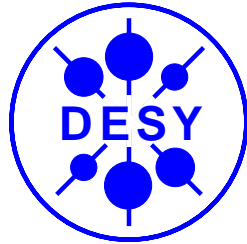


Searches for Neutral Higgs Bosons at LEP

Alexei Raspereza, DESY-Zeuthen



DESY Seminar, 30.10.2001



- Standard Model and Higgs Mechanism
- LEP Performance
- Higgs Boson at LEP (Production and Decays)
- Analysis Procedures
- Search for SM Higgs Boson
- Flavour Independent Search for hZ process
- Search for Higgs Bosons of MSSM
- Conclusion

Milestones of Standard Model Triumph

- Observation of gluons by TASSO at PETRA, DESY (1980)
- Discovery of W^\pm and Z by UA1 and UA2 at $p\bar{p}$ collider, CERN (1983)
- SM successfully passed precision tests at LEP, HERA, SLAC, TEVATRON
- Discovery of top-quark at TEVATRON $p\bar{p}$ collider, Fermilab (1995)

What is missing ?

3 fermion families

5 representations of $U(1)_Y \times SU(2)_L$ group

$$\begin{pmatrix} \nu \\ \ell \end{pmatrix}_L \quad \begin{pmatrix} u \\ d \end{pmatrix}_L \quad \ell_R \quad u_R \quad d_R$$

Gauge Invariance

Kinetic fermion term

$\mathcal{L}_f = i\bar{\Psi}\gamma^\mu\partial_\mu\Psi$ should be invariant under

$$\Psi \longrightarrow U(\mathbf{x})\Psi = (1 + Y\alpha(\mathbf{x}) + \vec{\beta}(\mathbf{x})\vec{\sigma})\Psi \implies$$

introduction of B_μ, \vec{W}_μ

$$\partial_\mu \rightarrow D_\mu = \partial_\mu + ig_1 Y B_\mu + ig_2 \vec{\sigma} \vec{W}_\mu$$

Gauge Boson transformation rules

$$B_\mu \rightarrow U(\mathbf{x})B_\mu U^{-1}(\mathbf{x}) + \frac{i}{g_1} \left[\partial_\mu U(\mathbf{x}) \right] U^{-1}(\mathbf{x})$$

$$W_\mu^i \rightarrow U(\mathbf{x})W_\mu^i U^{-1}(\mathbf{x}) + \frac{i}{g_2} \left[\partial_\mu U(\mathbf{x}) \right] U^{-1}(\mathbf{x})$$

Boson kinetic term

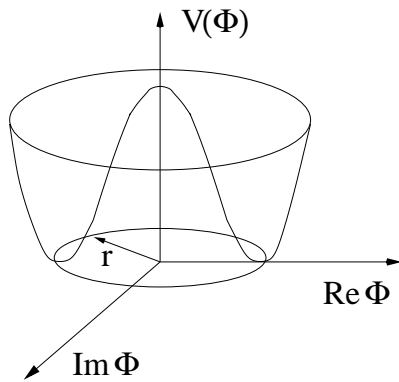
$$\mathcal{L}_B = -\frac{1}{4}(\vec{W}_{\mu\nu}\vec{W}^{\mu\nu} + B_{\mu\nu}B^{\mu\nu})$$

$$\vec{W}_{\mu\nu} = \partial_\mu \vec{W}_\nu - \partial_\nu \vec{W}_\mu + g_2 \vec{W}_\mu \times \vec{W}_\nu$$

$$B_{\mu\nu} = \partial_\mu B_\nu - \partial_\nu B_\mu$$

Gauge invariance forbids mass terms

$$m\bar{\Psi}\Psi \quad \frac{1}{2}m^2 B_\mu B^\mu \quad \frac{1}{2}m^2 \vec{W}_\mu \vec{W}^\mu$$



$$\Phi = \begin{pmatrix} \Phi^+ \\ \Phi^0 \end{pmatrix}$$

$$\mathcal{L}_H = \partial_\mu \Phi^\dagger \partial^\mu \Phi - V(\Phi^\dagger \Phi)$$

$$V(\Phi^\dagger \Phi) = \lambda \left(\Phi^\dagger \Phi - \frac{\mu^2}{2\lambda} \right)^2$$

Vacuum state: $\Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ \nu \end{pmatrix}, \quad \nu = \mu/\sqrt{\lambda}$

Physical state: $\Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ \nu + h(x) \end{pmatrix}$

$\mathcal{L}_H =$

$$\begin{array}{l} \frac{1}{2} \partial_\mu h \partial^\mu h \quad \} \text{propagation term} \\ -\lambda \nu^2 h^2 \quad \} \text{Higgs mass term} \\ -\lambda \nu h^3 \quad \} \\ -\frac{\lambda}{4} h^4 \quad \} \text{Higgs self-interaction} \end{array}$$

$$m_h = \sqrt{2\lambda\nu}$$

Boson masses

$$D_\mu \Phi^\dagger D^\mu \Phi \implies$$
$$\left. \begin{aligned} & \frac{1}{8} (g_2 \nu)^2 W_\mu^+ W^{-\mu} \\ & \frac{1}{8} \nu^2 (g_1^2 + g_2^2) Z_\mu Z^\mu \\ & \frac{1}{4} g_2^2 \nu W_\mu^+ W^{-\mu} h \\ & \frac{1}{4} \nu (g_1^2 + g_2^2) Z_\mu Z^\mu h \end{aligned} \right\} \begin{array}{l} \text{W and Z mass terms} \\ \text{W and Z interaction with h} \end{array}$$

$$W^\pm = W^1 \pm iW^2$$
$$\begin{pmatrix} A \\ Z \end{pmatrix} = \begin{pmatrix} \cos \theta_w & \sin \theta_w \\ -\sin \theta_w & \cos \theta_w \end{pmatrix} \begin{pmatrix} B \\ W^3 \end{pmatrix}$$

$$m_W = \frac{1}{2} \nu g_2$$

$$m_Z = \frac{1}{2} \nu \sqrt{g_1^2 + g_2^2}$$

$$\cos \theta_w = \frac{m_W}{m_Z}$$

Fermion masses

$$\mathcal{L}_Y = Y^u \bar{Q}_L \Phi^* u_R + Y^d \bar{Q}_L \Phi d_R + Y^\ell \bar{L}_L \Phi \ell_R + \text{h.c.} \implies$$

$$\left. \begin{aligned} & \frac{Y^f}{\sqrt{2}} (\bar{\Psi}_L \Psi_R + \bar{\Psi}_R \Psi_L) \nu \\ & \frac{Y^f}{\sqrt{2}} (\bar{\Psi}_L \Psi_R + \bar{\Psi}_R \Psi_L) h \end{aligned} \right\} \begin{array}{l} \text{Fermion mass terms} \\ \text{Fermion interaction with h} \end{array}$$

$$m_f = \frac{Y^f \nu}{\sqrt{2}}$$

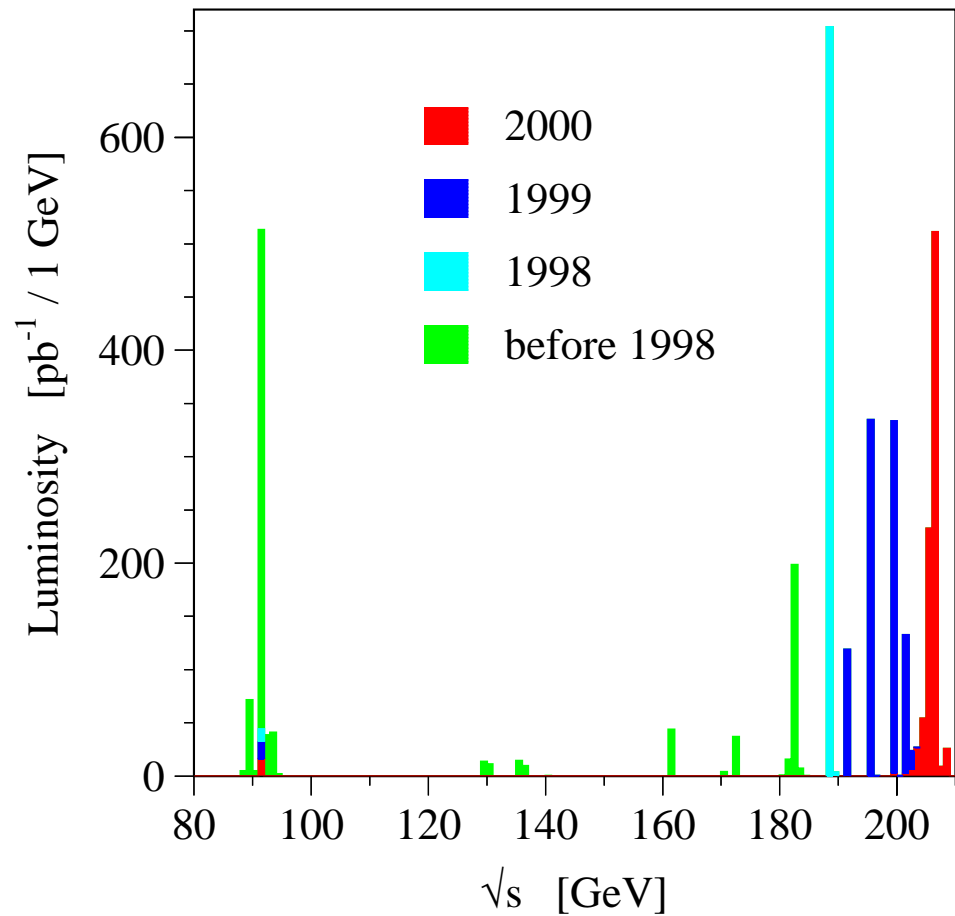
The running of LEP

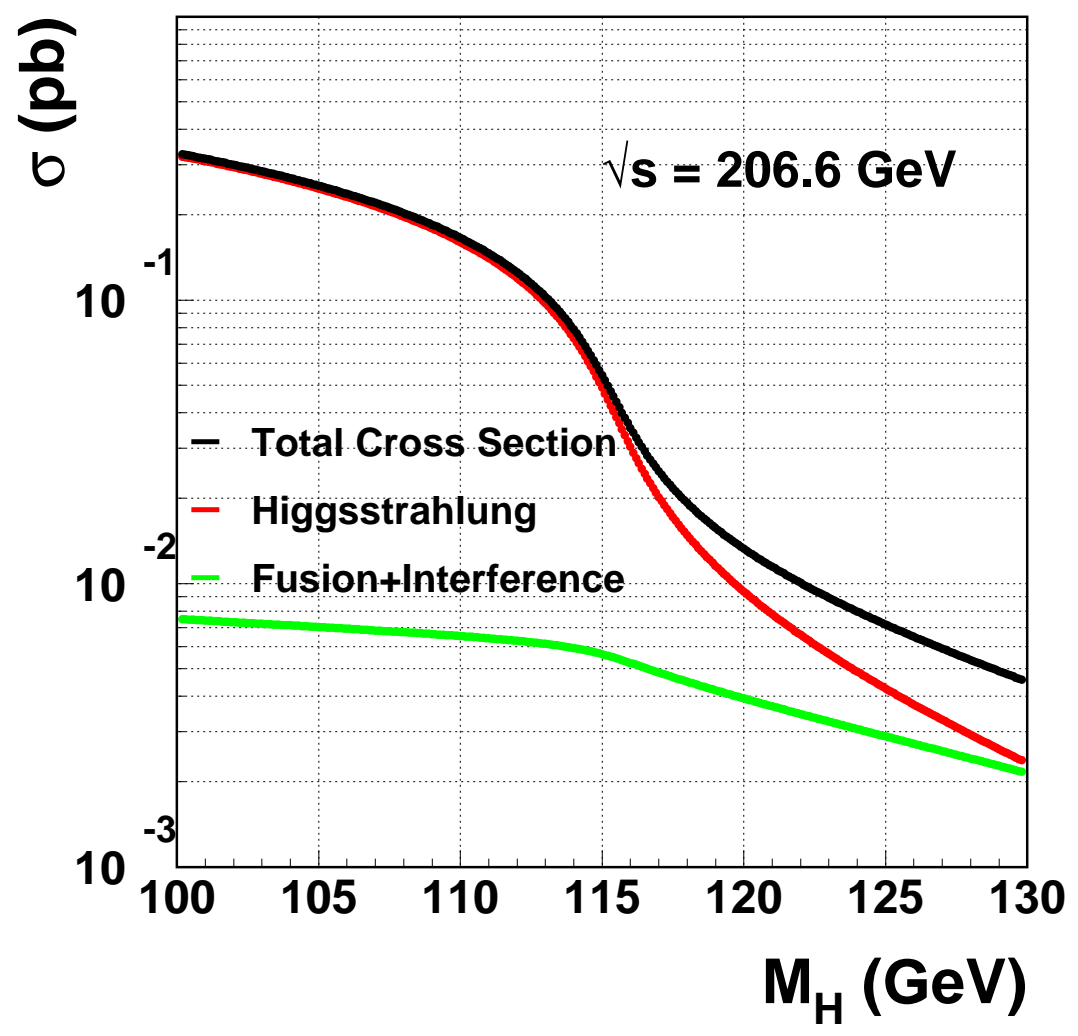
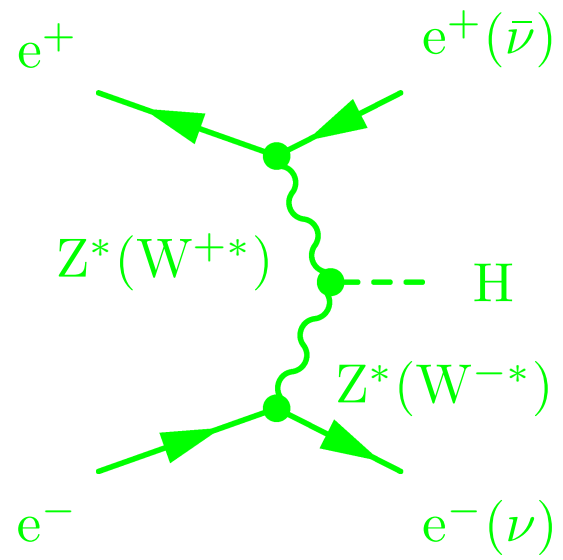
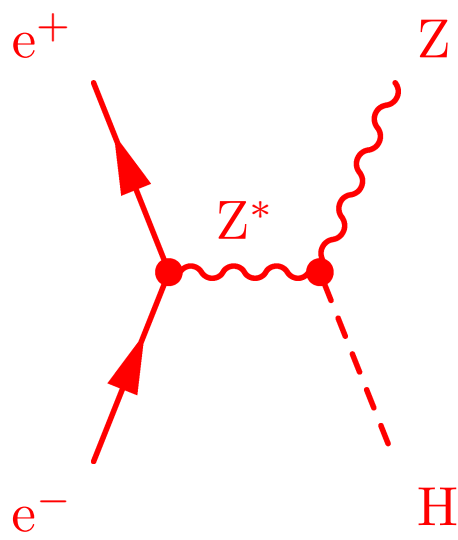
In 12 years of running, LEP has produced e^+e^- collisions in four interaction points:

