

Field Calculations for Bunch Compressors (non-1D and 1D)

1 Calculation Methods

- 1.1 Direct Time Domain Calculation
- 1.2 Retarded Sources
- 1.3 Tracking
- 1.4 1D-Approach

2 Effects

- 2.1 Compression Work
- 2.2 Surface Impedance
- 2.3 Longitudinal Field
- 2.4 Transverse Effects
- 2.5 Transients
- 2.6 Shielding

1 Calculation Methods

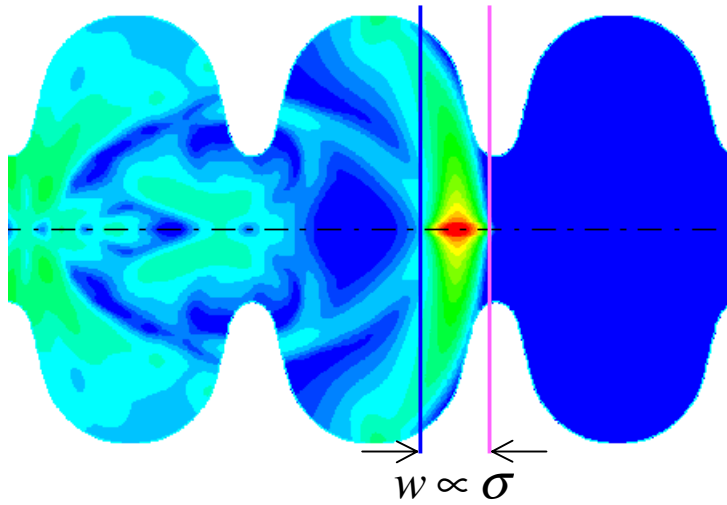
1.1 Direct Time Domain Calculation

e.g. FDTD - method with particle tracking
PIC codes (Particle In Cell)

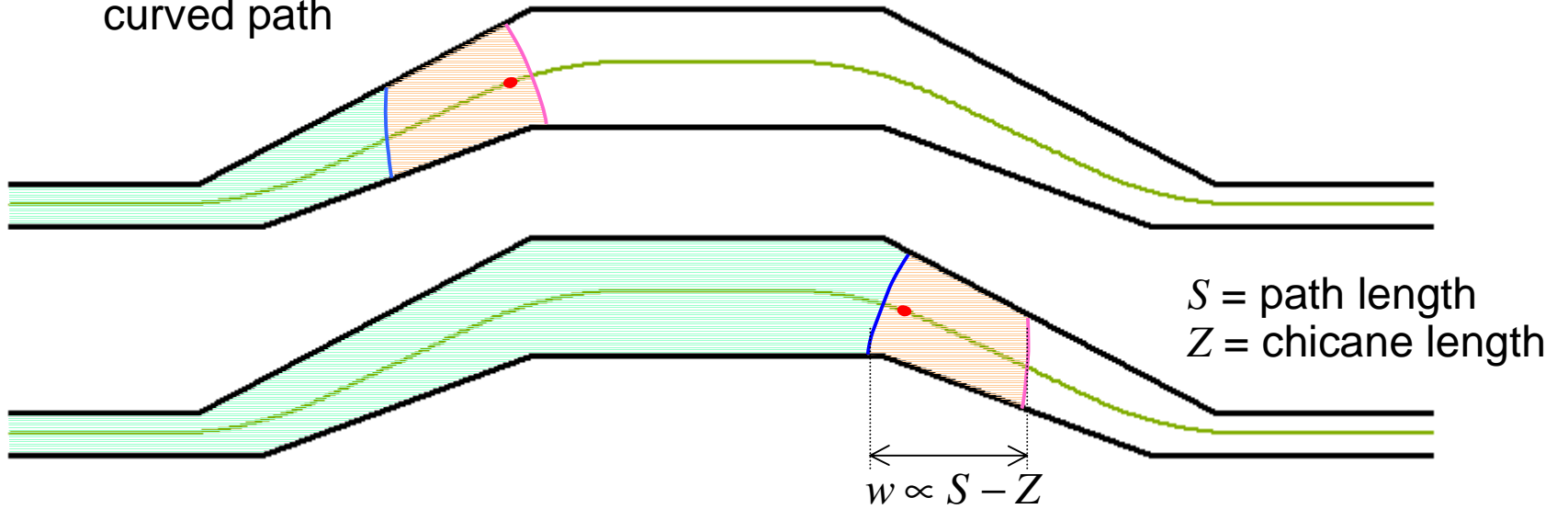
- a) calculation window
- b) dispersion

a) calculation window

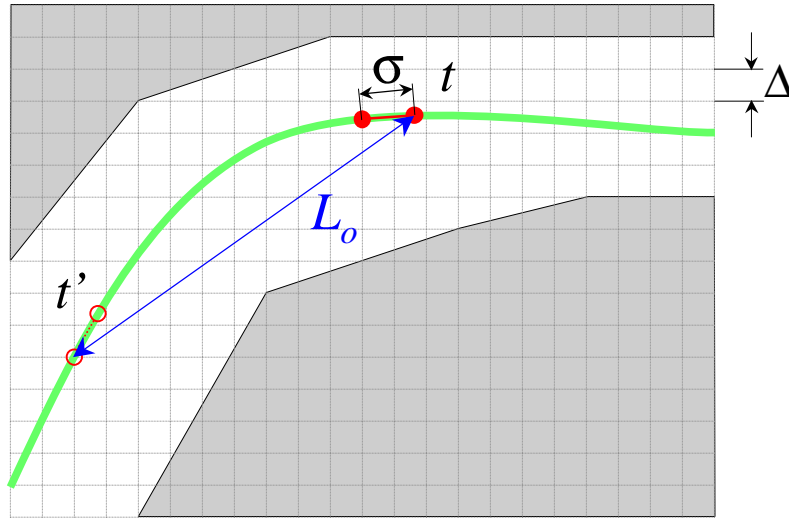
linear motion
($v = c_0$)



curved path



b) dispersion



numerical integration:

$$\|\vec{k}\| = \frac{\omega}{c_0} + O(\Delta^N)$$

Δ = step width

N = order of integration (e.g. FDTD $N=2$)

$$\text{error} \propto \frac{L_o \Delta^N}{\sigma^{N+1}} \quad \text{e.g. FDTD: } \frac{L_o \Delta^2}{\sigma^3} \lesssim 1$$

typical interaction length: $L_o = \sqrt[3]{24 R_0^2 \sigma}$

e.g. $\sigma = 100 \mu\text{m}$, $R_0 = 1 \text{ m}$, $\Rightarrow \Delta \propto 3 \mu\text{m}$

algorithms with low numerical dispersion needed:

α) compensation for one direction of propagation

β) higher order algorithms $\text{error} \propto \frac{L_o \Delta^N}{\sigma^{N+1}}$

1.2 Retarded Sources

a) point charge (Liénerd-Wiechert fields)

exact solution known;

problems: very fast time structure of 1-particle field $\tau \propto R_0/\gamma^3 c_0$
singularity at source point

} very many particles needed;

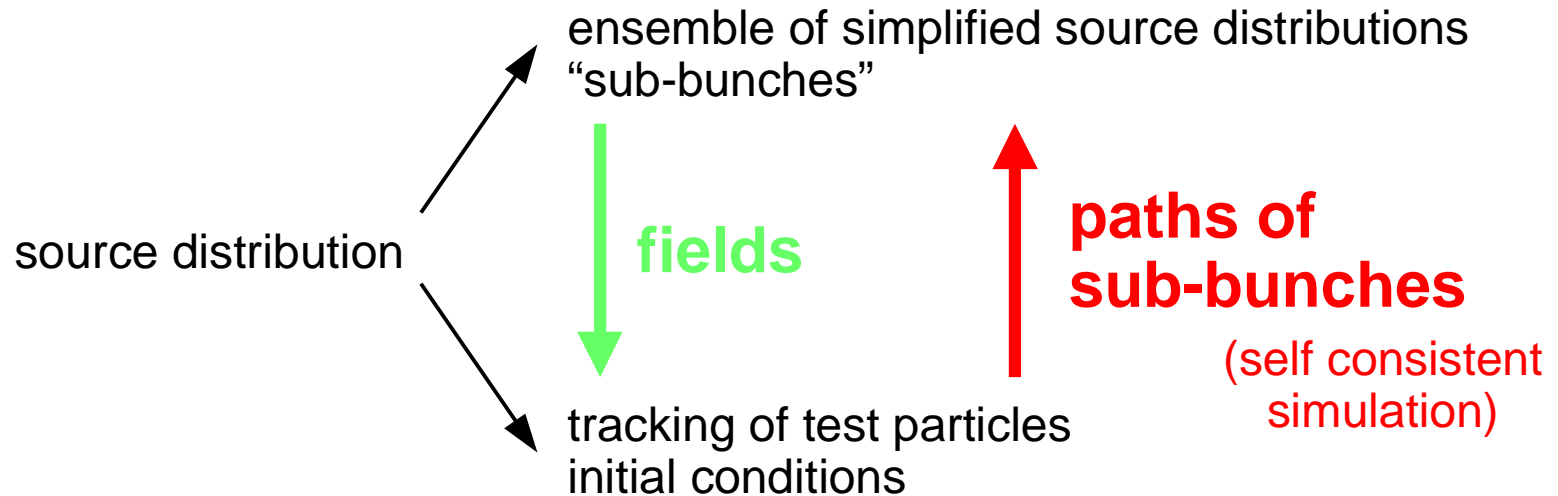
b) distributed sources

eg. $\vec{J}(\vec{r}, t)$ and $\rho(\vec{r}, t)$ known

$$\left. \begin{aligned} \Phi(\vec{r}, t) &= \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}', t')}{\|\vec{r} - \vec{r}'\|} dV' \\ \vec{A}(\vec{r}, t) &= \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}', t')}{\|\vec{r} - \vec{r}'\|} dV' \end{aligned} \right\} \begin{aligned} \vec{E}(\vec{r}, t) &= -\nabla\Phi - \partial_t \vec{A} \\ \vec{B}(\vec{r}, t) &= \nabla \times \vec{A} \end{aligned}$$

problems: \vec{J} and ρ have to be known for all points (x, y, z, t) ;
3d integration for every observation point

c) simplified distributed sources



computer codes: N.N. (R.Li, Jefferson Lab)
TraFiC4 (M.Dohlus, T.Limberg DESY, A.Kabel SLAC)

