

Some Considerations on

CSR Microbunching Instability

( SASE in a Bunch Compressor ? )

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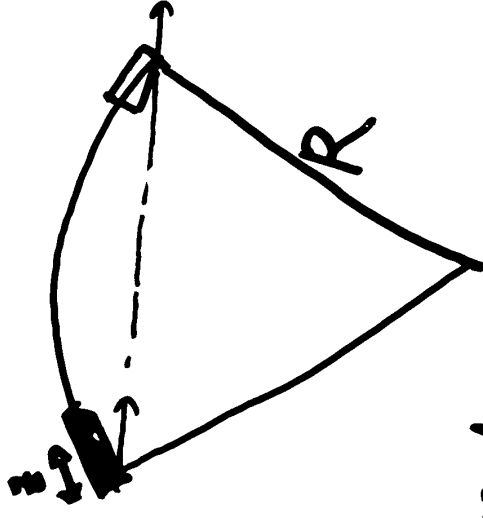
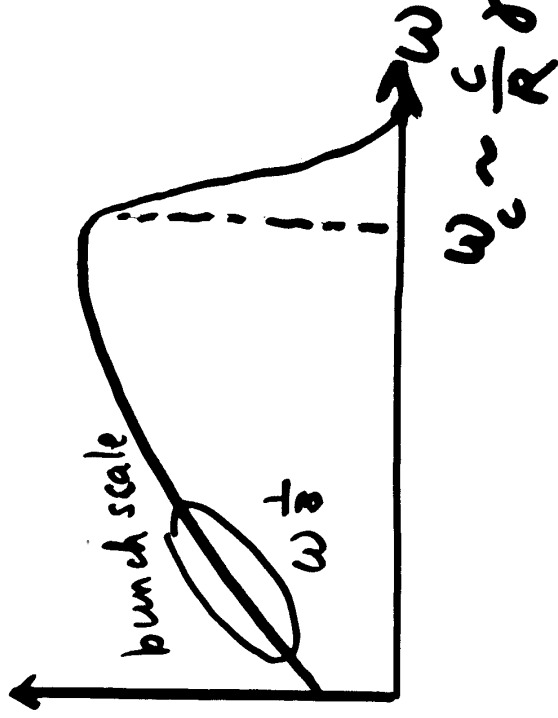
## Introduction

- Bunch Compressor (BC) is design to be a very effective buncher ( $\sigma_s \rightarrow \frac{\sigma_s}{10}$ )  
 $\Rightarrow$  also very effective to grow microbunching!  
due to synchrotron radiation "wake"
- Radiative interaction in a BC is similar to FEL interaction

{ start up from noise (SASE)  
high gain  
saturation

# Synchrotron Radiation Impedance/Wake

$$\frac{dP}{d\omega}$$



Radiative interaction  
within the bunch

⇒ impedance  $Z_{||} \sim \omega^{1/3}$

$$\text{CSR "wake" } W \sim \left( -\frac{1}{Z} \frac{\partial}{\partial z} \right)$$

$z$ : bunch coordinate

(Murphy, Derbenev at el. 1995)

- Steady state, no shielding. 1D line charge

# Micro bunching

- BC (chicane): energy modulation  $\rightarrow$  density modulation  
"optical klystron", "HG HG"

$$\frac{dZ}{dS} = \frac{R'_{56}(S)}{R_{56}} \delta \quad (R'_{56} \equiv \frac{dR_{56}}{dS})$$

- SR impedance/wake: density mod.  $\rightarrow$  energy mod.

$$\frac{d\delta}{dS} \propto \int_{-\infty}^z \frac{1}{(z-z')^{1/3}} dz' \frac{dN_b}{dz}$$

loop for instability (Borland/Emma, Schneidmiller, Heifetz/Stupakov)

- Stability criterion ( $\sim$  Boussard's microwave stability)

$$\left| \frac{Z''}{N} \right| \sim \left( \frac{\omega R}{c} \right)^{-2/3} < \frac{IA \delta^2 \sigma^2}{I \left( \frac{1}{R_{56}} \right)} \left| R'_{56} \right|$$

CSR  $Z'' \sim N^{1/3}$

What's the source of modulation?