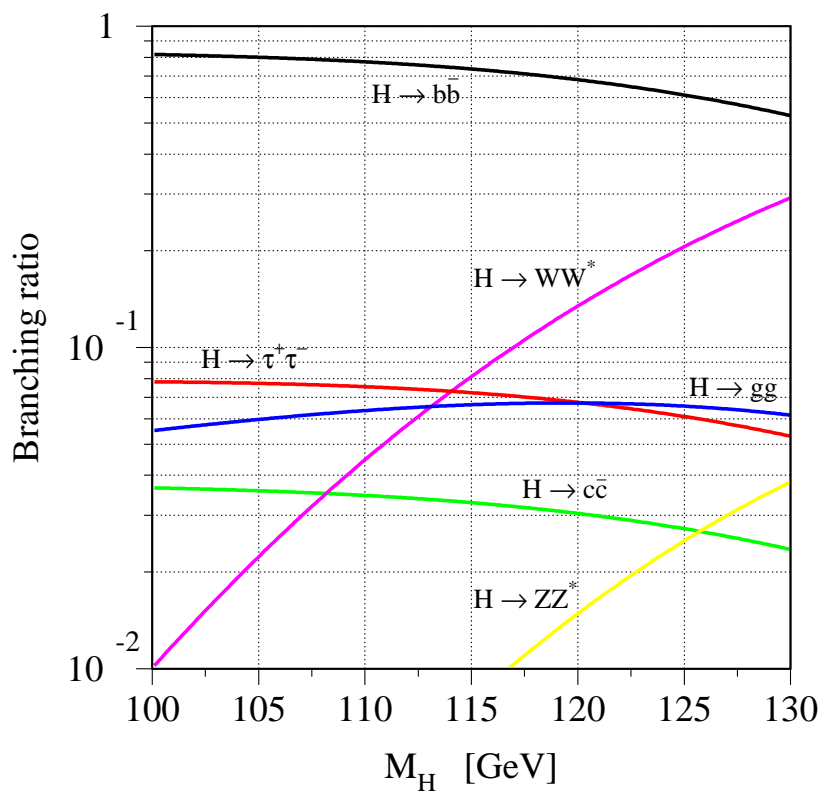
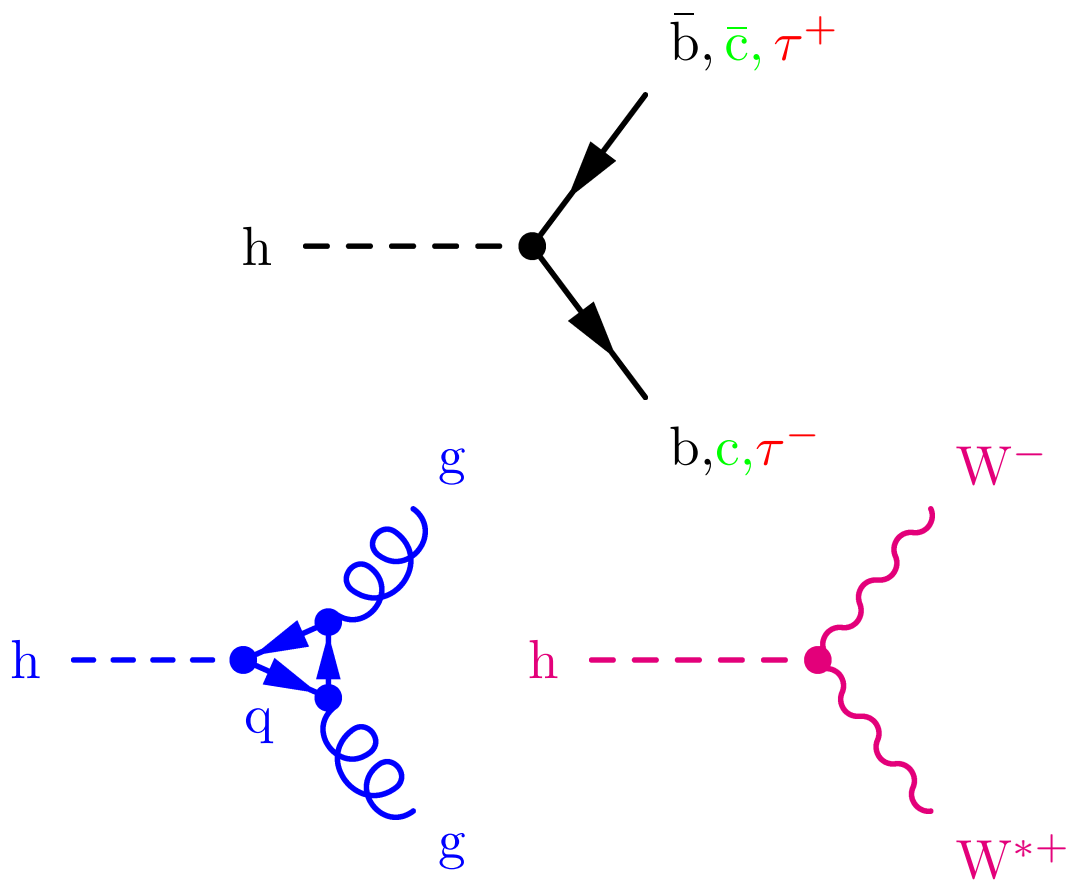
ALEPH, DELPHI, L3 and OPAL

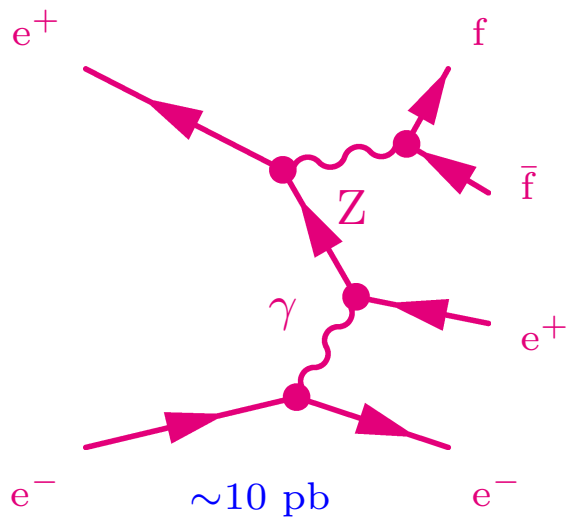
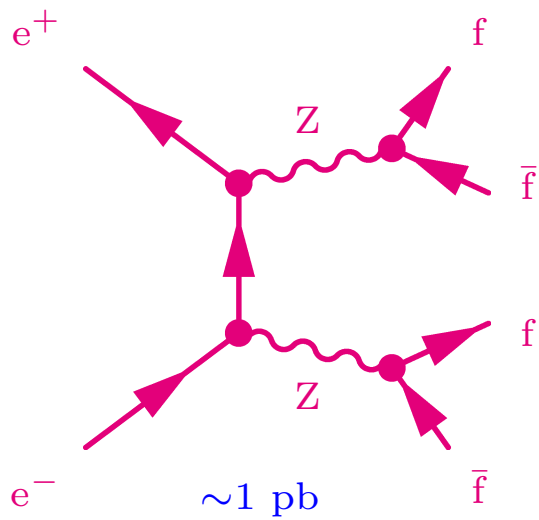
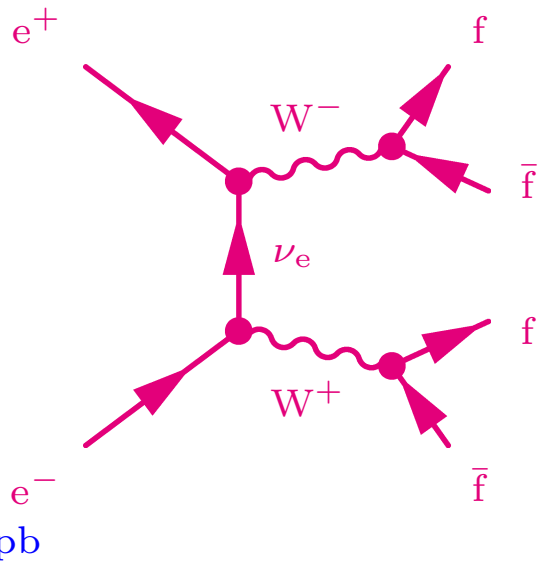
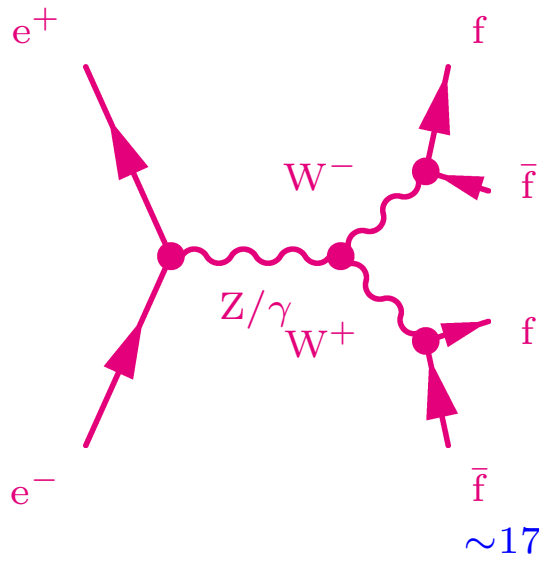
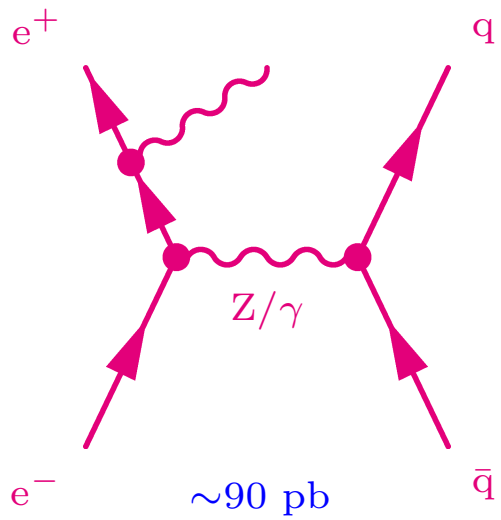
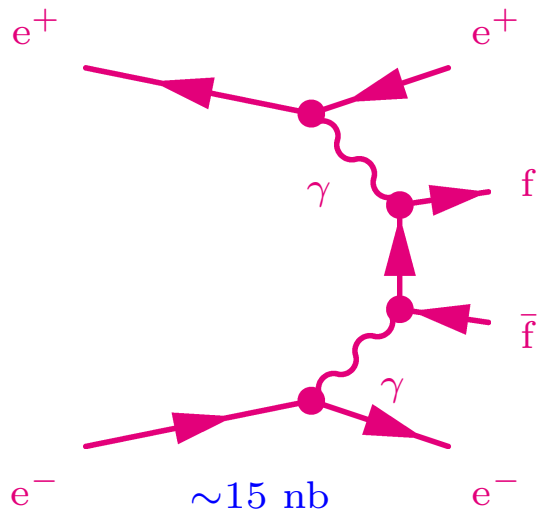
in a wide range of centre-of-mass energies:

$$\sqrt{s} \simeq 90 - 210 \text{ GeV}$$

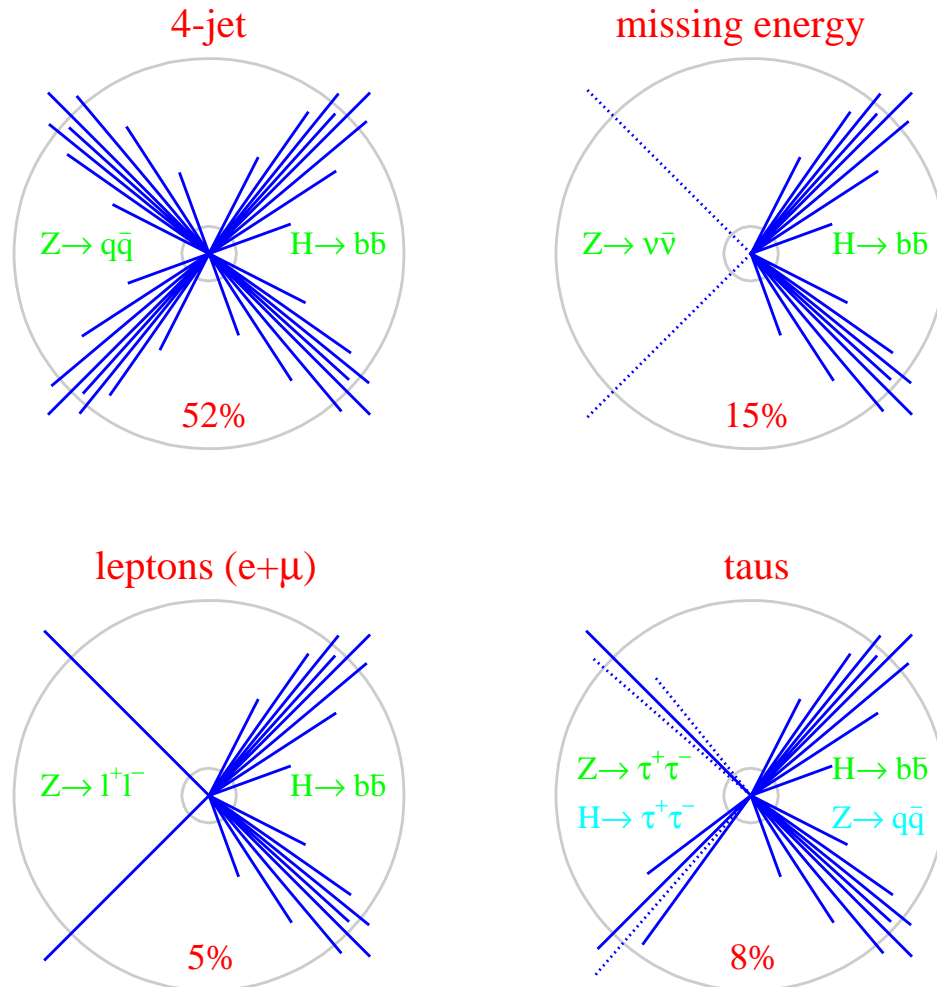








About 80% of the final states exploited

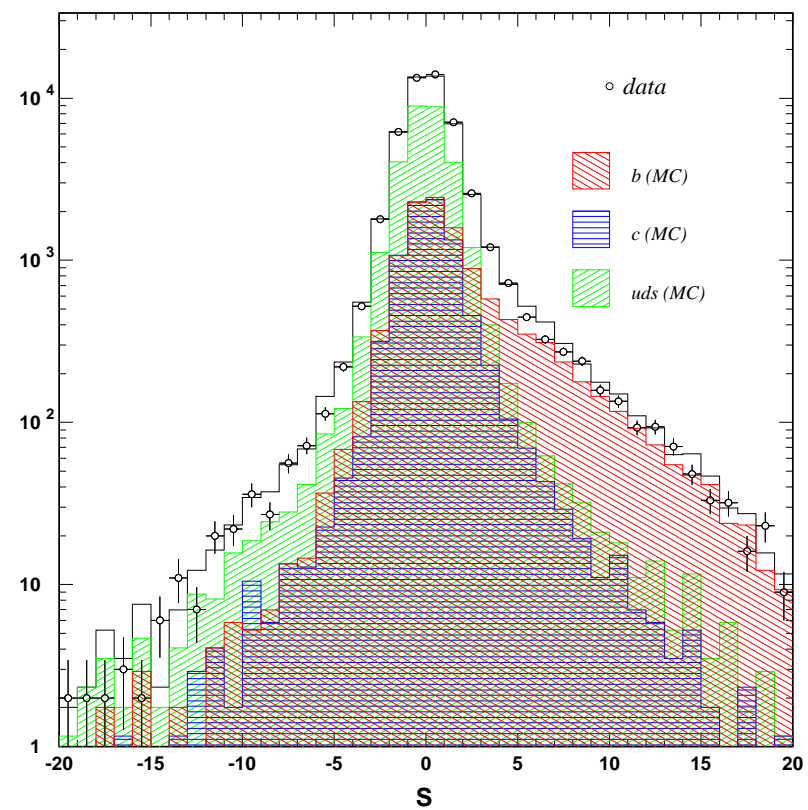
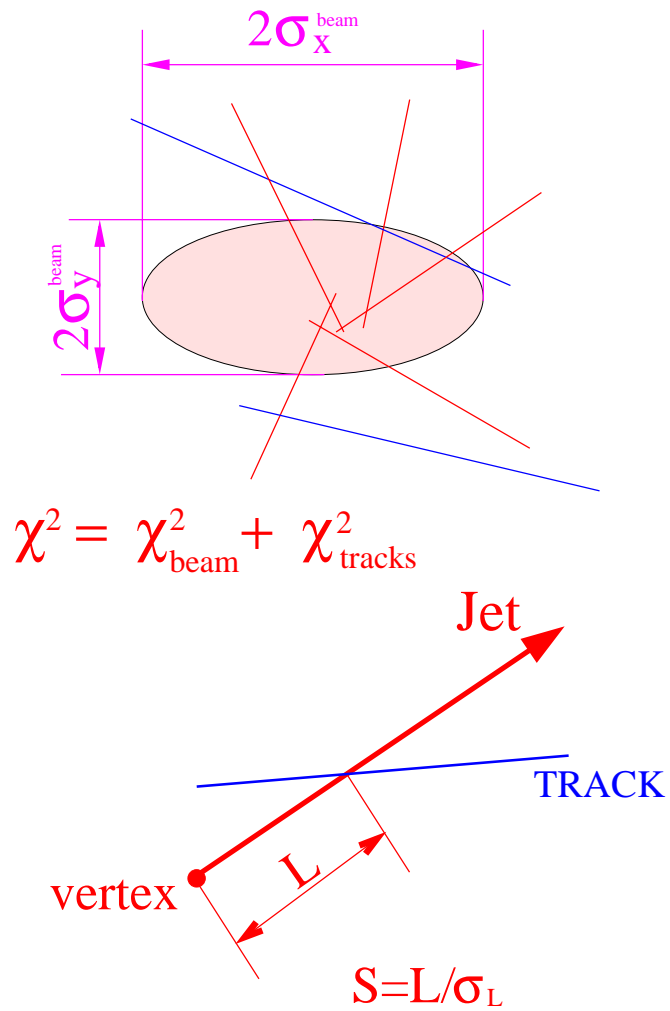


Expected signal events for a $m_H = 115$ GeV Higgs:

4-jet	= 6.5
missing energy	= 1.9
taus	= 1.0
leptons (e+μ)	= 0.6
All channels	= 10

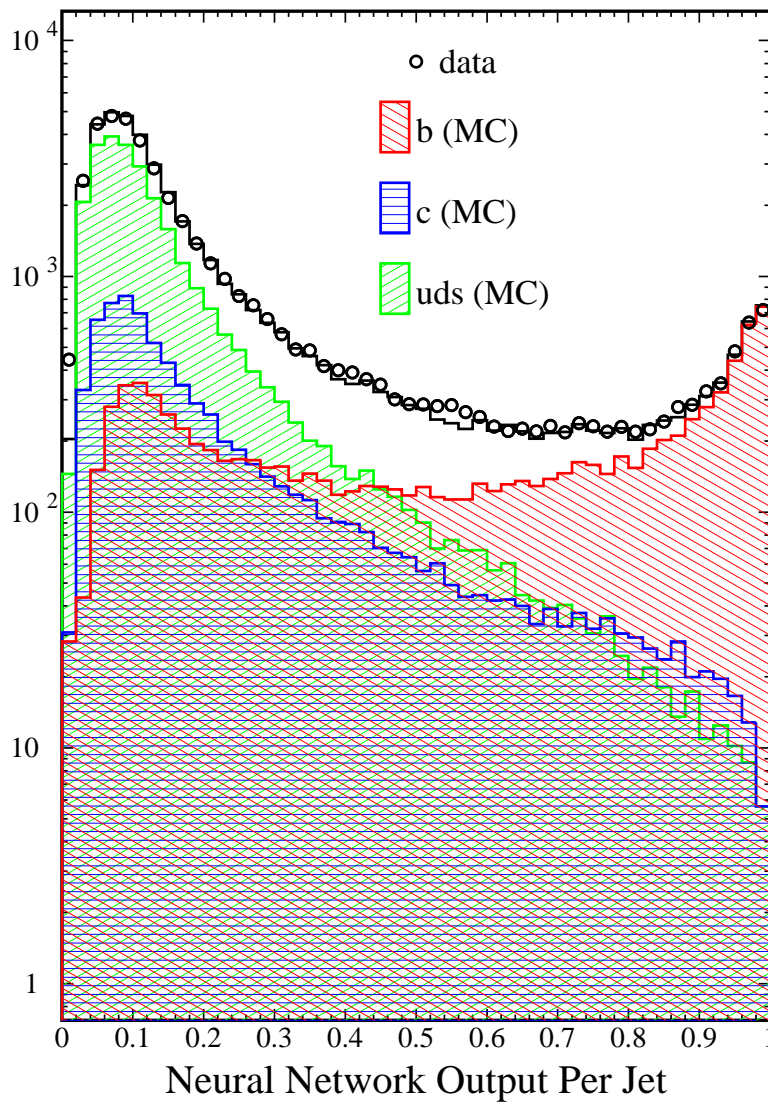
b-jet features

- $\tau_b \sim 1.6$ ps;
 $\tau_c = 0.2 - 1.0$ ps, $\tau_\tau \sim 0.3$ ps;
 $\tau(K_S^0) \sim 89$ ps, $\tau(\Lambda) \sim 260$ ps
- secondary vertexes
- large multiplicity
- (may be) accompanying leptons
- higher masses (smaller boosts)



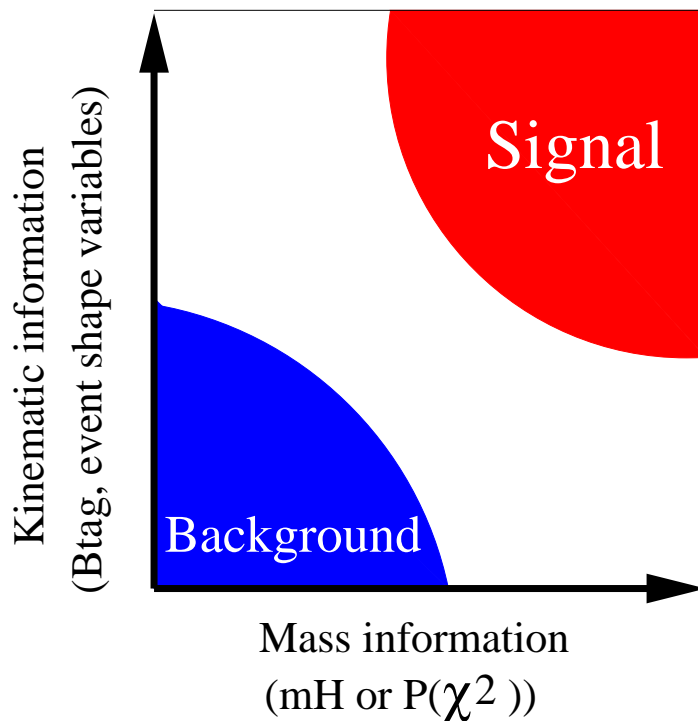
decay length information
 jet multiplicity
 jet shape variables
 information about prompt leptons
 information about secondary vertexes

\Rightarrow NNet



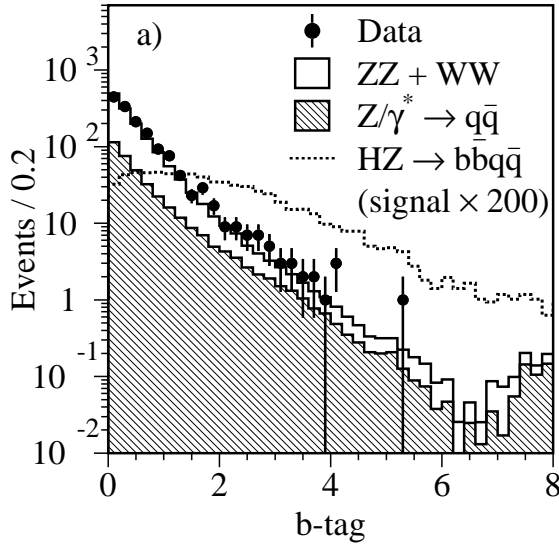
Analysis steps

- preselection of high multiplicity hadronic events
- fine-tuned signal-to-background discrimination based on NNet, Likelihood or cut-based technique
- kinematic fit \implies access to the mass information
 - 4C fit (energy-momentum conservation)
 - 5C fit ($4C \oplus m_{qq, \ell\ell} = m_Z$)
 - 6C fit ($4C \oplus m_{qq, \ell\ell} = m_Z \oplus m_{q'q'} = m_H$)missing energy and tau channels: neutrino and tau momenta are undefined \rightarrow fits with less number of constraints
- construction of final discriminant

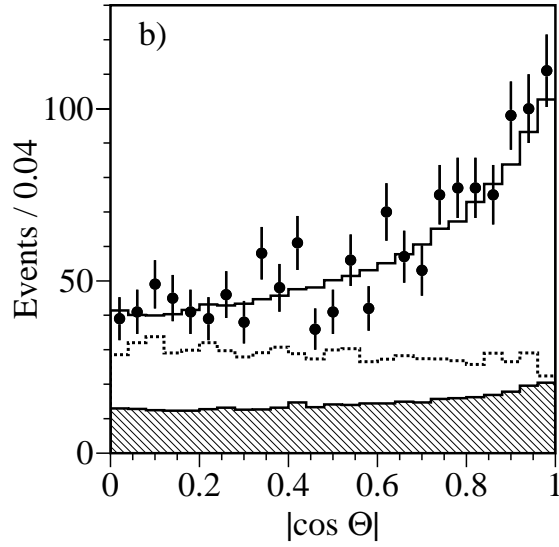


Signatures: high multiplicity, 4jet topology, full energy

event btag

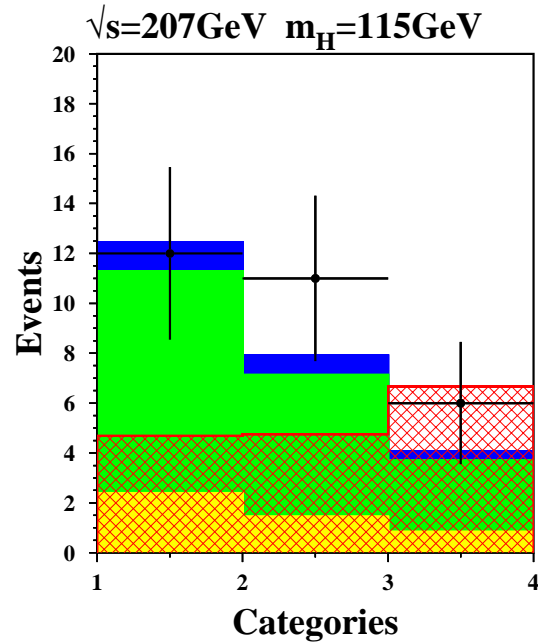
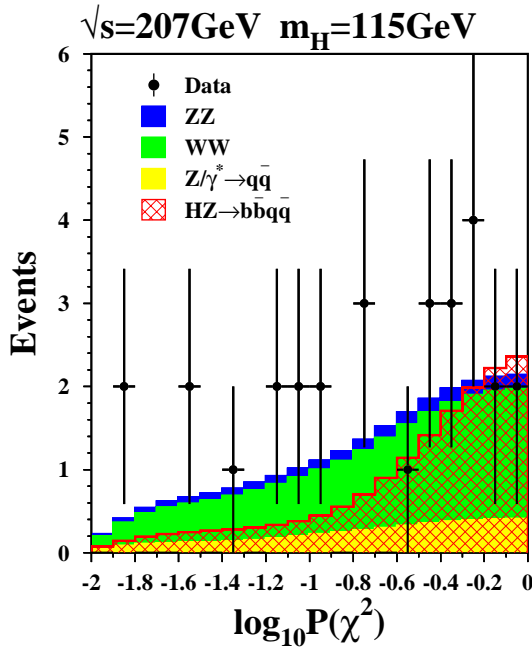