1.3 Tracking

a) partial cancellation of transverse forces

Transverse Effects of Microbunch Radiative Interaction
Y.S.Derbenev, V.D.Shiltsev

$$x'' + (K^2 - n)x = \frac{K\Delta\mathcal{E} + F_x}{\mathcal{E}}$$

$$K\Delta\mathcal{E} = K\Delta\mathcal{E}_0 + Kq_0 \int \text{CSR} dt + Kq_0(\Phi_0 - \Phi)$$

$$F_x = \dots + q_0 K \vec{v} \cdot \vec{A}$$

$K(s) = 1/R_0(s)$ = curvature

$\Delta\mathcal{E}$ = energy offset

$F_x(s,x,t)$ = transverse force

Φ = scalar potential

\vec{A} = vector potential

singular for 1D beams ($\sim \ln(r)$)
singularities cancel !

approach 1 (TRAFIC4): calculate forces, solve EOM

$$x'' + (K^2 - n)x = \frac{K\Delta\mathcal{E} + F_x}{\mathcal{E}}$$

avoid singularities:

real transverse beam dimensions have to be taken into account

“cancellation”:

F_x calculated by field solver

$\Delta\mathcal{E}$ calculated by EOM solver

approach 2 (R.Li): modified LHS, EOM

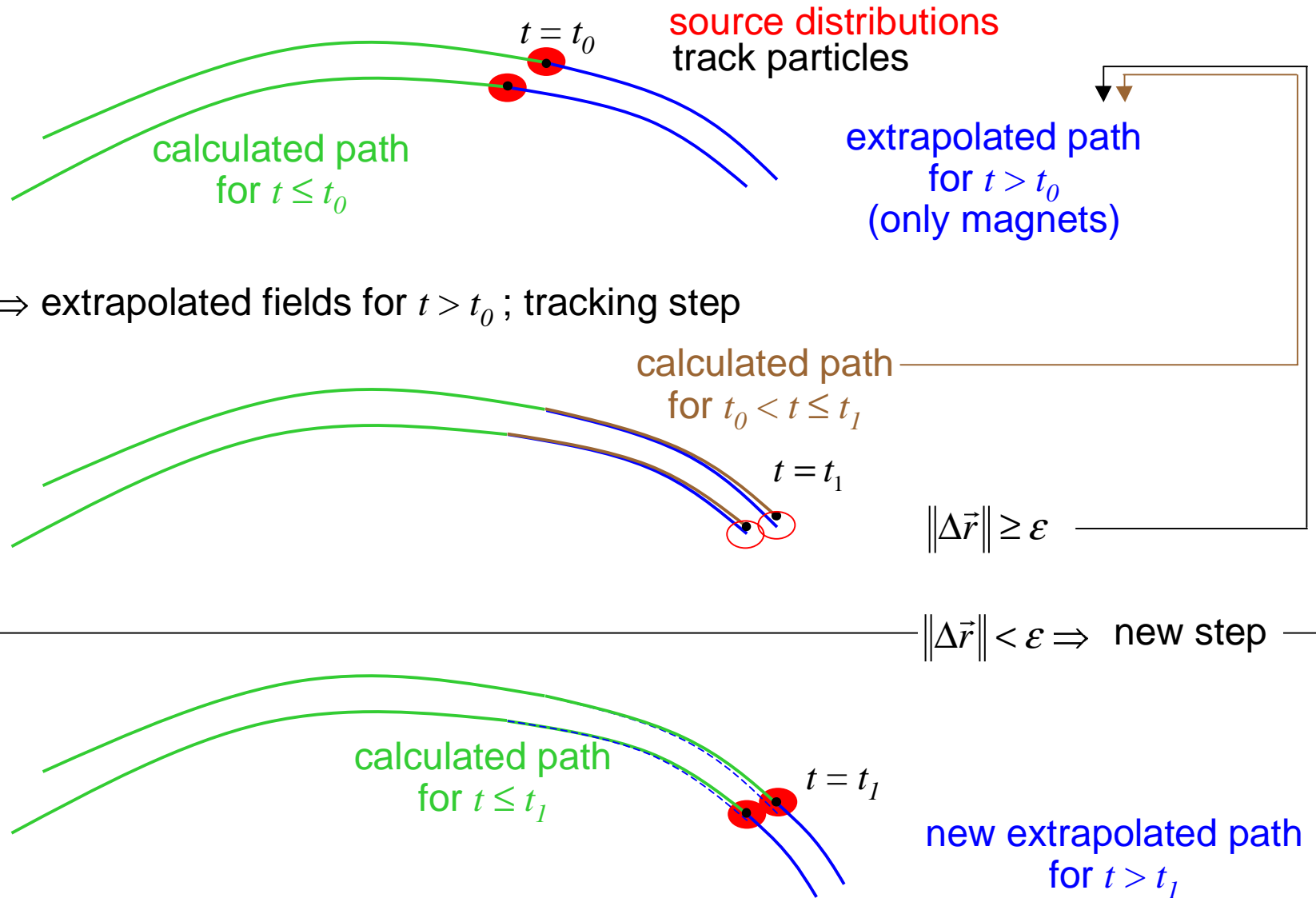
$$x'' + (K^2 - n)x = \frac{K\Delta\hat{\mathcal{E}} + \hat{F}_x}{\mathcal{E}}$$

$$K\Delta\hat{\mathcal{E}} = K\Delta\mathcal{E}_0 + Kq_0\Phi_0 + Kq_0 \int \text{CSR} dt$$

$$\hat{F}_x = \dots + q_0 K \vec{v} \cdot \vec{A} - Kq_0\Phi$$

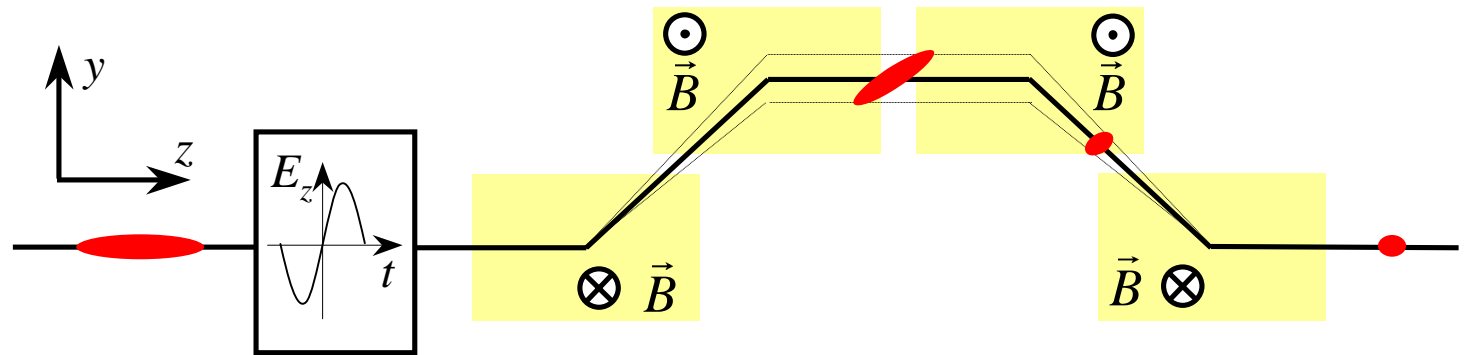
“cancellation” calculated by field solver

b) self consistent tracking



example: bunch compressor

principle:

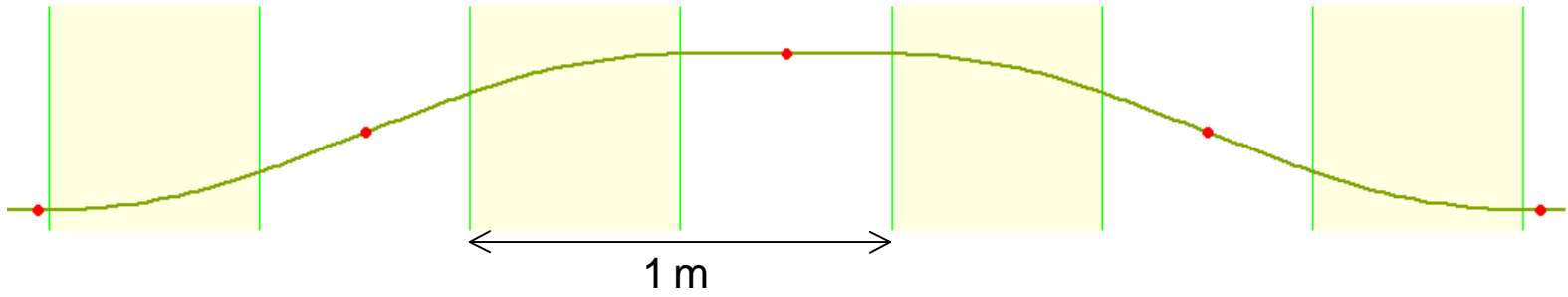
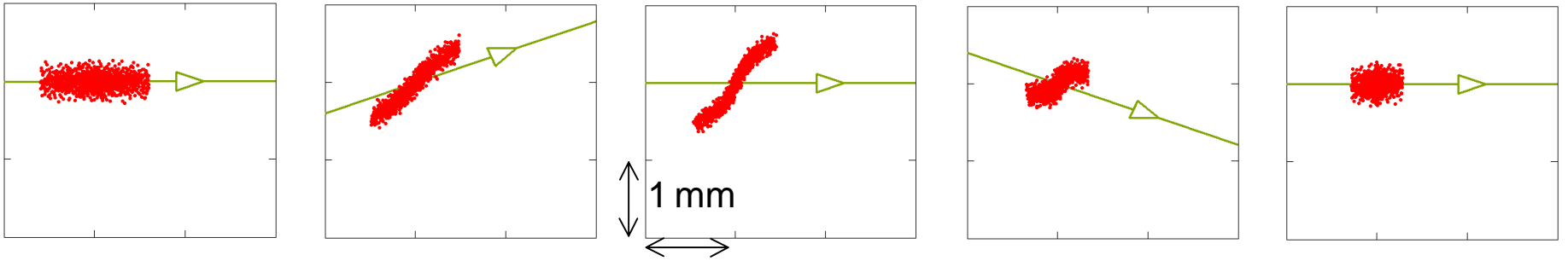