- a non-smooth bunch shape has high frequency content

- Shot noise  $\rightarrow$  density fluctuation  
(incoherent synchrotron radiation)

at modulation wavelength  $\lambda$

noise bunching  $b \sim \frac{1}{\sqrt{N_\lambda}} \leftarrow$  # of particles in  $\lambda$

$$\sim 10^{-3} - 10^{-4}$$

- Similar to SASE, but broad-band v.s. narrow-band

SASE  $b_{\text{eff}} \sim \frac{1}{\sqrt{N_{e_c}}} \leftarrow$  # of particles in coherence length  
 $l_c \sim 100\lambda$

# High Gain Regime

• like FEL, introduce scaling "p" (stability  $\sigma_s > p$ )

$$p \sim \sqrt{\frac{I}{I_A} \frac{1}{\gamma} \frac{1}{(-R_{56})}} \left( \frac{\lambda}{R} \right)^{\frac{1}{3}} \sim 10^{-4} \quad (\lambda \sim 10 \text{ to } 1 \mu\text{m})$$

• PC Gun simulation: incoherent  $\sigma_s \sim 10^{-5}$  @ 250 MeV

$\Rightarrow \sigma_s \ll p$  (very unstable against microbunching)

• High-gain approximation

instability occurs much faster than bunch compression

$\Rightarrow$  use the instantaneous I,  $\lambda$  for the instability calculation

# Gain length

- When  $\sigma_8 \ll \rho$ , gain length  $\frac{\lambda}{\rho(-R_{56})} \sim \frac{\lambda^{2/3}}{(-R_{56})^{1/2}}$
- Compare to FEL gain length  $\frac{\lambda_u}{\rho} \sim \frac{(2\sigma^2)\lambda}{\rho}$