angular information



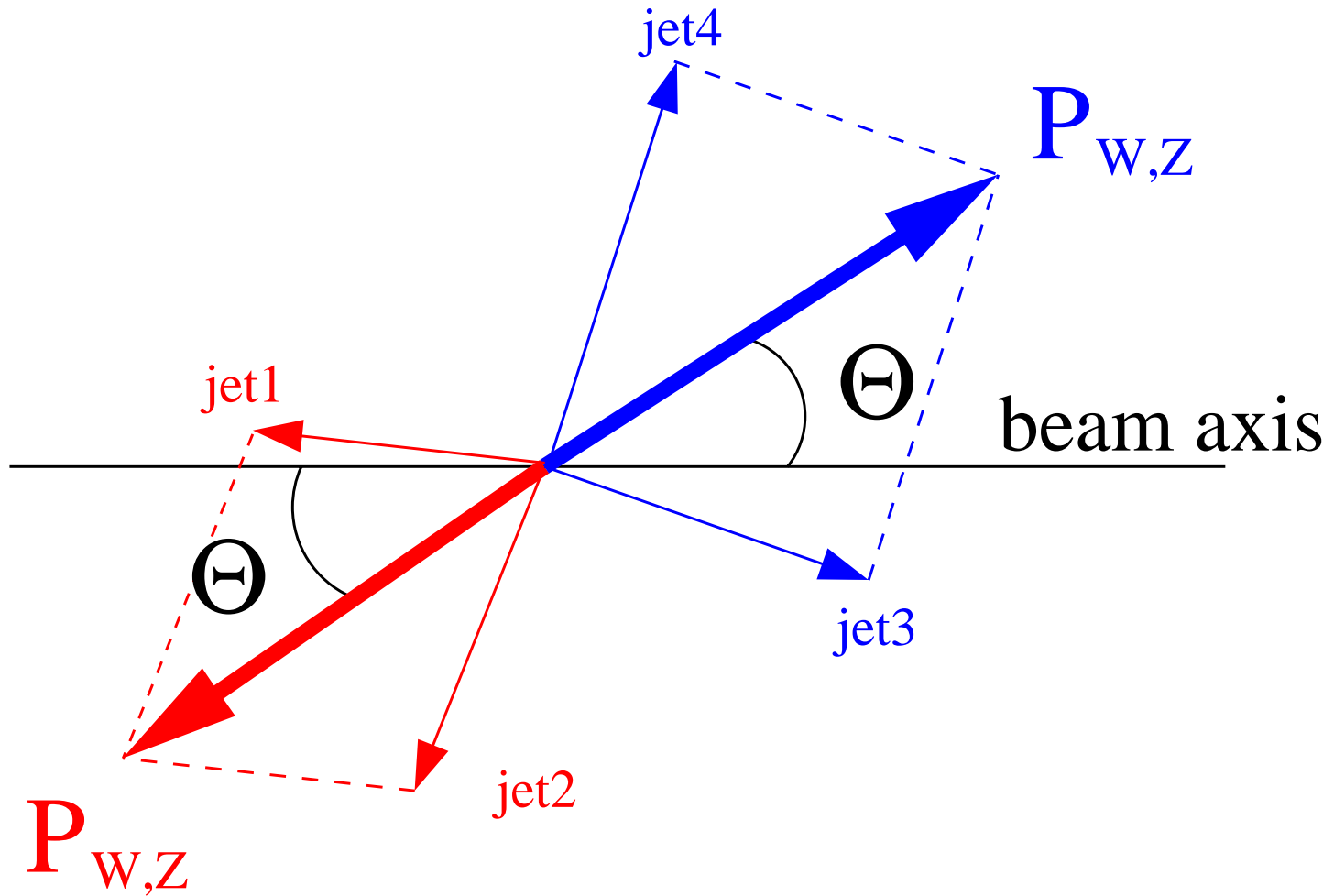
mass information in L3

$$\chi^2 = \min \left\{ \left(\frac{\Sigma_i^{4C} - (m_H + m_Z)}{\sigma_\Sigma} \right)^2 + \left(\frac{\Delta_i^{4C} - |m_H - m_Z|}{\sigma_\Delta} \right)^2 \right\}$$

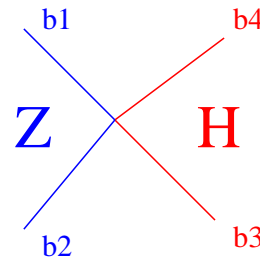
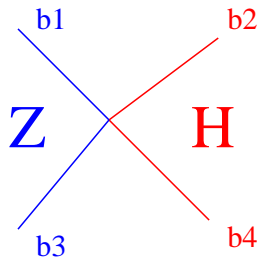
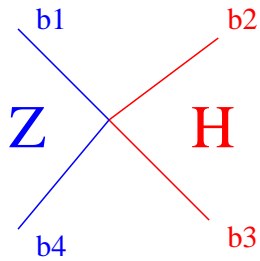


Mass variable, btag, category \rightarrow Likelihood

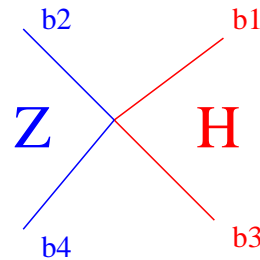
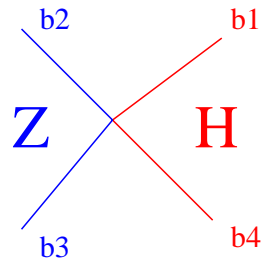
Angular information in L3 Hqq analysis



Jet tag based event category in L3 Hqq analysis



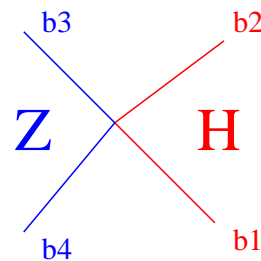
Category 1



Category 2

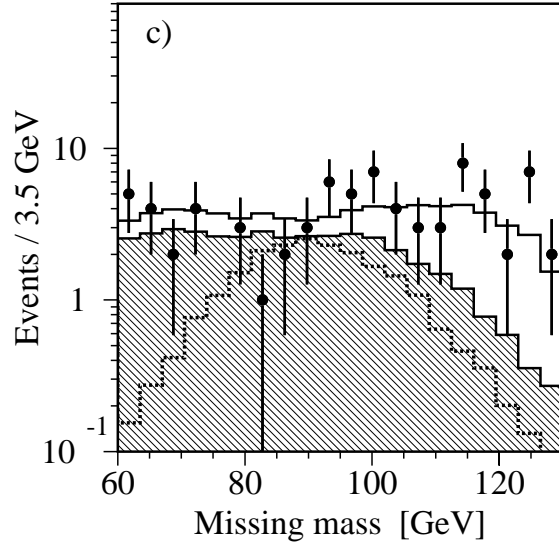
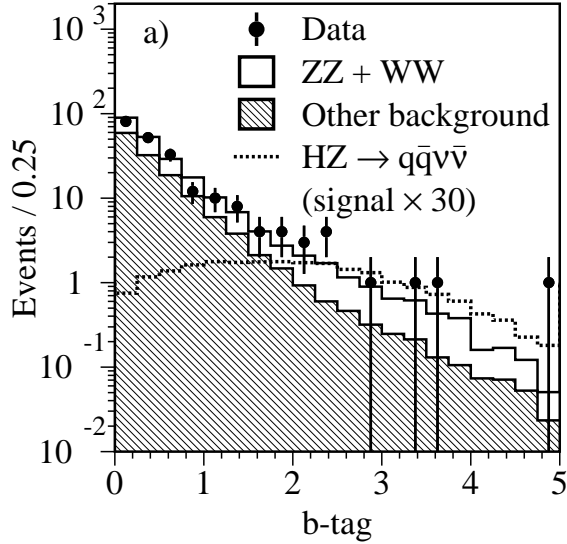
$$b1 > b2 > b3 > b4$$

b_i - btag value of jet i

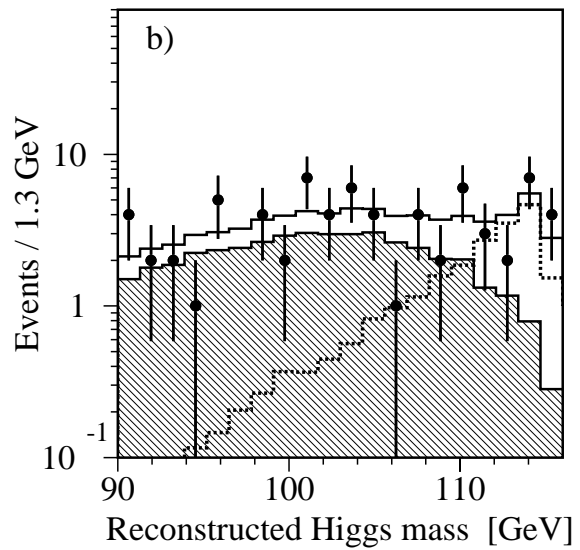
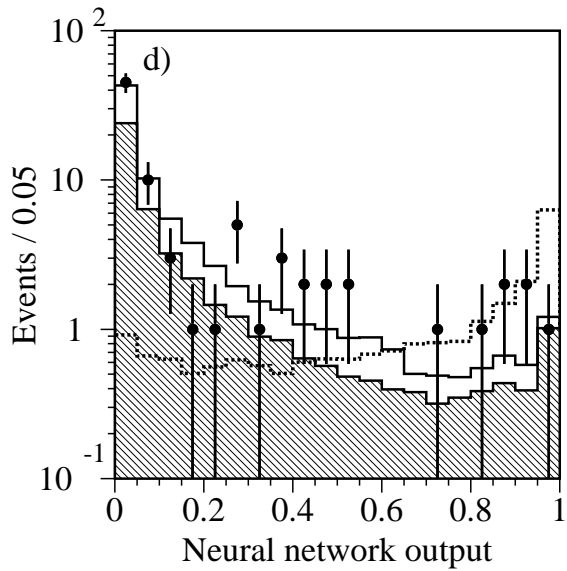


Category 3

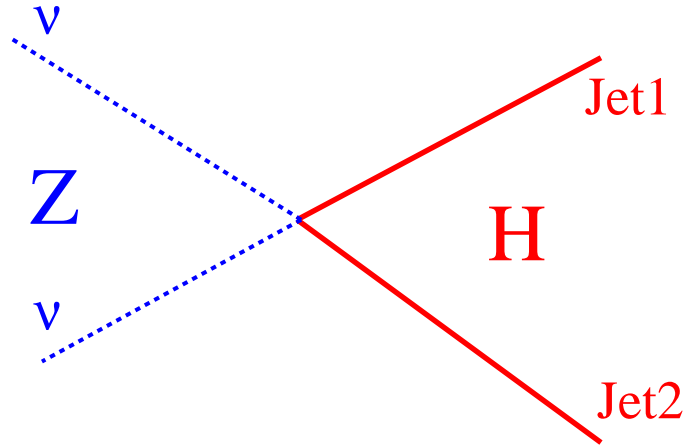
Signatures: 2 jets, missing energy, $M_{\text{mis}} = M_Z$



Kinematics \rightarrow NNet



NNet output + m_H^{rec} \rightarrow Final discriminant



$$E_{\text{mis}} = \sqrt{s} - E_{\text{jet1}} - E_{\text{jet2}}$$

$$\vec{P}_{\text{mis}} = -\vec{P}_{\text{jet1}} - \vec{P}_{\text{jet2}}$$

$$M_{\text{mis}} = \sqrt{E_{\text{mis}}^2 - P_{\text{mis}}^2}$$

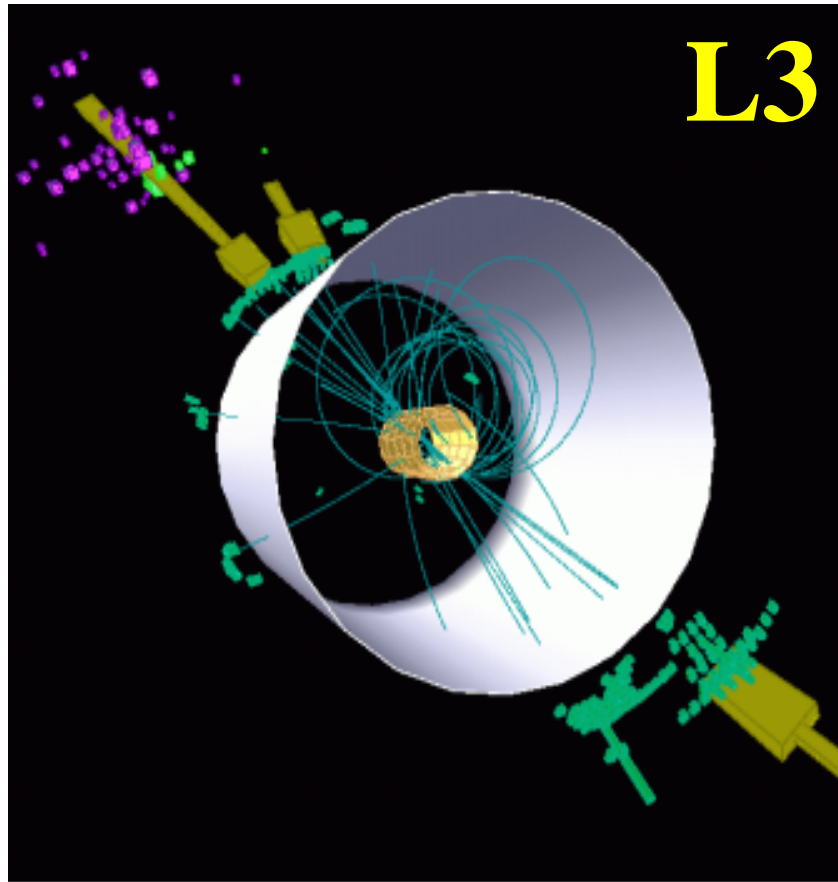
Constraint : $M_{\text{mis}} = M_Z$

Kinematic fit

↓

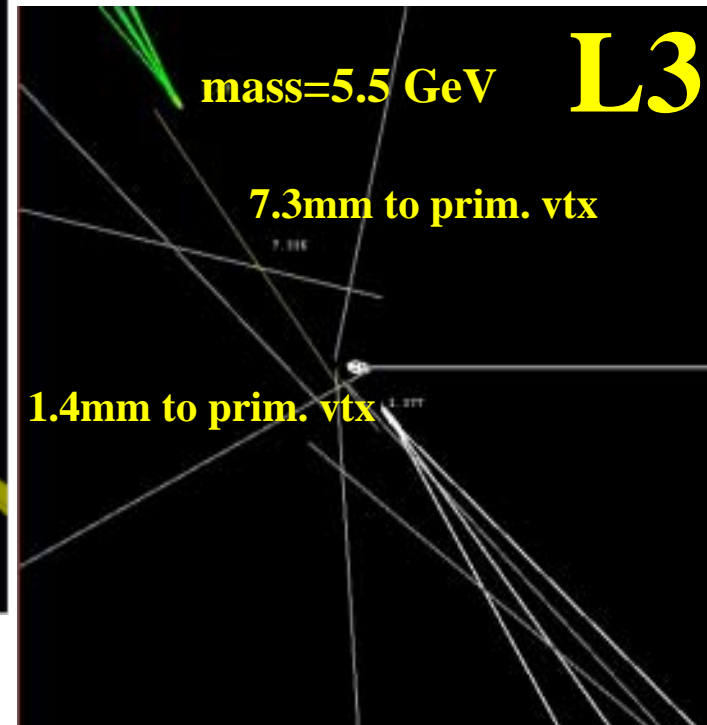
$$M_{\text{H}}^{\text{rec}} = \sqrt{(E_{\text{jet1}} + E_{\text{jet2}})^2 - (\vec{P}_{\text{jet1}} + \vec{P}_{\text{jet2}})^2}$$

