta = 1 \quad \dots \quad \Delta = 1+1+1+1 = 4$$

Quantum theory: Scaling dimension encoded

in

$$D = D_0 + \lambda D_1 + \lambda^2 D_2 + \dots$$

Dilaton operator diagonal + perturbative corrections

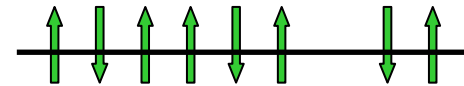
diag(Δ) Feynman diagrams

Degenerate
Perturbation
theory

Part II: Dilaton operator of SYM

Space of (single trace) composite fields \sim
state space of one-dimensional spin chain :

e.g. $\text{Tr} (Z \bar{Z} Z Z \bar{Z} Z \dots \bar{Z} Z)$



$$Z = \Phi_1 + i\Phi_2, \quad \bar{Z} = \Phi_1 - i\Phi_2$$

$n = \# \text{ of spins}$

D_1 acts like Hamiltonian of spin chain model
(e.g. of Heisenberg model) on fields in SYM

$$D_1 \sim H \sim \sum \sigma_i \sigma_{i+1}$$

Diagonalizable for small n
& in the limit $n \gg 1$ (TBA)

Part II: Spectrum for string in AdS

For a particle moving freely in some background X
→ spectrum of Laplace operator $H = \Delta_X$

e.g. real line: $H = -\partial^2 \Rightarrow E = p^2$ **AdS harder (bound states)**

Analogous infinite dimensional problem for a
string in AdS not solved (spectrum unknown)

String spectrum known • $R/\ell_s \gg 1$ (flat space)

• semi-classical limit (large quantum numbers)

e.g. rotating strings $\Rightarrow E = E_J(R)$ $J \gg 1$ **spin J**

Part II: Tests of AdS/CFT duality

Spectrum in AdS \leftrightarrow Scaling dim. in SYM

Semi-classical limit: String states with large J

\leftrightarrow long composite operators in gauge theory.

$$E_J(\rho) \sim E_J^{(0)} + E_J^{(1)} \rho + \dots \leftrightarrow D \sim D_0 + D_1 \lambda + \dots$$

$\rho = R/\ell_s$

D_1 acts like Hamiltonian of spin chain model

(e.g. of Heisenberg model) on fields in SYM

classical spinning strings \leftrightarrow diag. of spin chain H (TBA)

Part III: Outlook and conclusions

AdS/CFT dualities generalize to field theories

- in other dimensions (e.g. $\text{AdS}_3/\text{CFT}_2$) **solvable ?**
- with less symmetry (less SUSY; w/o scale inv.)
 $S^5 \rightarrow X^5$; confining theories

Open problem: solution of perturbative string theory in AdS spaces (spectrum, interaction).

\leftrightarrow 2D σ -model on super spaces \leftrightarrow **cond mat.**

$$S_{\text{Pol}}(X) \sim \int dx dt G_{ab}(X) \partial X^a \partial X^b \quad \text{AdS}_3 \sim \text{PSU}(1,1|2) \dots$$

Part III: Further aspects of duality

Tempting to read duality in opposite direction:

Gauge theory \rightarrow stringy extension of gravity

cosmological back-
grounds, black holes

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