problem: EM fields generated by the bunch

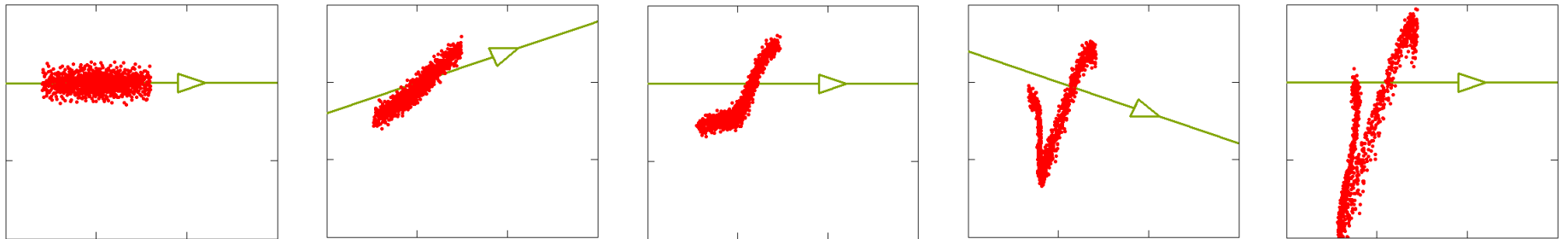
e.g. energy loss of the bunch due to CSR

$\Delta E = f(\text{particle position})$, bump is not closed \rightarrow emittance growth

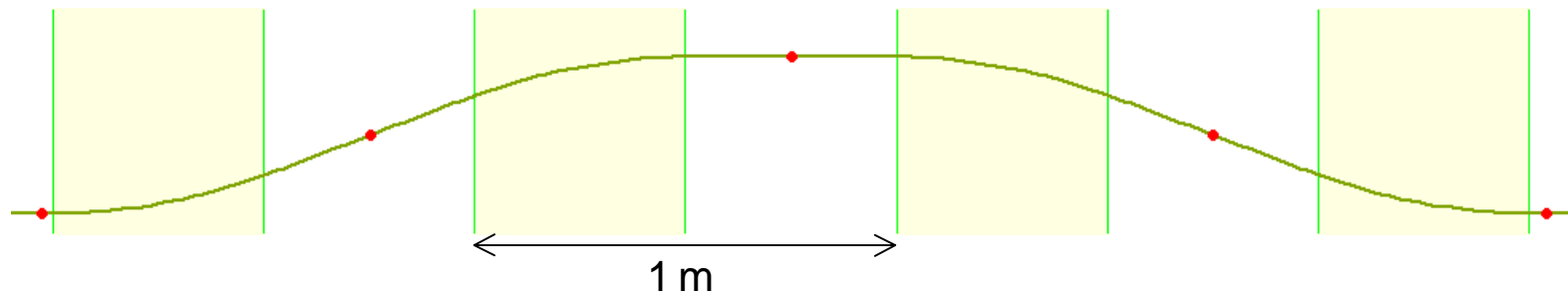
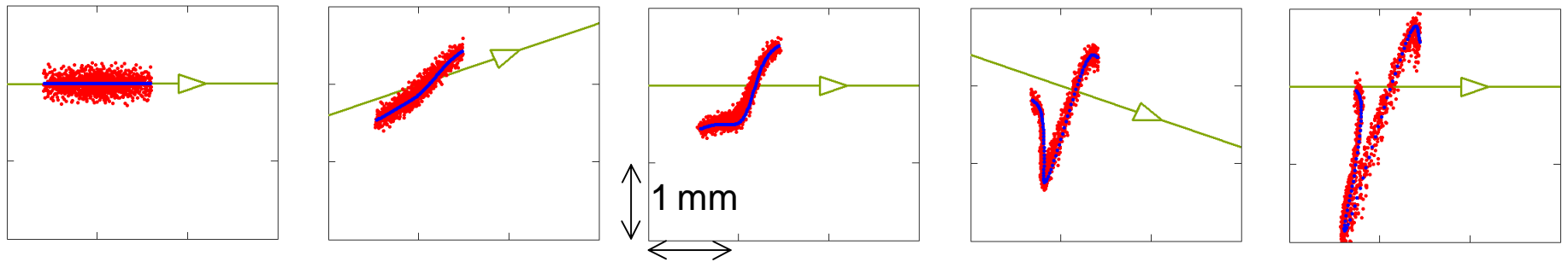
without self-interaction



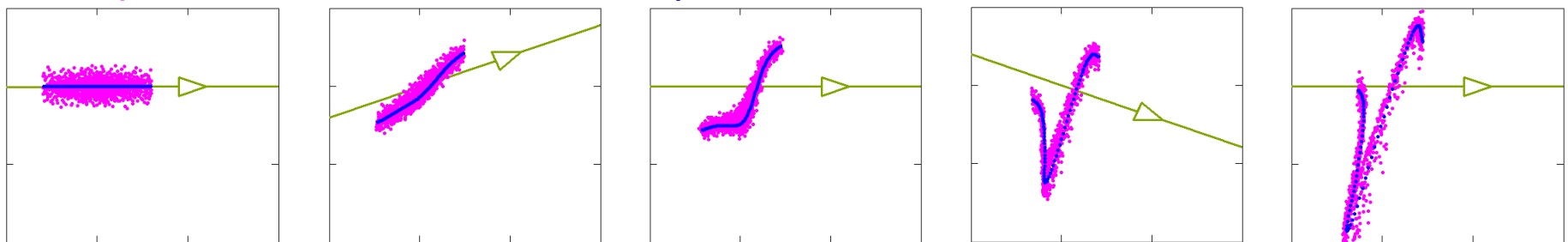
with self-interaction



80 particles, self consistent
1000 particles, self consistent

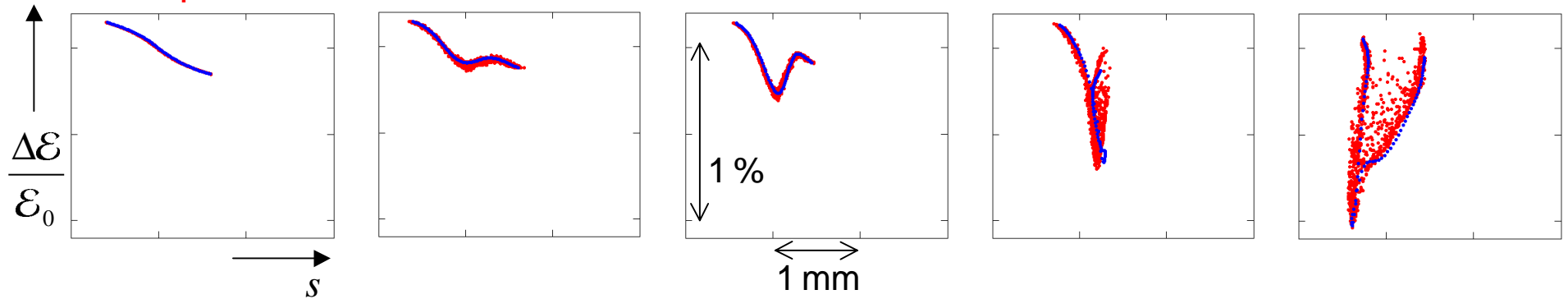


80 particles, self consistent
1000 particles, tracked in field of 80 particles

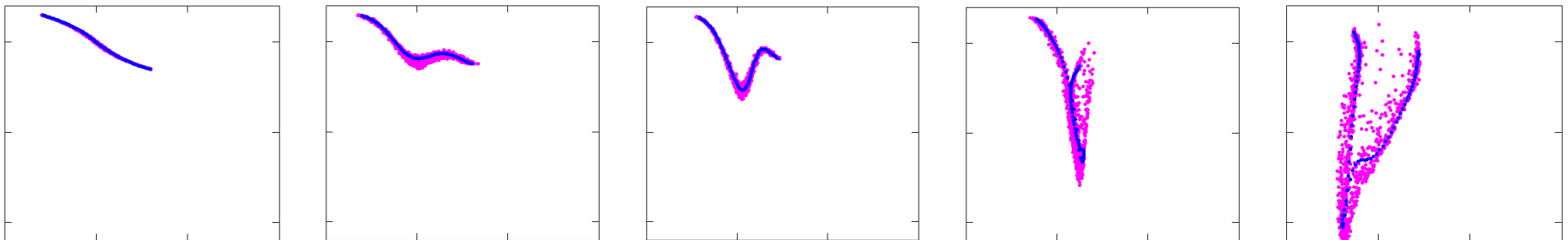


longitudinal phase space

80 particles, self consistent
1000 particles, self consistent



80 particles, self consistent
1000 particles, tracked in field of 80 particles



1.4 1D-Approach

e.g. Elegant
M. Borland

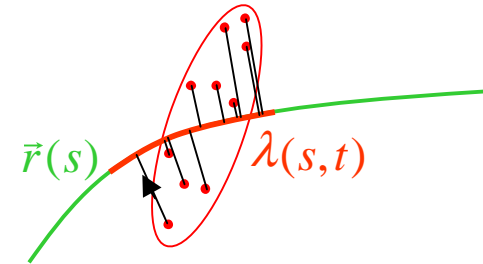
neglect transverse forces

neglect transverse beam dimensions

neglect Φ : $\Delta\hat{\mathcal{E}} = \Delta\mathcal{E}_0 + q_0 \int \text{CSR} dt$

neglect deformation of retarded distribution: $\lambda(s, t') \approx \lambda(s - c_0 t)$