most significant H $\nu\nu$ candidate

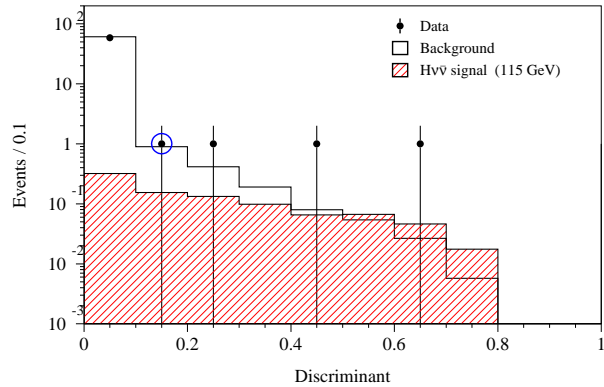


measured H mass=114 GeV
H mass resolution ~3 GeV

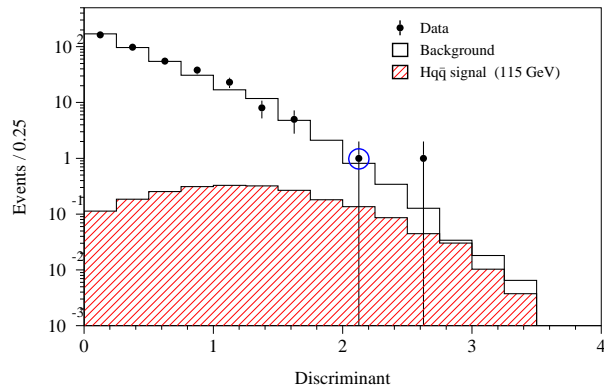
Secondary vtx's view



Combination of channels

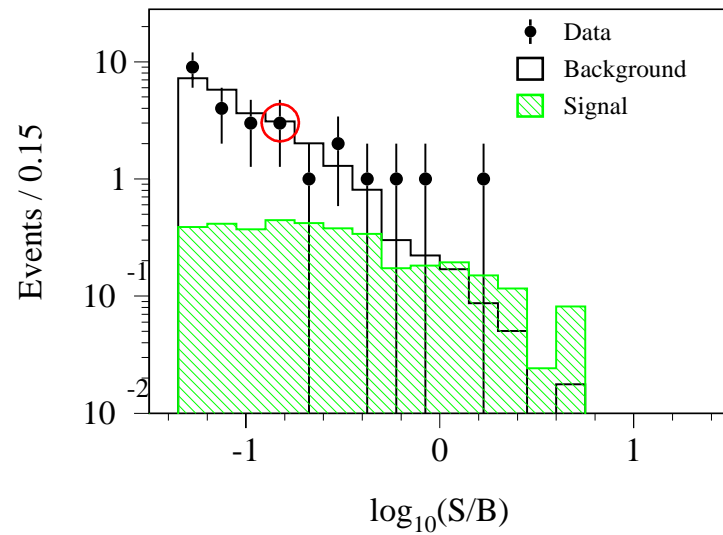


Combine channels with equal s/b ratio



LEP treated as a single analysis:

$$4 \text{ experiments} \times 4 \text{ channels} \times 15 \text{ beam energies} = 240 \text{ analyses combined}$$



The shape of the distributions depends on the Higgs mass hypothesis !

Statistical Method: $-2\ln(Q)$

Likelihood Ratio test-statistic:

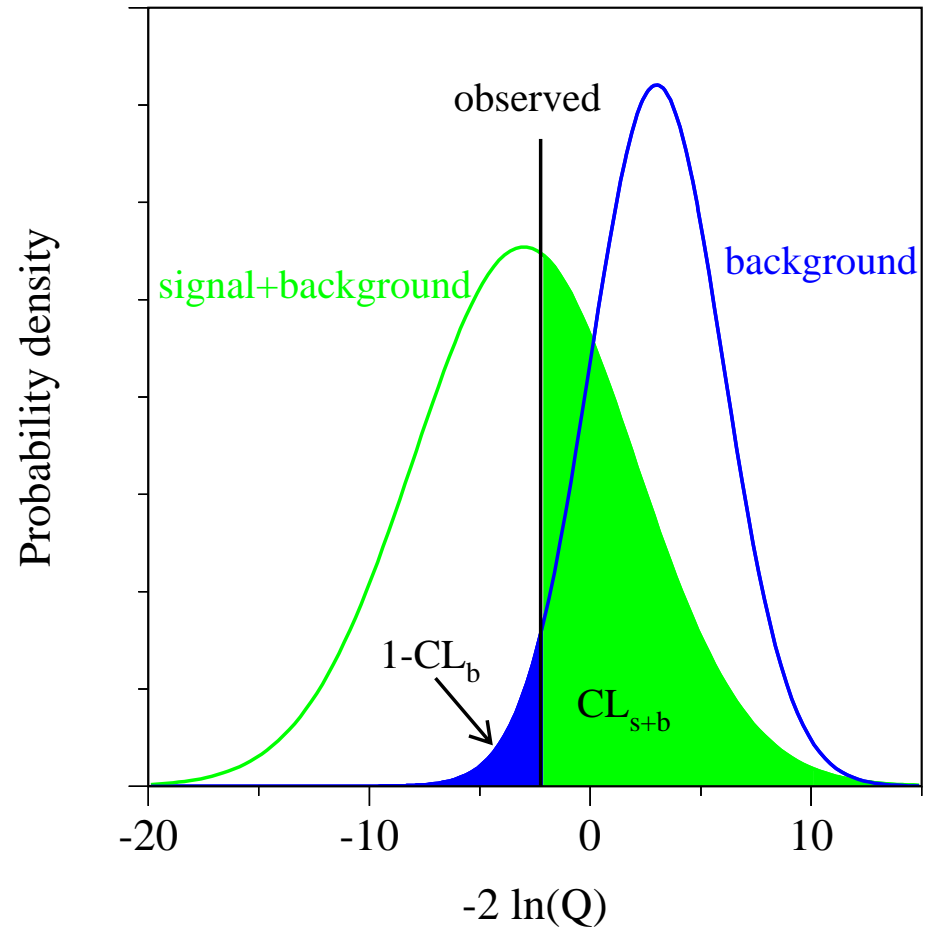
$$Q = \frac{\mathcal{L}(s + b)}{\mathcal{L}(b)}$$

Each bin (i) in the final variable is treated as a **Poisson counting experiment:**

$$\ln(Q) = -s_{tot} + \sum_{i=1}^N n_i \ln \left(1 + \frac{s_i}{b_i} \right)$$

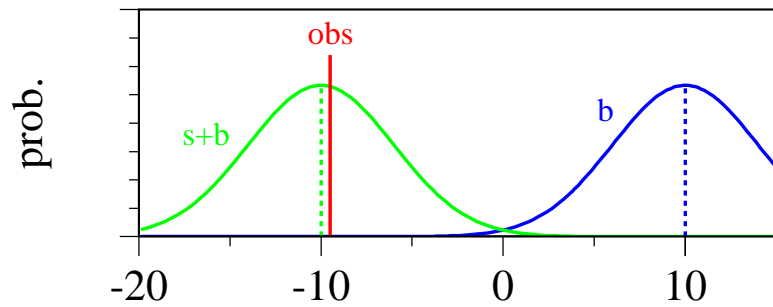
In the high statistics limit:

$$-2 \ln(Q) \rightarrow \Delta\chi^2$$

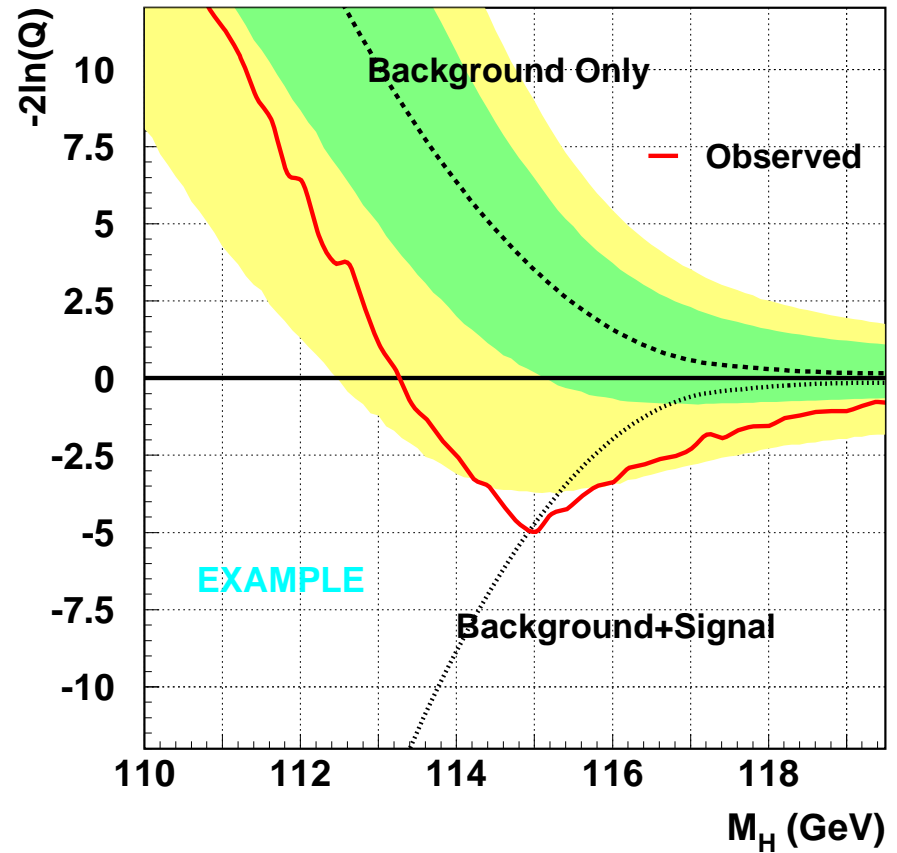
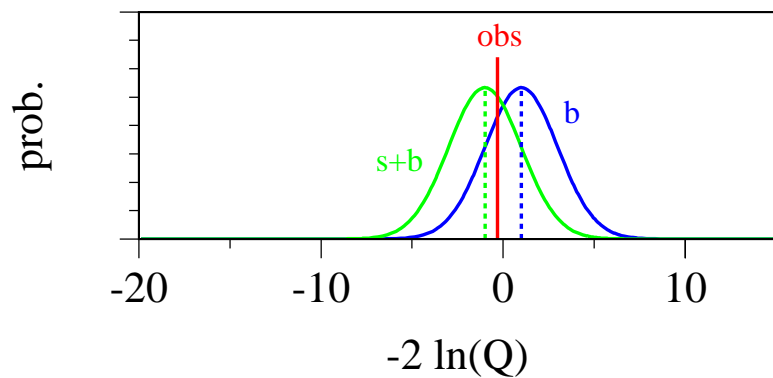