BC (Chicane) undulator

$$(-R_{56}) \sim \frac{\sigma^2}{R} \gg \frac{1}{\sigma^2}$$

Bunching: path length difference

velocity difference



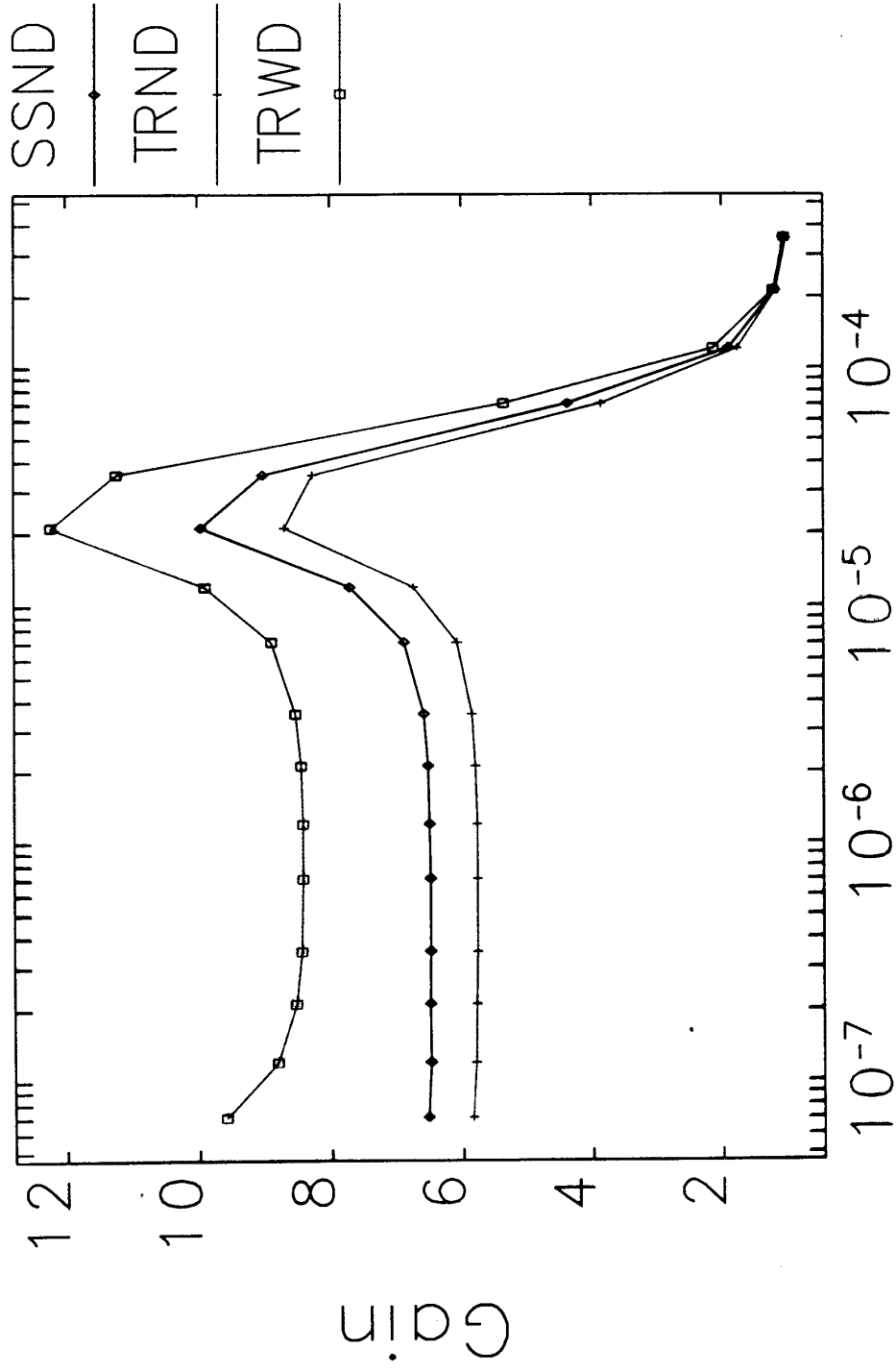
much more effective

FEL instability can

when  $\sigma_8 < \rho$  (micros)

occur even at 1 Å

# Elegant Simulations (M. Borland), no compression



Input Modulation Amplitude (rms)

Group bc1-07Nov01

LCLS BC1,  $\lambda = 10 \mu\text{m}$ ,  $G_{\text{theory}} = 6$

# Numerical Calculations

	LEUTL	LCLS BC1	BC2	Test-chicane
energy [GeV]	0.15	0.25	4.5	5
charge [nC]	0.5	1	1	1
$\sigma_s$ [ $\mu\text{m}$ ]	540	830	190	200
incoherent $\sigma$	6e-5	1.2e-5	3e-5	2e-6
chirp [ $\text{m}^{-1}$ ]	12	21.4	40	36
bending R [m]	0.8	2.5	12.2	10.35
length per bend [m]	0.192	0.203	0.4	0.5
$R_{56}$ [mm]	65	36	22	25
$(1 + hR_{56})^{-1}$	4.4	4.4	8.4	10
$-R'_{56}$	0.17	0.089	0.0275	0.025
wavelength $\lambda$ [ $\mu\text{m}$ ]	10/1	10/1	3	1
scaling $\rho$ (at $s=0$ )	1.9e-4/8.6e-5	1.5e-4/7.2e-5	5.3e-5	3.8e-5
relative Gain	500/1	150/1e11	$\sim 1$	2e6
saturation $\sigma$		3.8e-4/1.9e-4		1.3e-4
$\Delta\epsilon_n$ [ $\mu\text{m}$ ]		$\sim 0.1$		$\sim 1$