(local rigid bunch approximation)



$$\text{CSR}(s, t) = \int \lambda'(\hat{s} - c_0 t') K(s, \hat{s}) d\hat{s}$$

$$\begin{bmatrix} x \\ x' \\ z \\ \Delta\hat{\mathcal{E}}/\mathcal{E}_0 \end{bmatrix}_{s+\delta s} = R(s, s+\delta s) \begin{bmatrix} x \\ x' \\ z \\ \Delta\hat{\mathcal{E}}/\mathcal{E}_0 \end{bmatrix}_s + \frac{q_0}{\mathcal{E}_0} \frac{\delta s}{c_0} \text{CSR}(s+z, s/c_0) \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix}$$

or: sub-bunch approach

$$\lambda(s, t) \approx \sum q_v \lambda_{(\text{sub})}(s - z_v(t) - c_0 t)$$

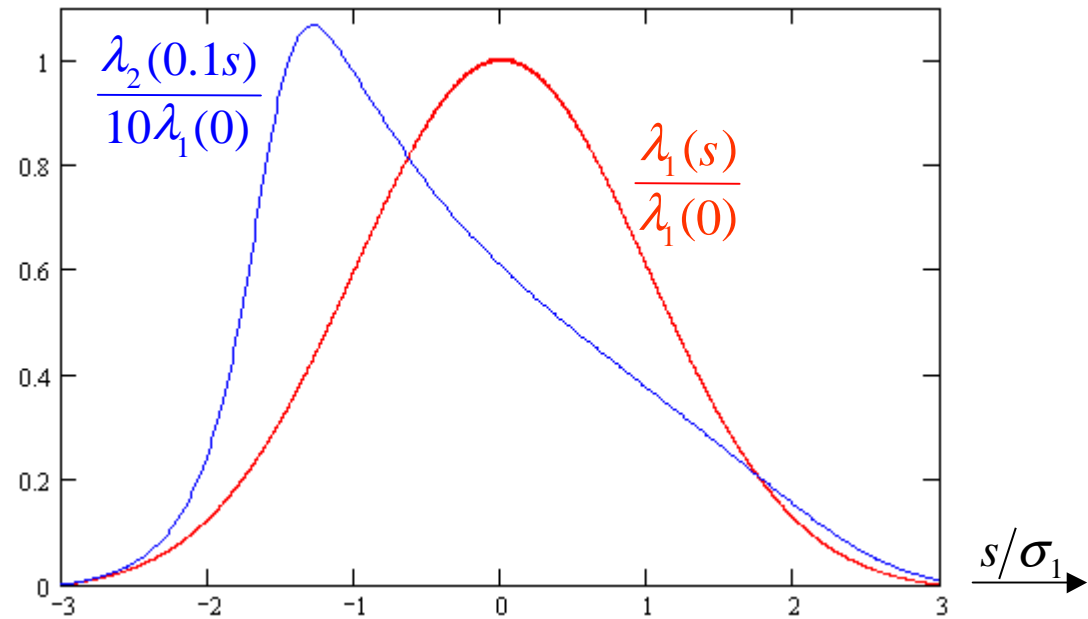
$$\int \lambda_{(\text{sub})}(s) ds = 1$$

$$\text{CSR}_{(\text{sub})}(s, t) = \int \lambda'_{(\text{sub})}(\hat{s} - c_0 t') K(s, \hat{s}) d\hat{s}$$

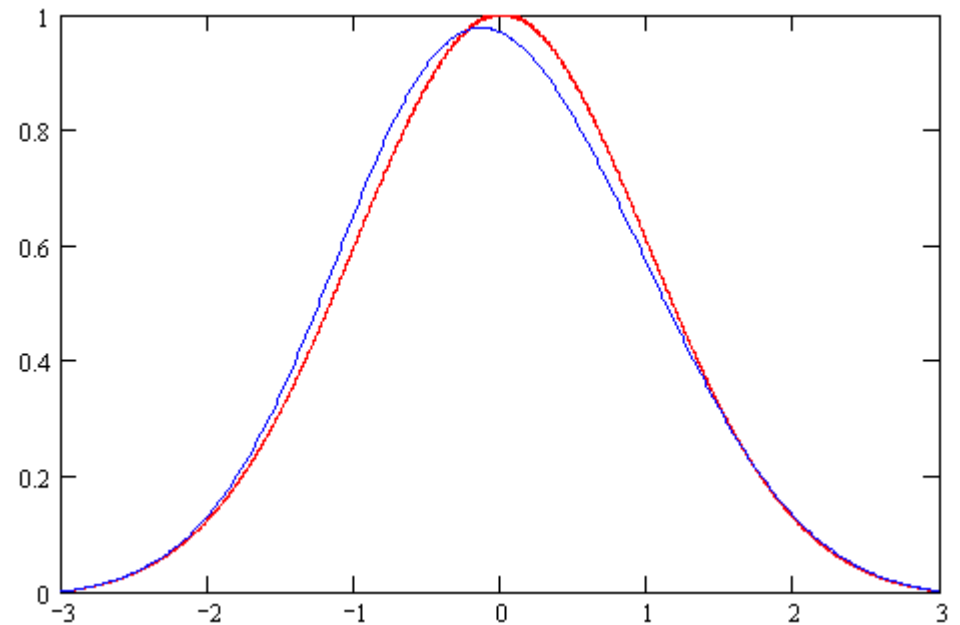
pseudo Green's function

$$\text{CSR}(s, t) \approx \sum q_v \text{CSR}_{(\text{sub})}(s, t + z_v/c_0)$$

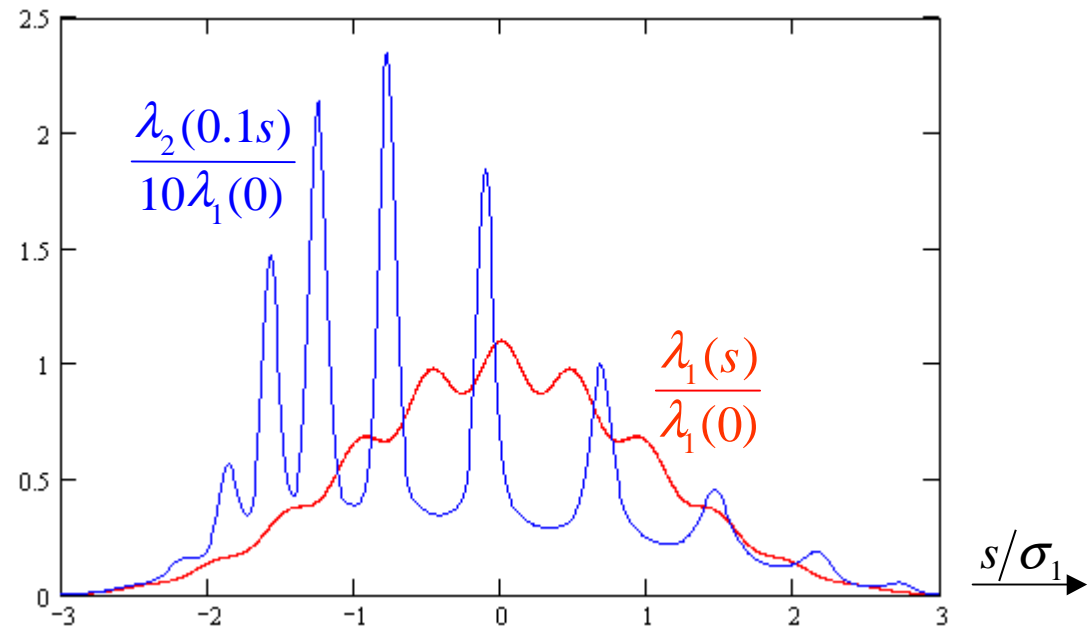
e.g. benchmark example
gaussian bunch
 $q = 1 \text{ nC}$, $\mathcal{E} = 500 \text{ MeV}$, $\sigma_1 = 200 \mu\text{m}$