More on the Likelihood Ratio

Low mass: high sensitivity

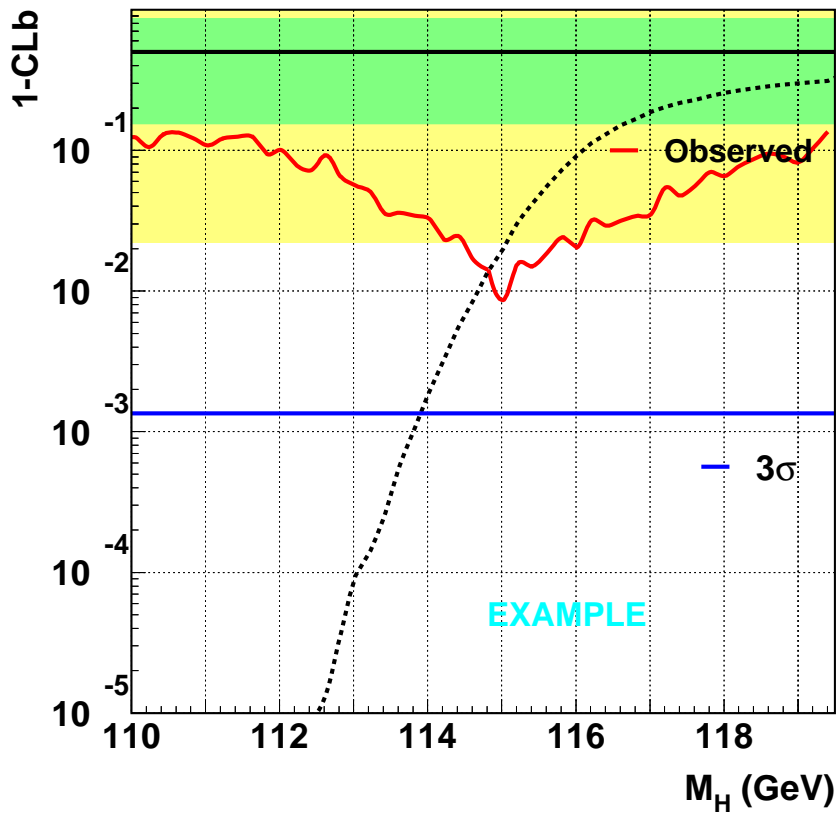


High mass: low sensitivity

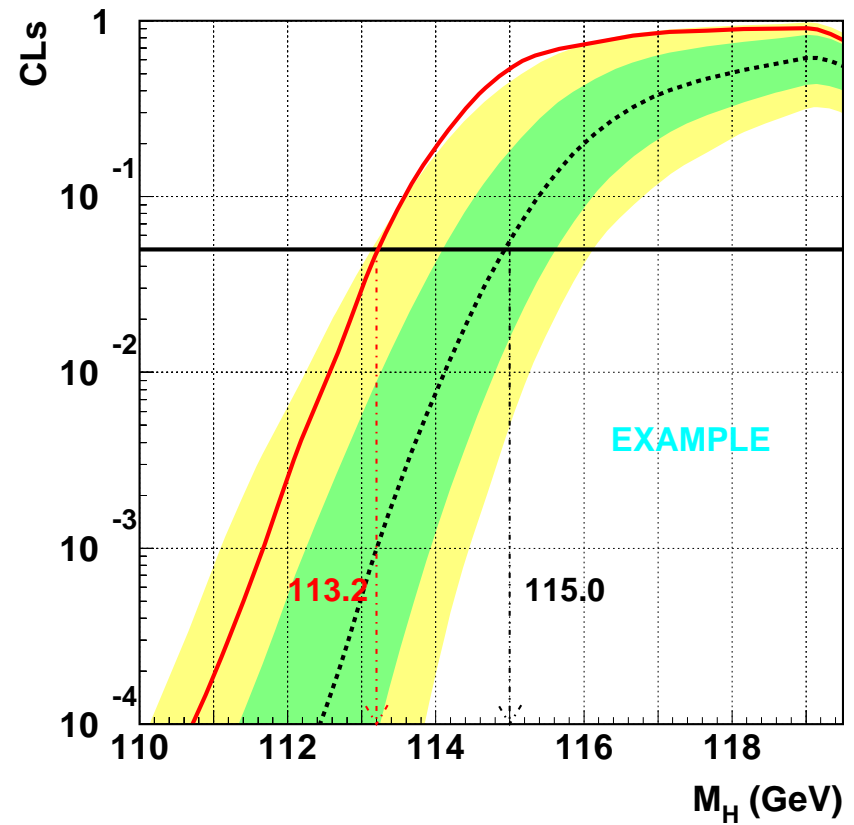


CL_b and CL_s

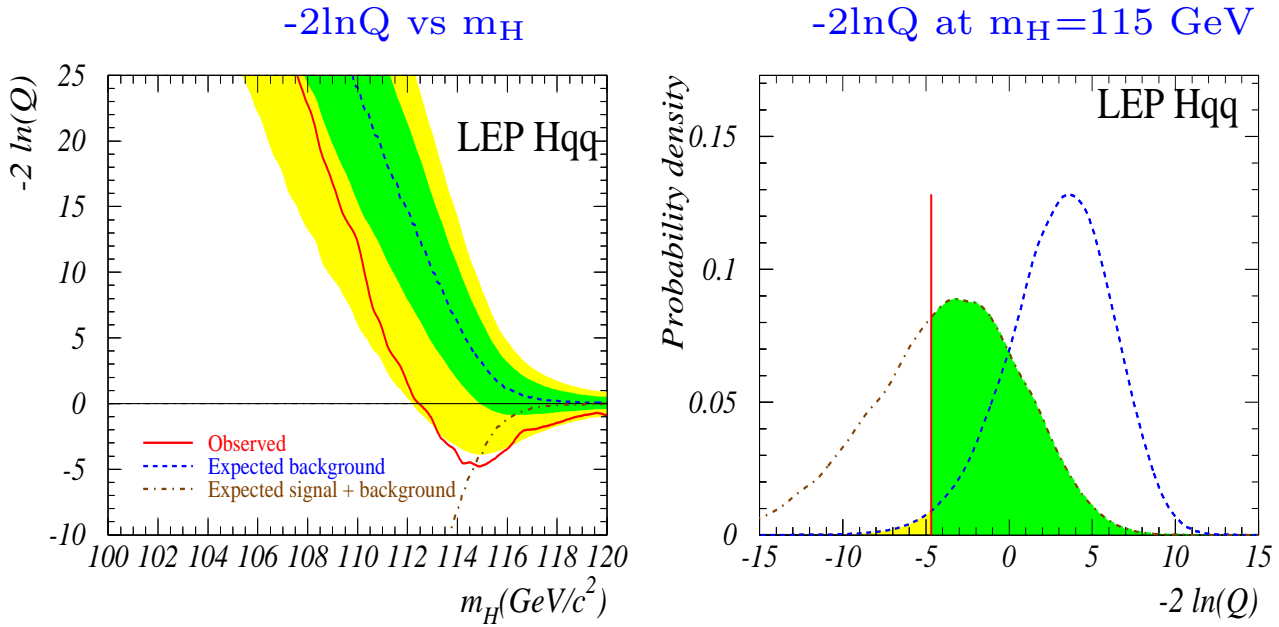
Discovery: no signal dependent



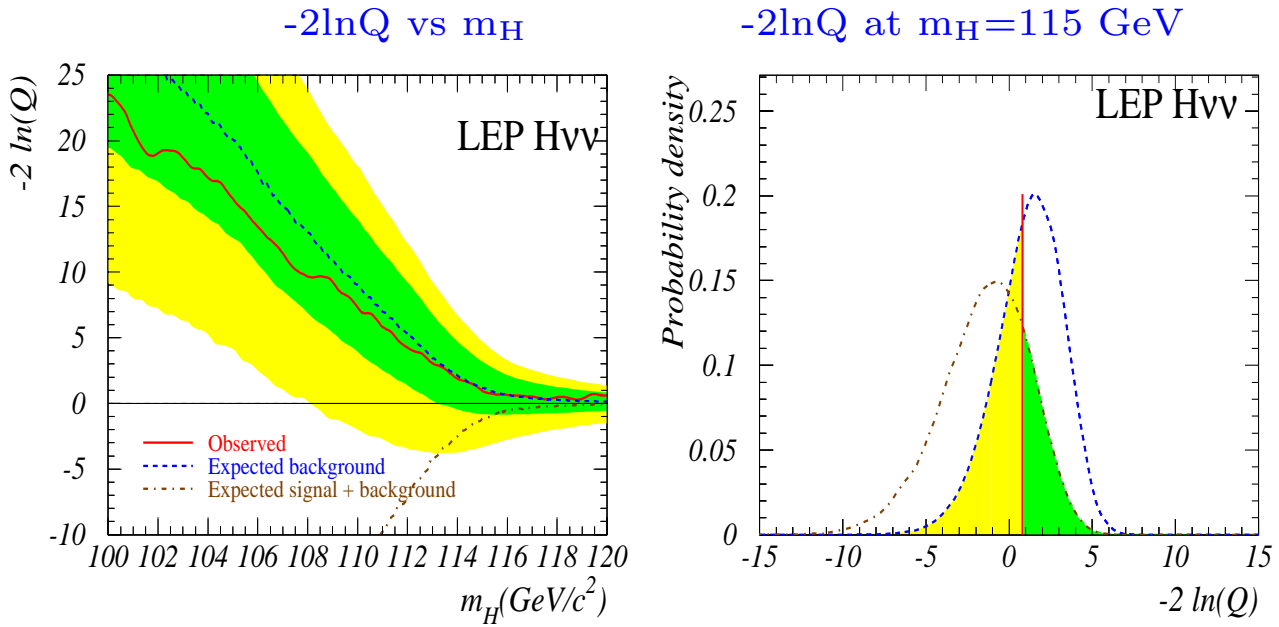
Exclusion: $CL_s = CL_{s+b}/CL_b$

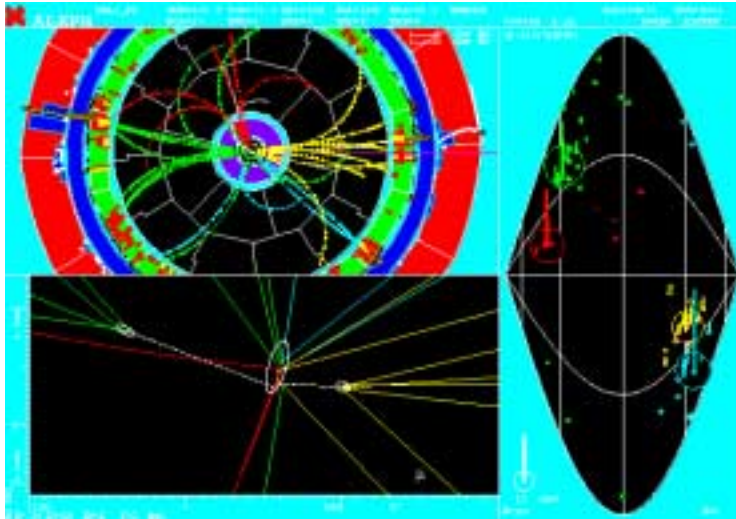


LEP-wide combination in Hqq channel



LEP-wide combination in H $\nu\nu$ channel





$$b_1 = 0.99$$

$$b_2 = 0.99$$

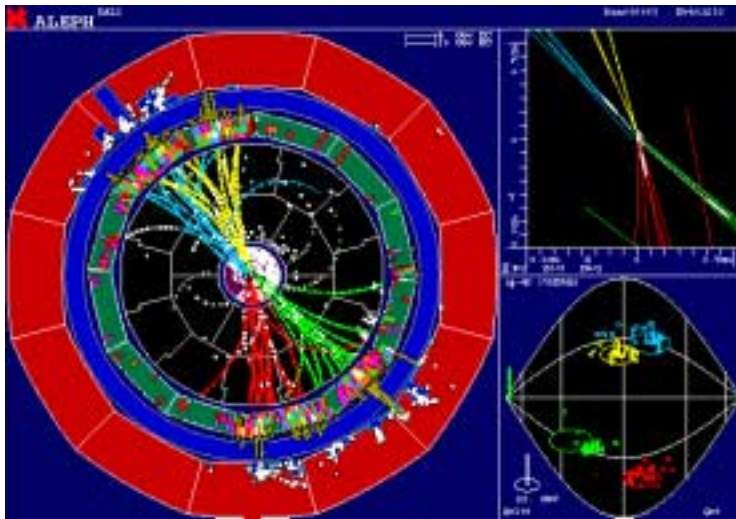
$$b_3 = 0.14$$

$$b_4 = 0.01$$

$$m_{12} = 104.2 \text{ GeV}$$

$$m_{34} = 101.3 \text{ GeV}$$

$$m_H^{\text{rec}} = 114.3 \text{ GeV}$$



$$b_1 = 0.99$$

$$b_2 = 0.99$$

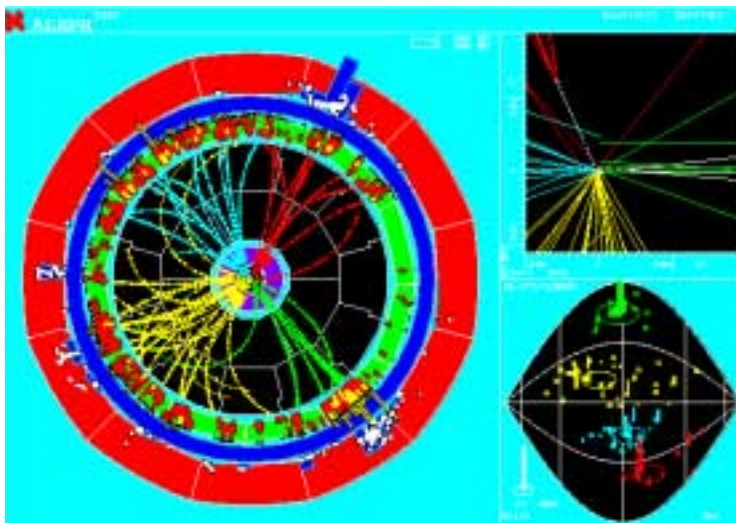
$$b_3 = 0.99$$

$$b_4 = 0.78$$

$$m_{12} = 109.2 \text{ GeV}$$

$$m_{34} = 94.9 \text{ GeV}$$