initial }

Table 1: CSR microbunching instability in bunch compressors

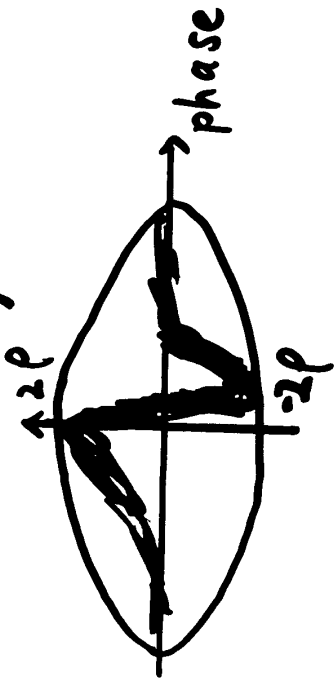
# Saturation

- Microbunching instability saturates when  $|b| \sim 1$

$$\text{noise } |b| \sim \frac{1}{10^3 \sim 10^4}$$

- All it takes a gain of  $10^3 \sim 10^4$

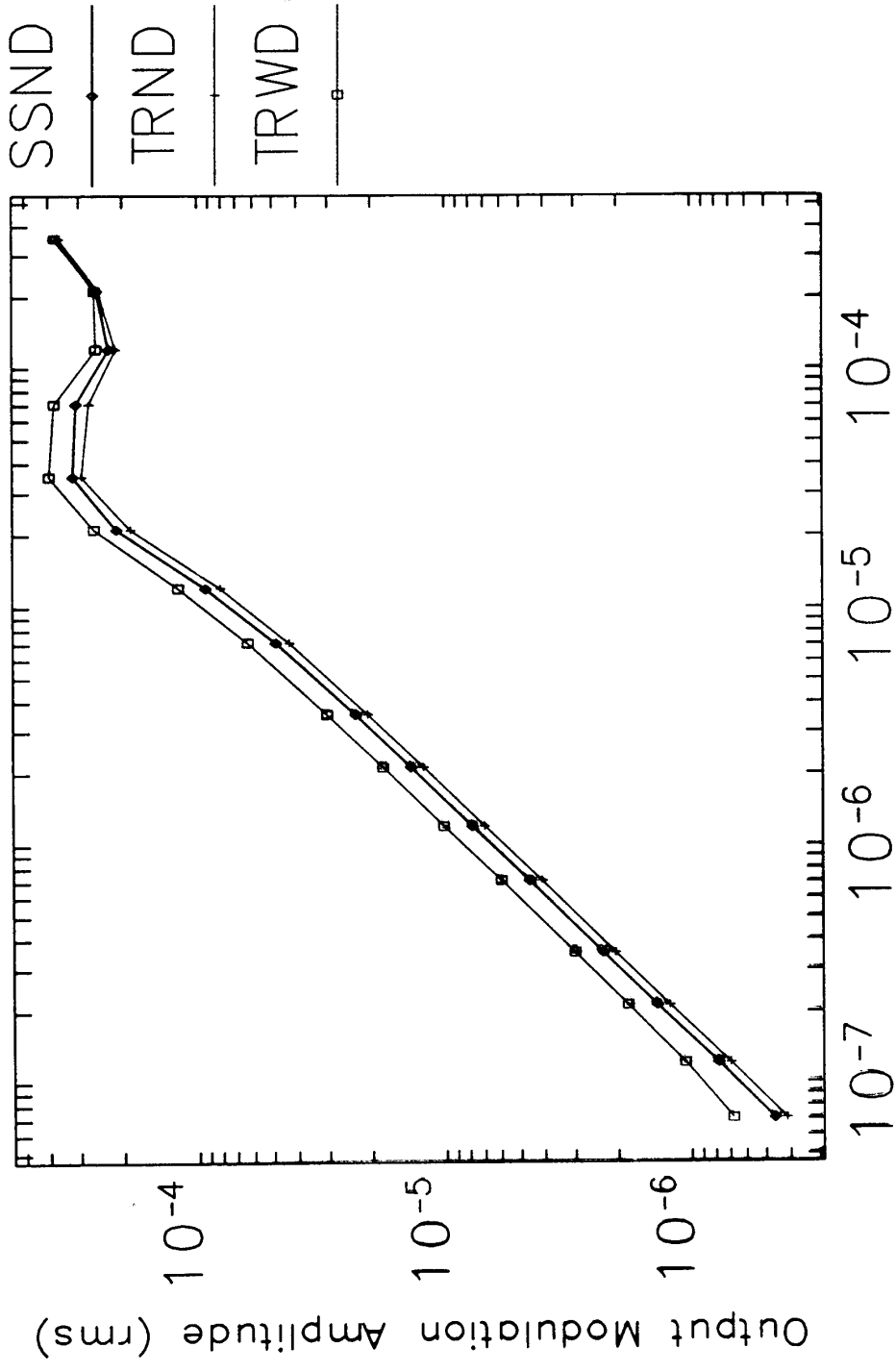
- At saturation, electrons trapped in bucket