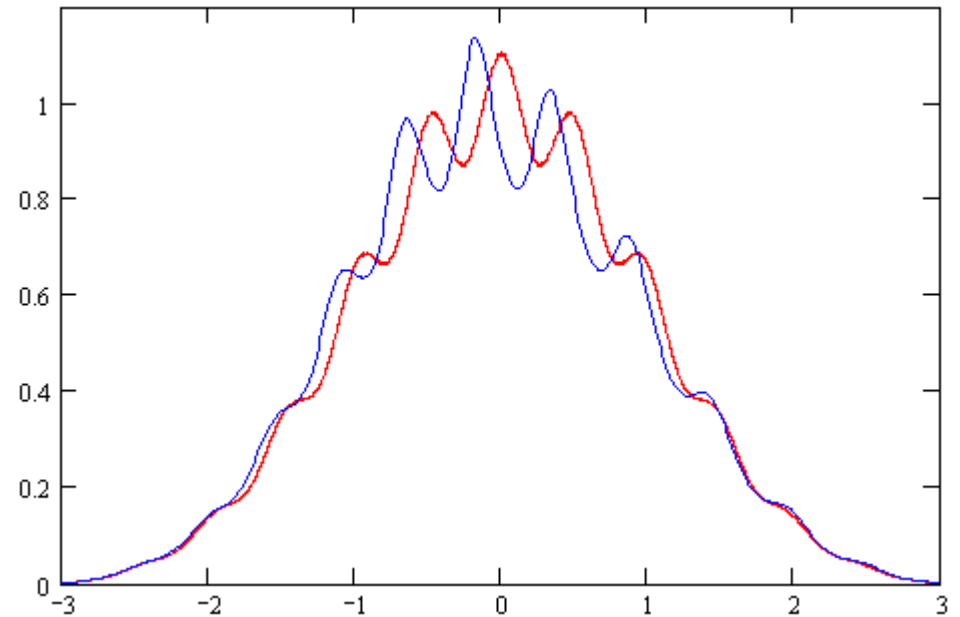
benchmark example
gaussian bunch
 $q = 1 \text{ nC}$, $\mathcal{E} = 5 \text{ GeV}$, $\sigma_1 = 200 \mu\text{m}$



modulated gaussian bunch
 $q = 1 \text{ nC}$, $\mathcal{E} = 500 \text{ MeV}$, $\sigma_1 = 200 \mu\text{m}$
initial modulation = 10 %
initial wavelength = $200 \mu\text{m}$



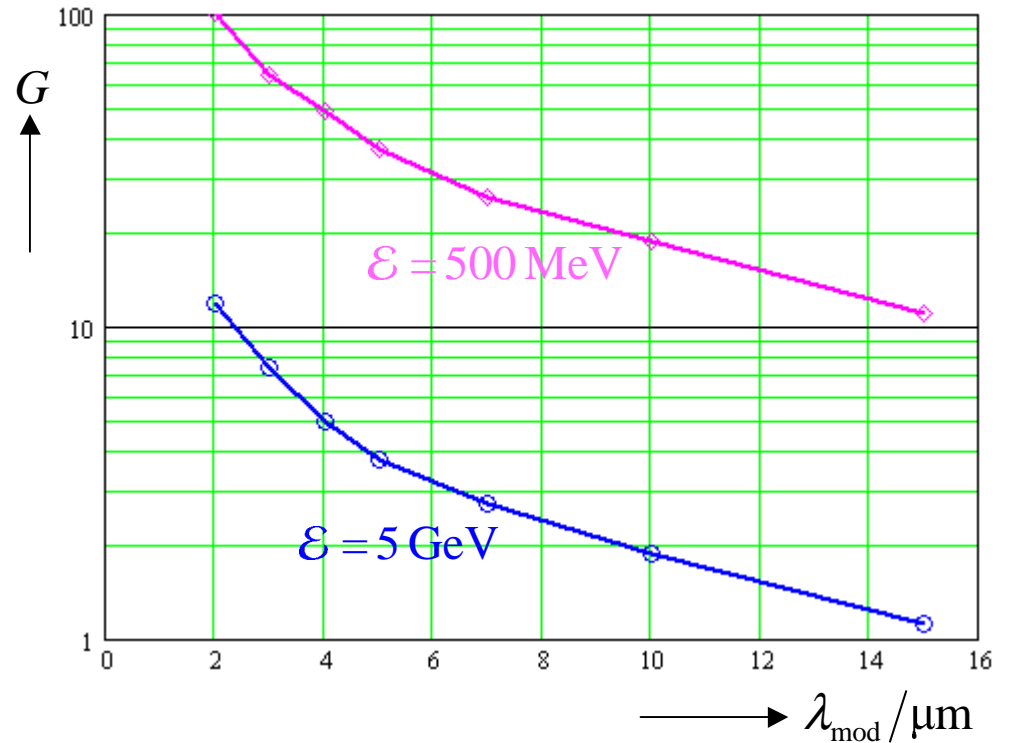
$\mathcal{E} = 5 \text{ GeV}$
same initial distribution



small signal gain

benchmark example

$$q = 1 \text{ nC}, \sigma_1 = 200 \mu\text{m}, \sigma_\varepsilon = 0$$



wavelength after compression

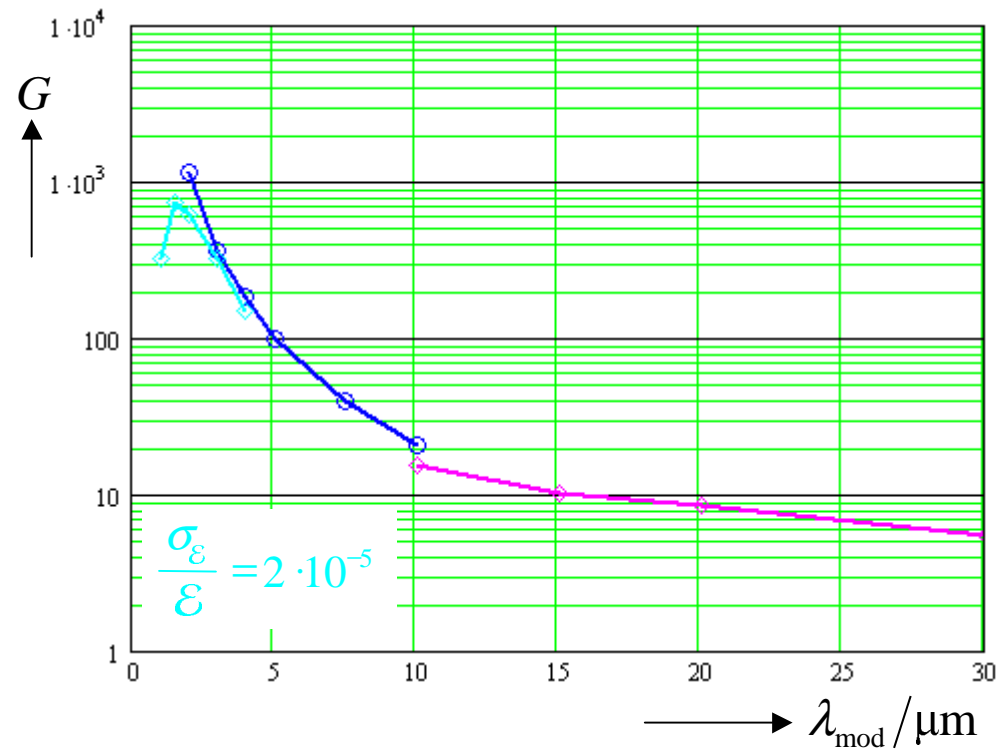
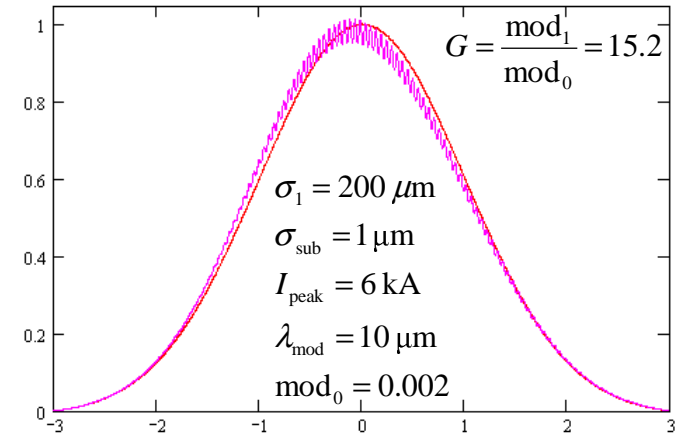
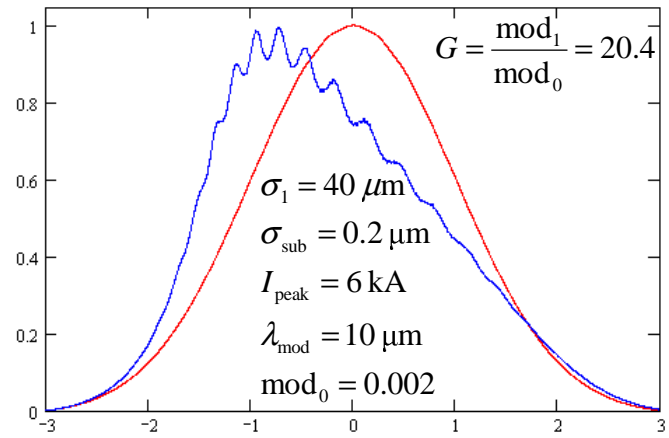
$$\sigma_{\text{sub}} = 0.2 \mu\text{m}$$

maximal gain limited by: uncorrelated energy spread $\lambda_{\text{max}} \approx 2\pi R_{56} \frac{\sigma_\varepsilon}{\varepsilon}$

e.g.: $R_{56} = 2.5 \text{ cm}, \frac{\sigma_\varepsilon}{\varepsilon} = 2 \cdot 10^{-5} \rightarrow \lambda_{\text{max}} \approx 3 \mu\text{m}$