$$m_H^{\text{rec}} = 112.9 \text{ GeV}$$



$$b_1 = 0.99$$

$$b_2 = 0.21$$

$$b_3 = 0.99$$

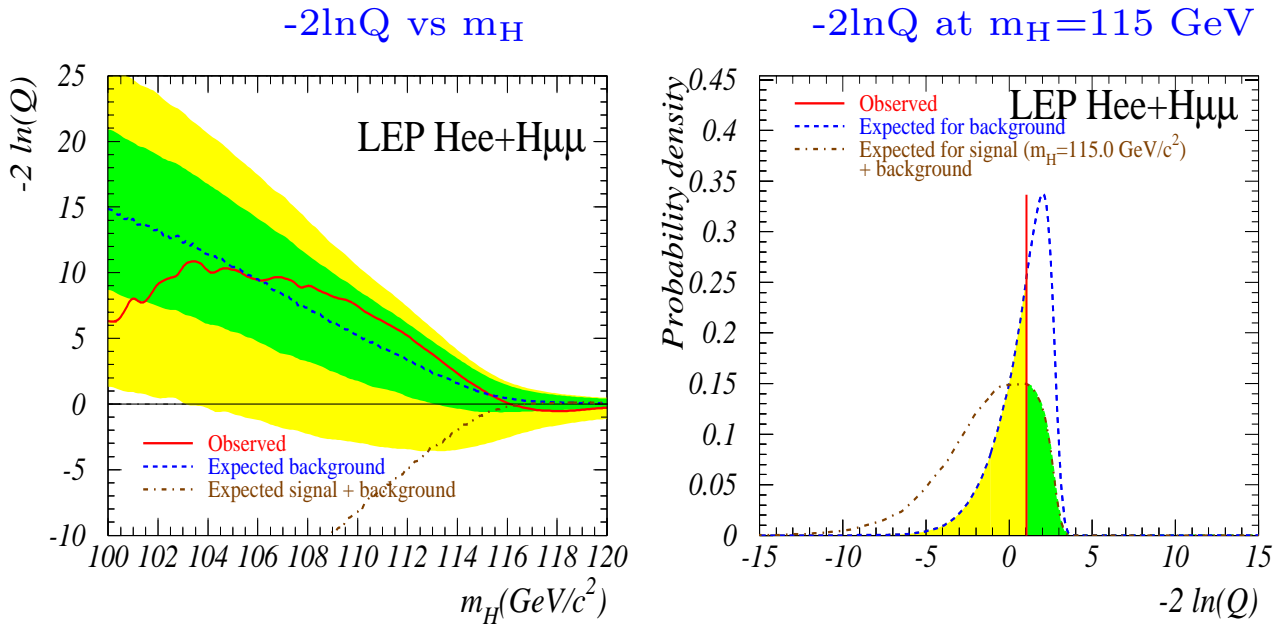
$$b_4 = 0.84$$

$$m_{12} = 104.9 \text{ GeV}$$

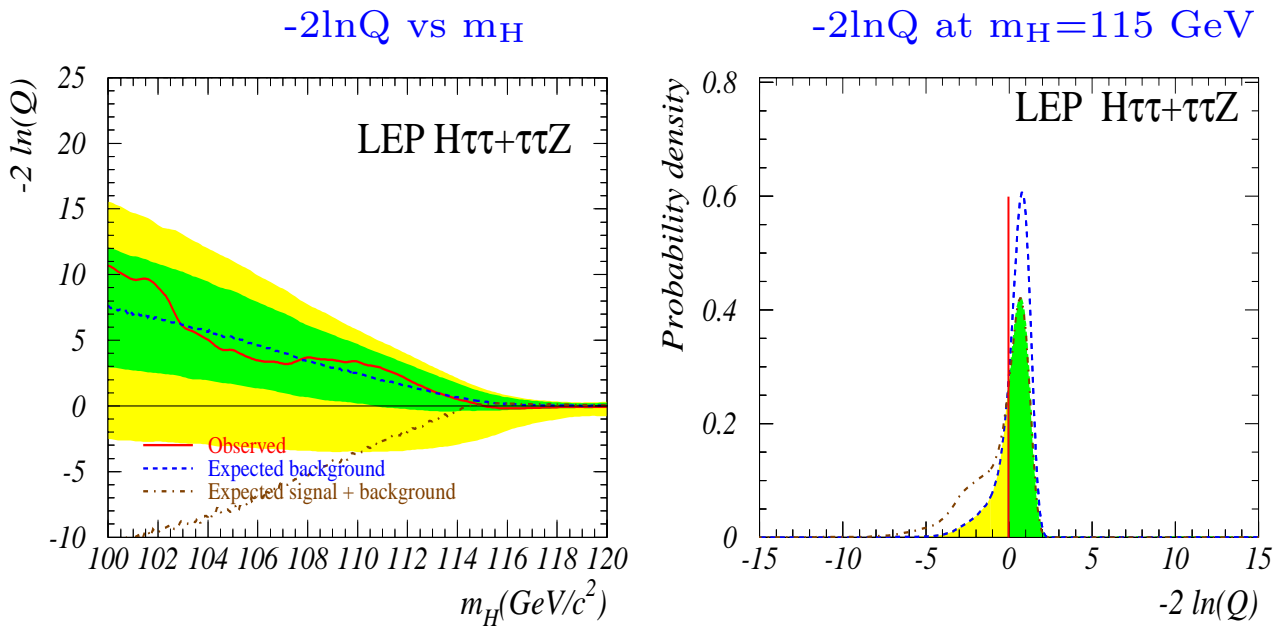
$$m_{34} = 96.3 \text{ GeV}$$

$$m_H^{\text{rec}} = 110.0 \text{ GeV}$$

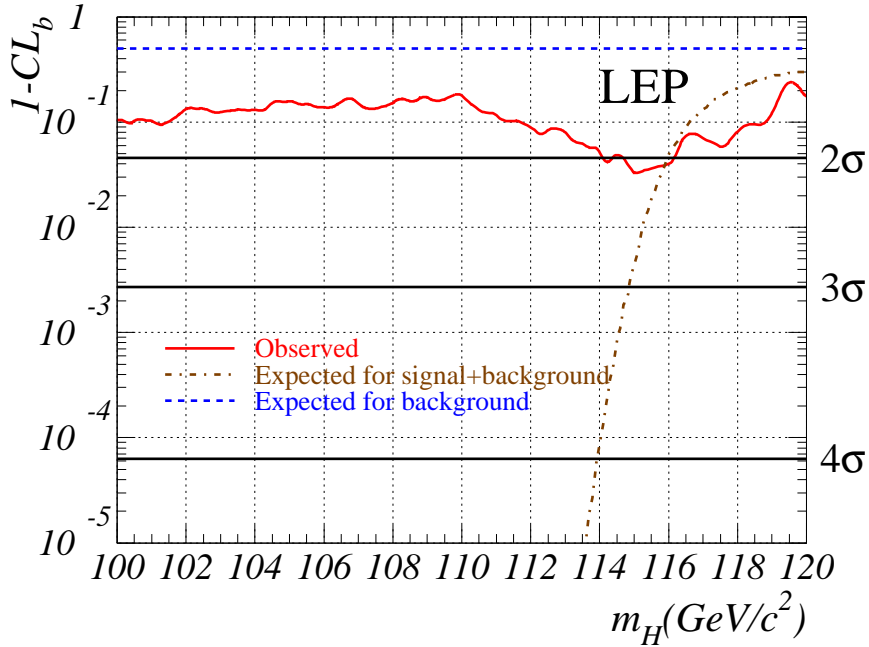
LEP-wide combination in $H\ell^+\ell^-$ channel



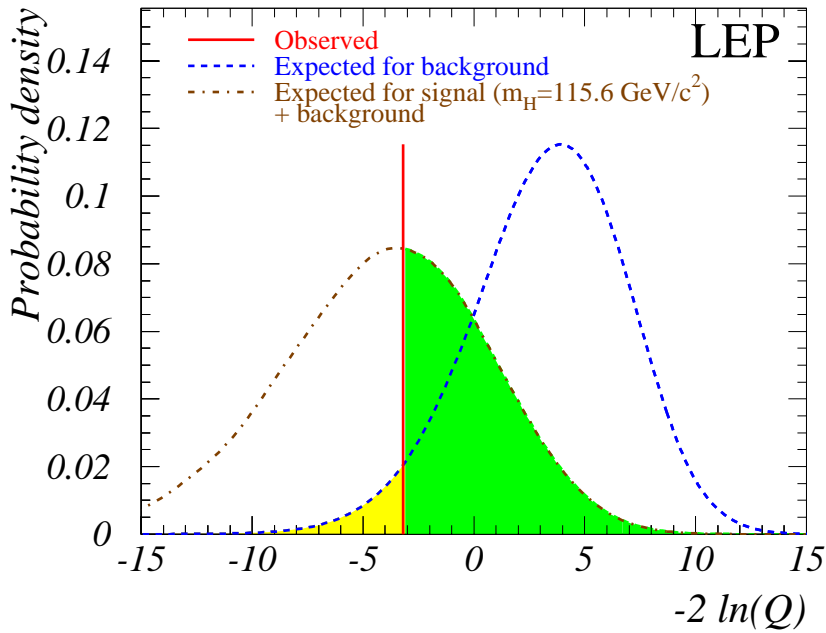
LEP-wide combination in tau-channels



$1 - CL_b$ vs m_H



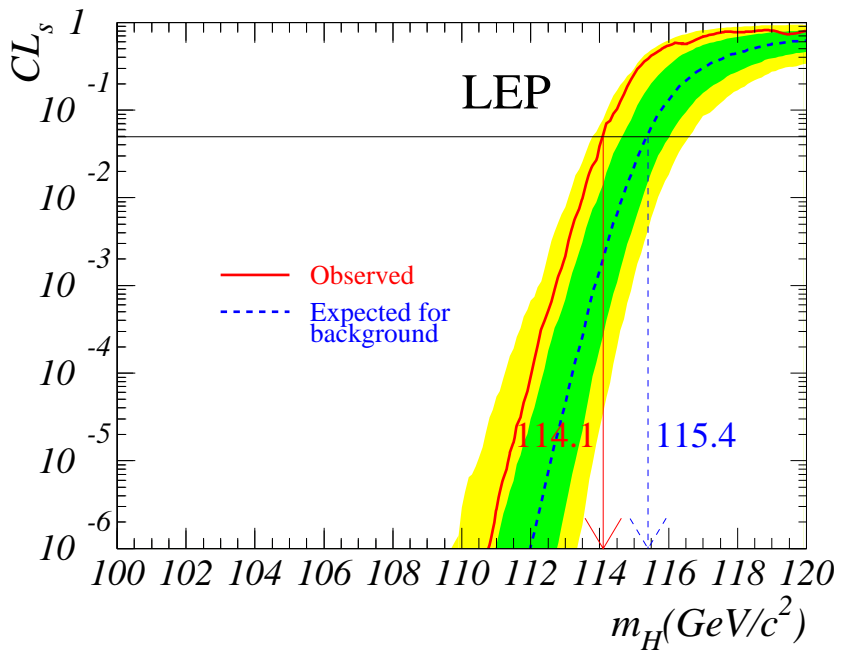
$-2\ln Q$ at $m_H=115.6$ GeV



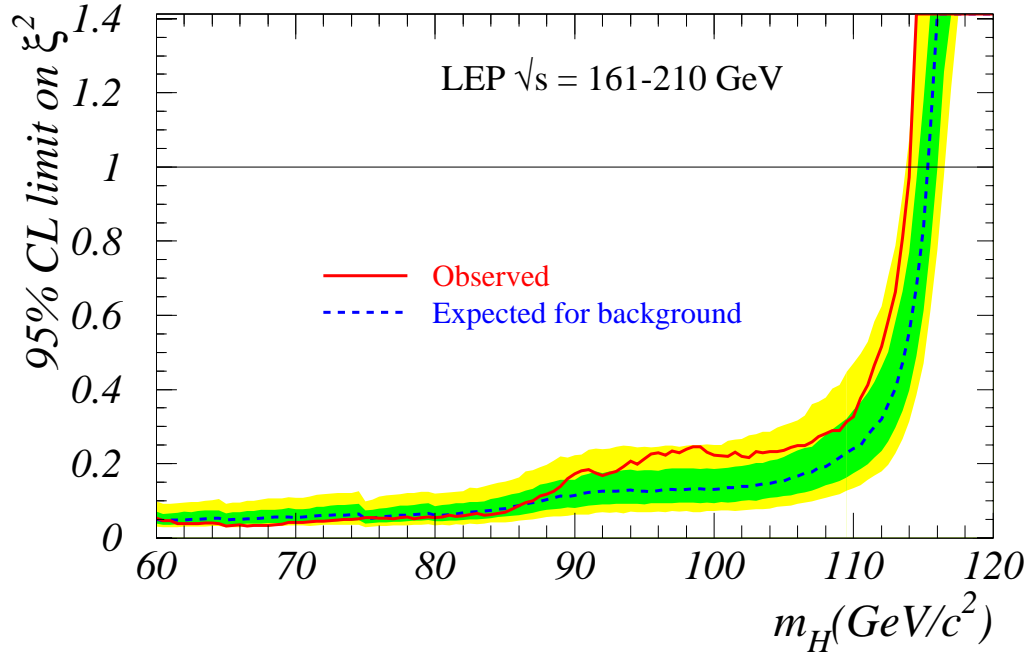
$1 - CL_b = 3.4\%$ at $m_H = 115.6$ GeV

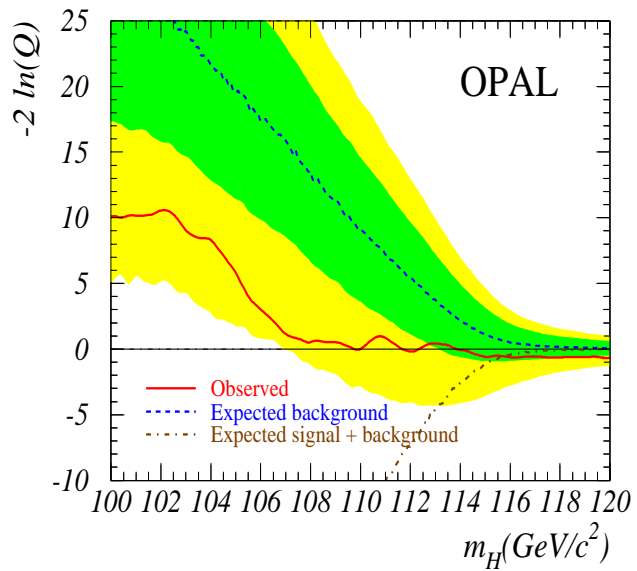
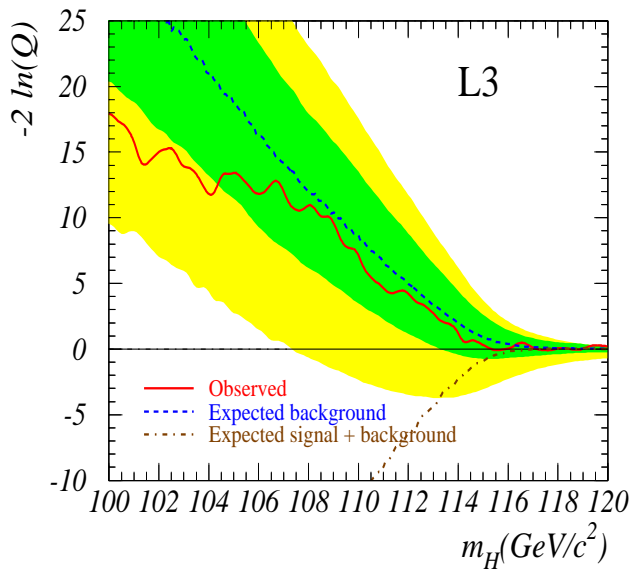
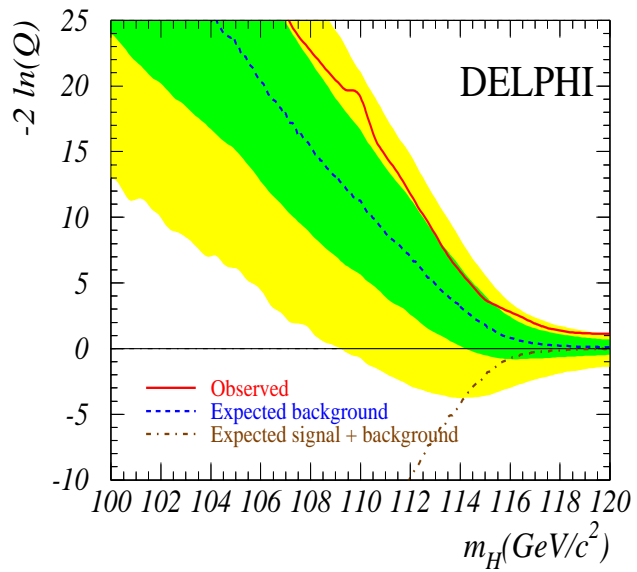
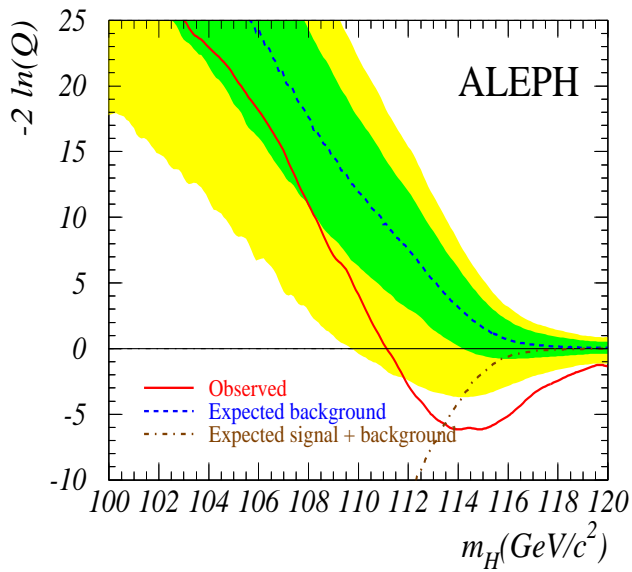
$CL_{s+b} = 44\%$ at $m_H = 115.6$ GeV