$$\sigma_s \sim 2\rho \sim 10^{-4}$$

←  $\lambda$  → sliced energy spread

# Elegant Saturation Simulation (M. Borland)



Input Modulation Amplitude (rms)  $10 \mu m$

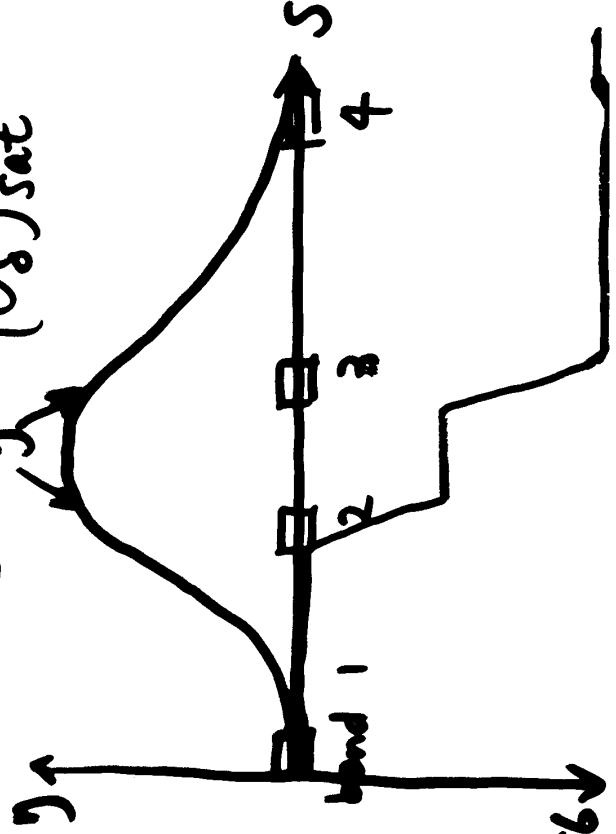
Group bc1-07Nov01 new BCI

# Emittance Growth

- Energy loss in the bend  $\xrightarrow{\text{dispersion}}$  emittance growth
- RMS size change (Emma/Brinkmann, PAC97)

$$\langle \Delta X \rangle_{\text{rms}} = \int R_{16} \frac{d\sigma_s}{ds} ds$$

$$= \int J \cdot (\sigma_s)_{\text{sat}} \sim \eta \cdot (2\rho)$$



$$\frac{\epsilon_n}{\epsilon_0} = \sqrt{1 + \frac{\gamma(1+\alpha)}{\epsilon_0 \beta}} \sqrt{\langle \Delta X \rangle^2}$$

$$-\int \frac{d}{ds} R_{56} ds = R_{56}$$

# Possible Cures

- Unlike FELs, don't want CSR microbunching
- Stability requires

$$\underbrace{\frac{\delta_x}{\delta_{BC}} \rho_{FEL}}_{10^{-3}} > \underbrace{\frac{\sigma_s}{(1+hR_{56})}}_{\text{increase}} > \underbrace{\rho_{BC} \sim \sqrt{\frac{I}{I_A} \frac{1}{\delta_{BC}} \frac{1}{(-R'_{56})}} \left(\frac{\lambda}{R}\right)^{\frac{1}{3}}}_{\text{decrease}}$$

wiggler before BC2 faster compression helps

what about BC1? less bends to grow instab.

(a laser wiggler)? but how effective