sub bunch length of numerical simulation

small signal gain,
no compression,
 $\mathcal{E} = 5 \text{ GeV}$



2 Effects

2.1 Compression Work

2.2 Surface Impedance

2.3 Longitudinal Field

2.4 Transverse Effects

2.5 Transients

2.6 Shielding

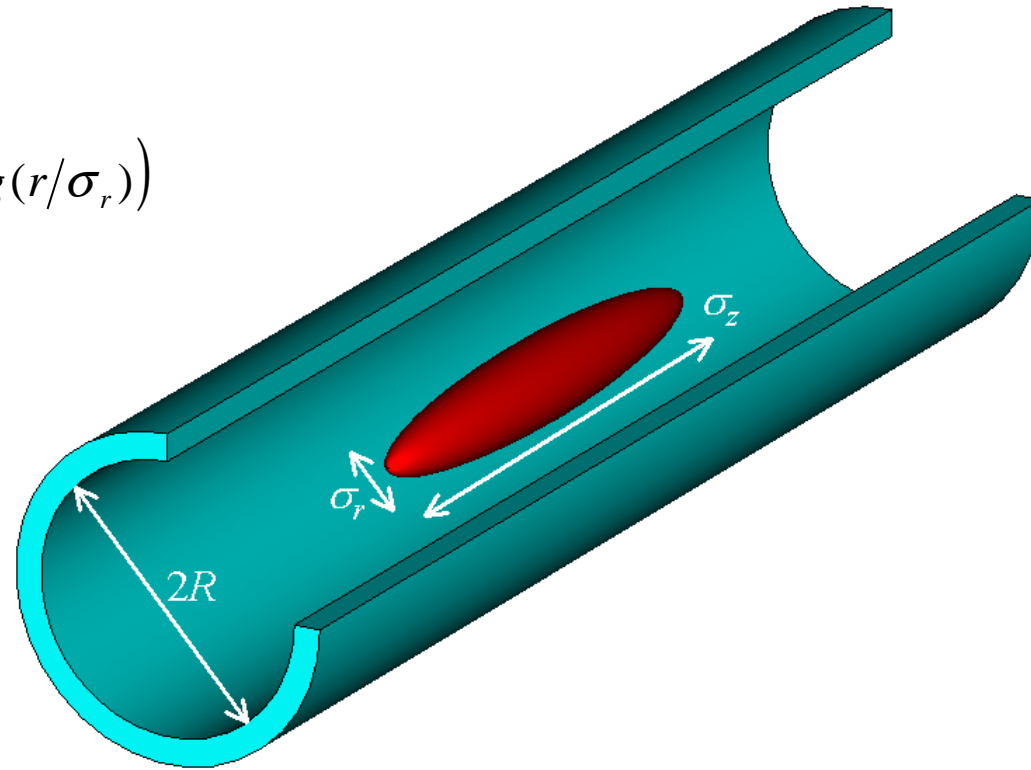
2.1 Compression Work

$$\gamma \rightarrow \infty$$

$$\rho(r, z) = \frac{q}{\sqrt{2\pi\sigma_z\sigma_r^2}} g(z/\sigma_z) g(r/\sigma_r)$$

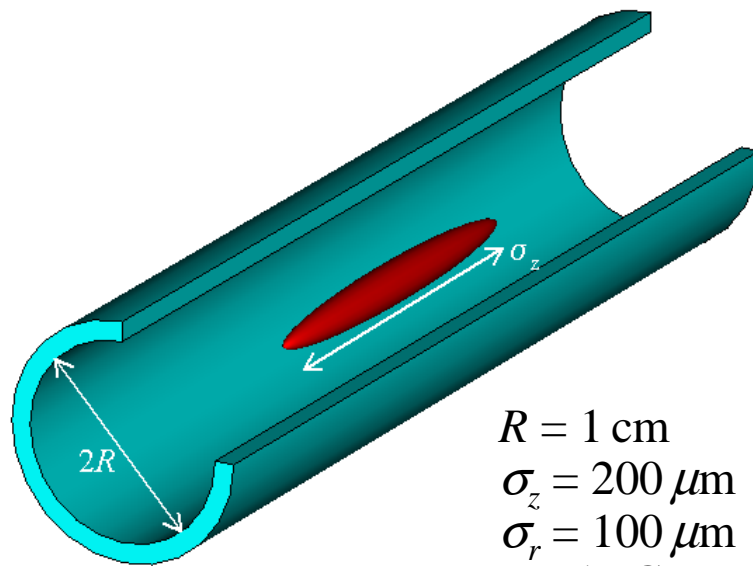
$$E_r(r, z) = \frac{q}{2\pi\epsilon_0 r \sigma_z} g(z/\sigma_z) \left(1 - \sqrt{2\pi} g(r/\sigma_r)\right)$$

$$B_\phi(r, z) = c_0 E_r(r, z)$$



$$\gamma \gg R/\sigma_z$$

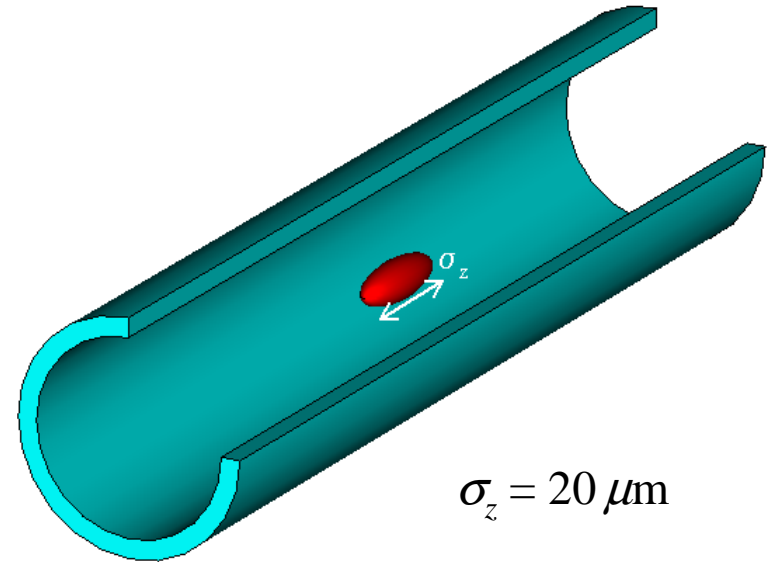
$$W_{tot} = W_e + W_m = \frac{q^2}{4\pi^{3/2}\epsilon_0\sigma_z} \ln\left(\frac{R}{1.5\sigma_r}\right)$$



$$\begin{aligned}
 R &= 1 \text{ cm} \\
 \sigma_z &= 200 \text{ } \mu\text{m} \\
 \sigma_r &= 100 \text{ } \mu\text{m} \\
 q &= 1 \text{ nC}
 \end{aligned}$$

$$W_{\text{tot}} = 0.107 \text{ mJ}$$

compression work



$$\sigma_z = 20 \text{ } \mu\text{m}$$

$$W_{\text{tot}} = 1.065 \text{ mJ}$$

change of potential energy

$$\Delta W_{\text{tot}} = 0.958 \text{ mJ}$$

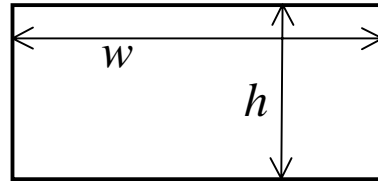
CSR: $R_0 = 10 \text{ m}$, $\sigma_z = 20 \text{ } \mu\text{m}$, $L = 0.5 \text{ m}$

$$P = 375 \text{ kW}$$

$$P L/c_0 = 0.625 \text{ mJ}$$

2.2 Surface Impedance

flat chamber



$$Z \approx \frac{1}{2w} \frac{Z_s}{1 + jk_0 \frac{h Z_s}{2 Z_0}}$$

CSR field:

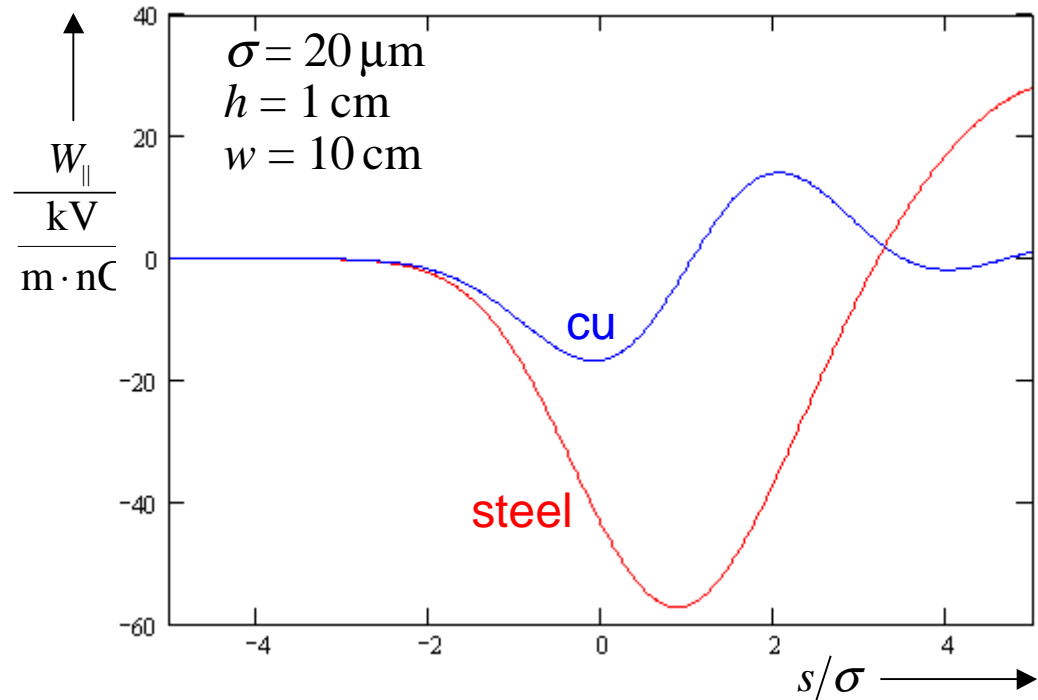
$$E_{\text{csr}} \propto \frac{0.044}{R_0^{2/3} \sigma^{4/3}} \frac{q}{\epsilon_0}$$

$$\sigma = 20 \mu\text{m}$$

$$R_0 = 10 \text{ m}$$

$$\frac{E_{\text{csr}}}{q} \propto 175 \frac{\text{kV}}{\text{m}} \frac{1}{\text{nC}}$$

resistive wall wake:



2.3 Longitudinal Field

$$\begin{aligned}\vec{E} \cdot \vec{e}_{\parallel} &= -\nabla\Phi \cdot \vec{e}_{\parallel} - \frac{\partial}{\partial t} \vec{A} \cdot \vec{e}_{\parallel} \\ &= -(\gamma^{-2} + \beta^2) \nabla\Phi \cdot \vec{e}_{\parallel} - \frac{\partial}{\partial t} \vec{A} \cdot \vec{e}_{\parallel}\end{aligned}\quad \nabla\Phi \cdot \vec{e}_{\parallel} = \frac{1}{v} \left(\frac{d\Phi}{dt} - \frac{\partial\Phi}{\partial t} \right)$$

$$\vec{E} \cdot \vec{e}_{\parallel} = \frac{1}{\gamma^2} \frac{1}{v} \frac{\partial\Phi}{\partial t} - \frac{1}{v} \frac{d\Phi}{dt} + \frac{1}{v} \frac{\partial}{\partial t} (\beta^2\Phi - \vec{A} \cdot \vec{v})$$

“space charge”
potential
“CSR”

Longitudinal CSR Field of a Thin Beam

rigid bunch model: $\lambda(s - vt)$

observer: \vec{r}_o

$\vec{r}(s)$

$\vec{v}_o = v\vec{e}_o$

$$E_{\text{CSR}} = \frac{1}{\beta c_0} \frac{\partial}{\partial t} (\beta^2 \Phi - \vec{A} \cdot \vec{v})$$

$$E_{\text{CSR}} = \frac{-\beta^2}{4\pi\epsilon_0} \int \lambda'(s + \beta R - vt) \frac{1 - \vec{e}_o \cdot \vec{e}}{R} ds$$

$$\vec{e}(s) = \frac{d}{ds} \vec{r}(s)$$

$$R(s) = \|\vec{r}_o - \vec{r}(s)\|$$

Longitudinal Field of a Thin Beam on a Circular Path

$$\vec{E} \cdot \vec{e}_\varphi = \frac{1}{\gamma^2} \frac{1}{v} \frac{\partial \Phi}{\partial t} - \frac{1}{v} \frac{d\Phi}{dt} + E_c G(c_0 t / \sigma)$$

with $E_c = \frac{1}{\sqrt[3]{3}(2\pi)^{3/2}} \frac{1}{R_0^{2/3} \sigma^{4/3}} \frac{q}{\epsilon_0}$

$$G(w) = \sqrt{2\pi} \int_0^\infty \frac{g'(x+w)}{\sqrt[3]{x}} dx$$

coherent radiated power:

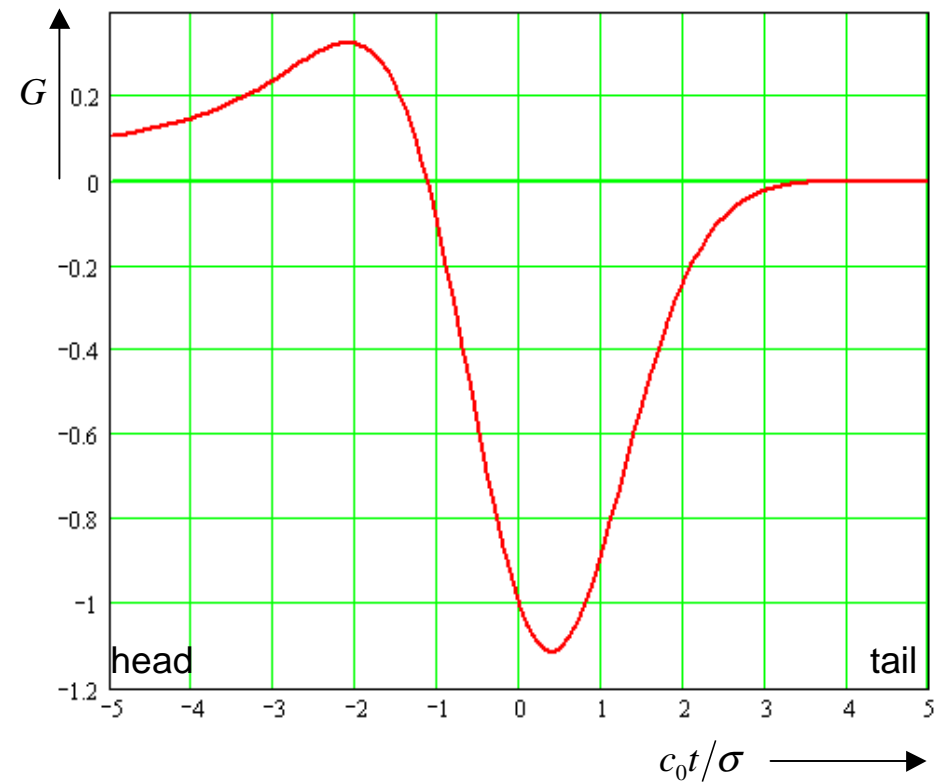
$$P \approx \frac{\Gamma(5/6)}{4\pi^{3/2} \cdot \sqrt[3]{6}} \frac{q^2 c_0}{\epsilon_0} \frac{1}{R_0^{2/3} \sigma^{4/3}}$$

e.g.: $q = 1 \text{ nC}$

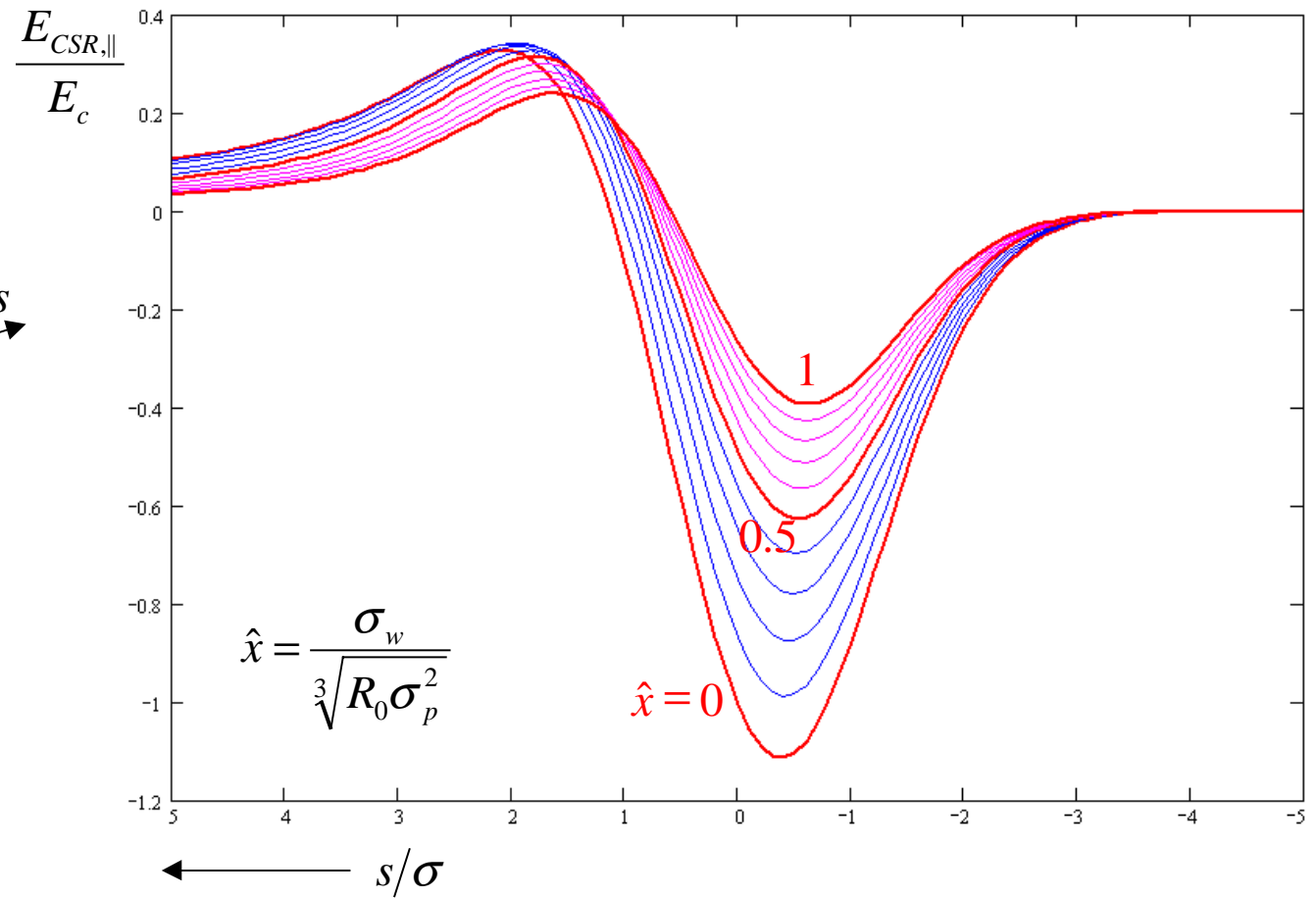
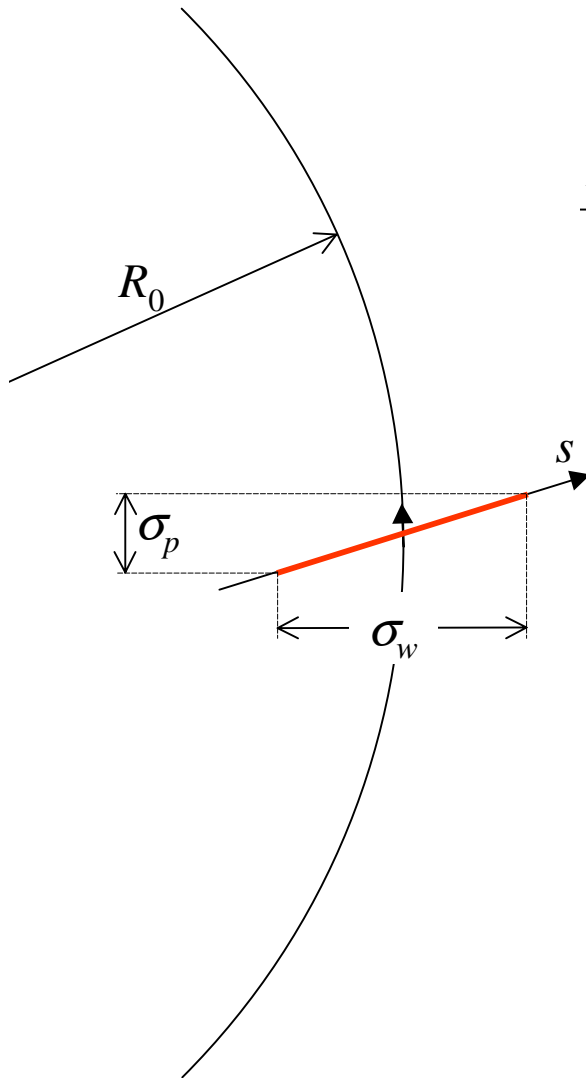
$R_0 = 10 \text{ m}$