CL_s vs m_H



95 % CL limit on $\xi = g_{HZZ}/g_{HZZ}^{\text{SM}}$





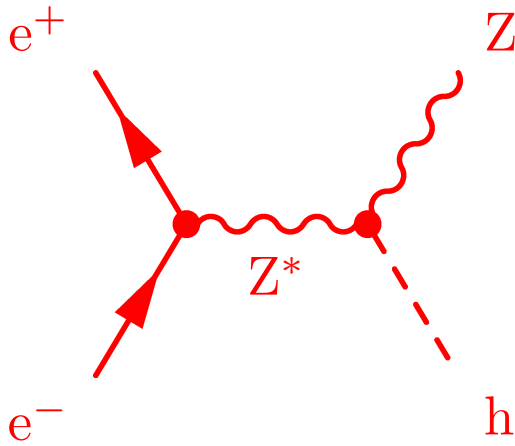
$$\Phi_1 = \begin{pmatrix} \Phi_1^+ \\ \Phi_1^0 \end{pmatrix} \quad \Phi_2 = \begin{pmatrix} \Phi_2^+ \\ \Phi_2^0 \end{pmatrix}$$

$$\langle \Phi_1 \rangle = \begin{pmatrix} 0 \\ \nu_1 \end{pmatrix} \quad \langle \Phi_2 \rangle = \begin{pmatrix} 0 \\ \nu_2 \end{pmatrix}$$

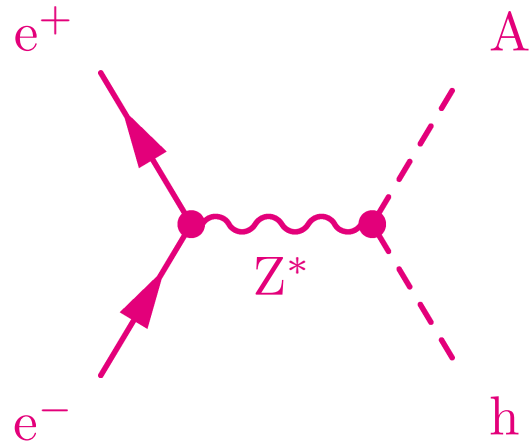
5 physical states: h, H, A, H^\pm

$$\tan \beta = \nu_1 / \nu_2$$

$$\begin{pmatrix} h \\ H \end{pmatrix} = \begin{pmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{pmatrix} \begin{pmatrix} \phi_1^0 \\ \phi_2^0 \end{pmatrix}$$



$$g_{hZZ} / g_{HZZ}^{\text{SM}} = \sin(\beta - \alpha)$$



$$g_{hAZ} / g_{HZZ}^{\text{SM}} = \cos(\beta - \alpha)$$

$$g_{hb\bar{b}} / g_{Hb\bar{b}}^{\text{SM}} = -\sin \alpha / \cos \beta$$

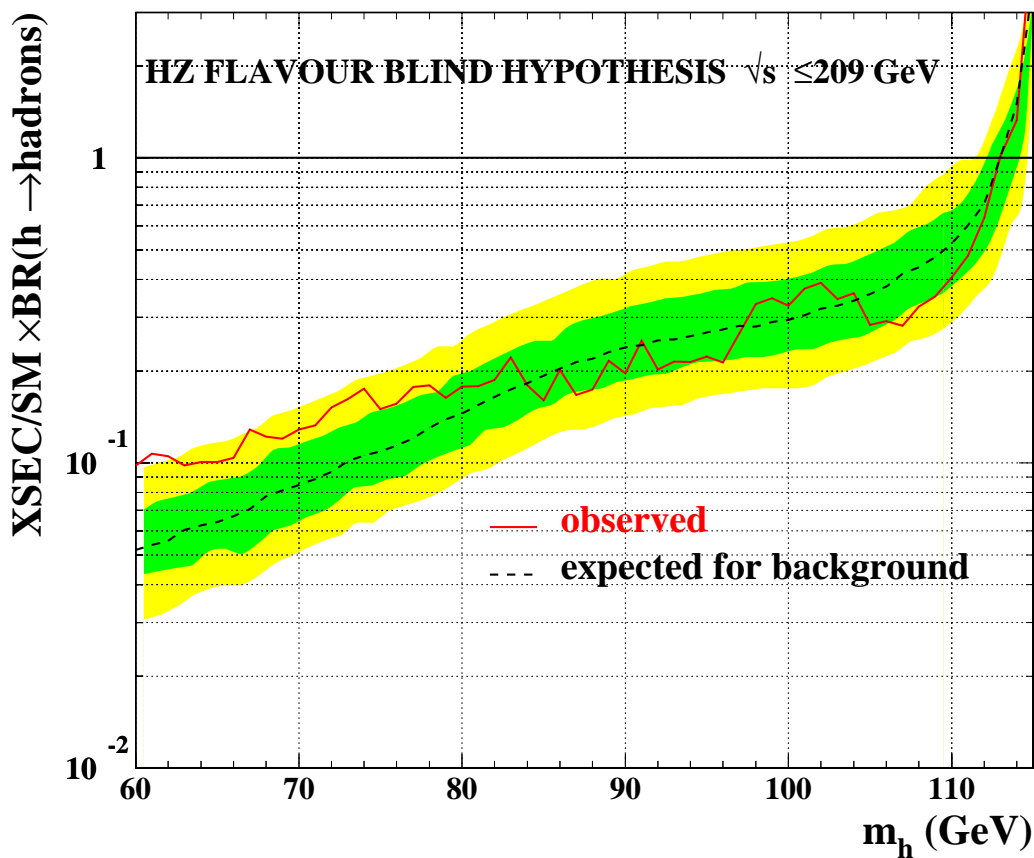
$$g_{Ab\bar{b}} / g_{Hb\bar{b}}^{\text{SM}} = \tan \beta$$

$$g_{hc\bar{c}} / g_{Hc\bar{c}}^{\text{SM}} = \cos \alpha / \sin \beta$$

$$g_{Ac\bar{c}} / g_{Hc\bar{c}}^{\text{SM}} = \cot \beta$$

- Analysis must be equally efficient for $h \rightarrow b\bar{b}, c\bar{c}, gg$ decays \implies
usage of any flavour-tag is disfavoured
- The weakest signal efficiency amongst $h \rightarrow b\bar{b}, c\bar{c}, gg$ is chosen \implies
model independence of search results

LEP PRELIMINARY



$m_h > 112.9$ GeV at 95 % CL assuming SM-like cross-section for hZ and $Br(H \rightarrow \text{hadrons}) = 100$ %
 $(m_h^{\text{exp}} > 113.0$ GeV)

At tree level

$$m_h^2 = \frac{1}{2} \left[m_A^2 + m_Z^2 - \sqrt{(m_A^2 + m_Z^2)^2 - 4m_Z^2 m_A^2 \cos^2 2\beta} \right]$$
$$m_h < m_Z$$

Radiative corrections $\implies m_h$ up to ~ 135 GeV

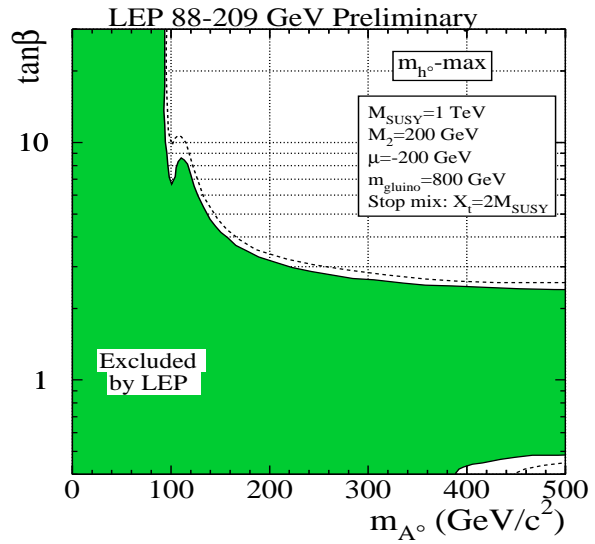
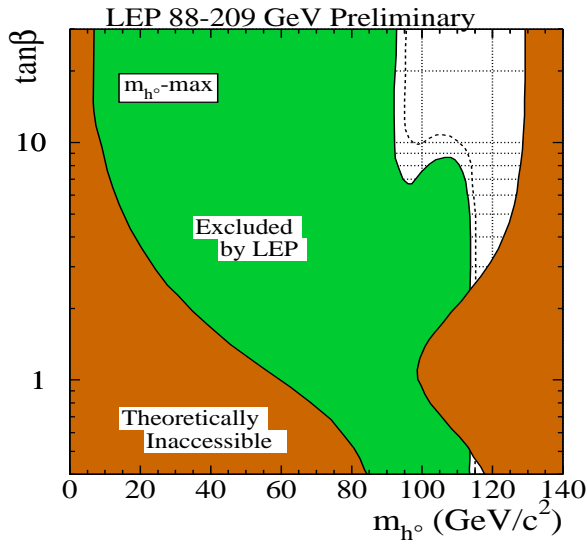
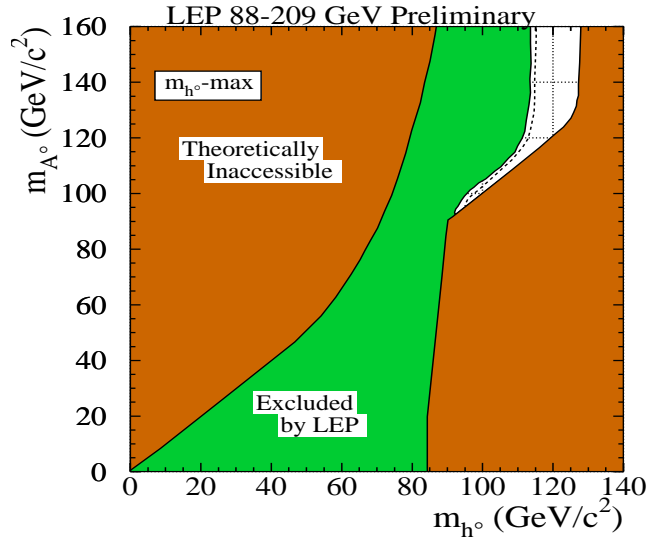
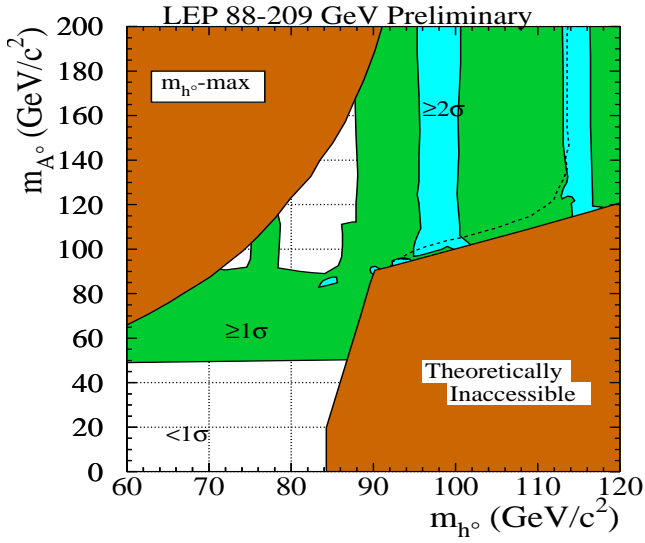
Interpretation of data within MSSM

- Basic theoretical assumptions \implies
reduction of the number of free MSSM parameters
without affecting Higgs phenomenology
- Particular sets of MSSM parameters \rightarrow
benchmark scenarios
- “mh-max” scenario
 $\rightarrow m_h(m_A, \tan \beta)$ acquires maximal value

M_{SUSY}	$= 1 \text{ TeV}$
M_2	$= 200 \text{ GeV}$
μ	$= -200 \text{ GeV}$
$X_t (\equiv A - \mu \cot \beta)$	$= 2M_{\text{SUSY}}$
$m_{\tilde{g}}$	$= 800 \text{ GeV}$

- scan over m_A and $\tan \beta$ is performed
- $m_h(m_A, \tan \beta)$ is calculated using 2-loop
diagrammatic approach (FeynHiggs program,
S.Heinemeyer, W.Hollik, G.Weiglein, Eur. Phys.
Jour. C9 (1999))

hZ search topologies are complemented with
 $hA \rightarrow b\bar{b}b\bar{b}, b\bar{b}\tau^+\tau^-$ channels



95 % CL mass limits

	observed	expected
$m_h >$	91.0 GeV	94.6 GeV
$m_A >$	91.9 GeV	95.0 GeV

$0.5 < \tan \beta < 2.4$ is excluded at 95 % CL

- LEP-wide searches for Neutral Higgs boson(s) within the framework of SM, general 2HDM and MSSM are performed.
- SM Higgs Searches
excess of $\sim 2.1 \sigma$ is observed at $m_H \sim 115$ GeV;
95 % CL mass limit $m_H > 114.1$ GeV is set.
- Searches for Higgs bosons of MSSM
mass limits $m_h > 91.0$ GeV and $m_A > 91.9$ GeV are derived at 95 % CL;
 $0.5 < \tan \beta < 2.4$ is excluded at 95 % CL (“mh-max” scenario).
- Flavour-independent search for hZ process
result is interpreted in terms of 95 % CL limits on hZ cross-section;
assuming SM like cross-section for hZ and $\text{Br}(H \rightarrow \text{hadrons}) = 100$ %
mass limit $m_h > 112.9$ GeV is set regardless of flavour content of h decay products.
- TEVATRON and LHC are next in the line to discover Higgs boson(s)
- Linear e^+e^- colliders - Higgs factory \Rightarrow precise determination of Higgs boson profile