$\sigma = 20 \mu\text{m}$

$P \approx 375 \text{ kW}$

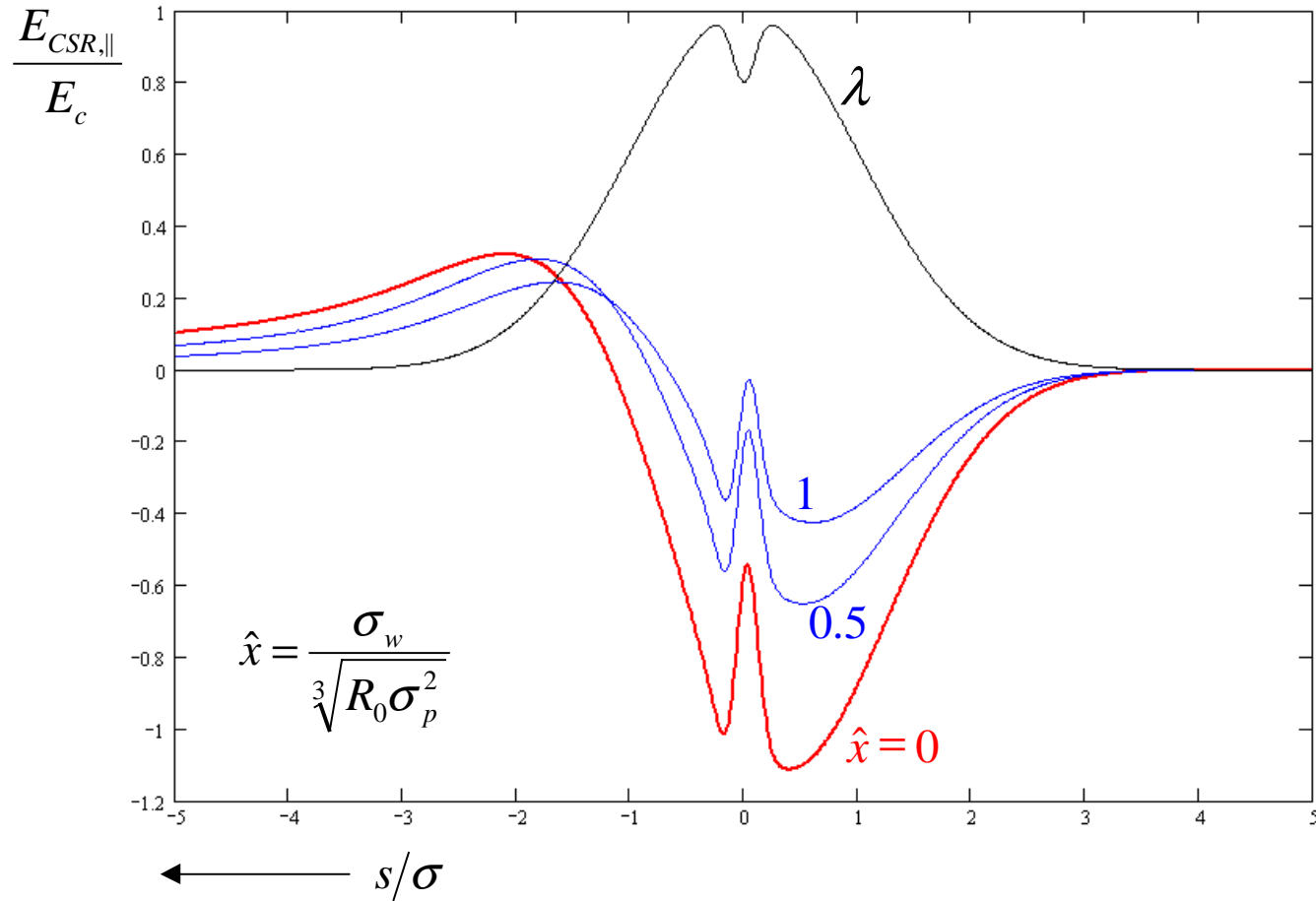


CSR Field of a Tilted Thin Beam



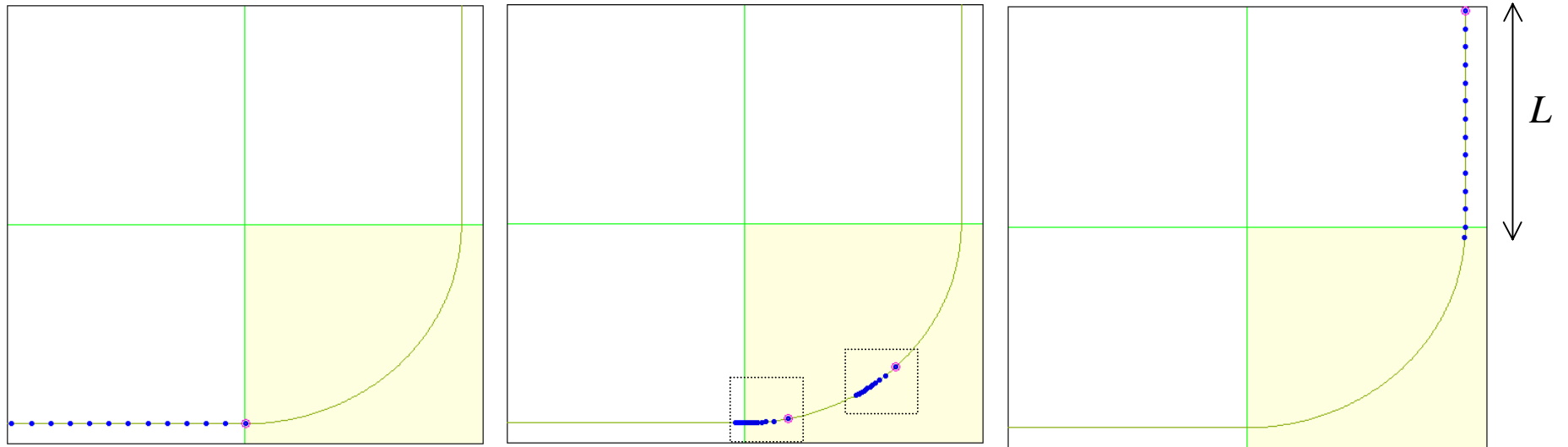
e.g. $R_0 = 10 \text{ m}, \sigma_p = 100 \mu\text{m}, \sigma_w = 2 \text{ mm} \rightarrow \hat{x} = 0.43$

... CSR Field of a Tilted Thin Beam

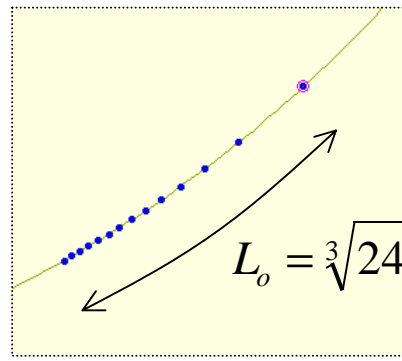
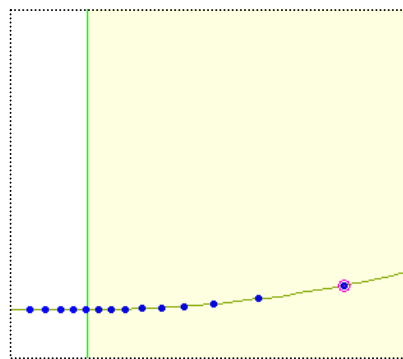


2.5 Transients

retarded particles, seen from the head particle:



$$L = \frac{l}{1-\beta} \approx 2\gamma^2 l$$



$$L_o = \sqrt[3]{24R_0^2 l} \quad \text{for } l \gg R_0/\gamma^3$$

Example:

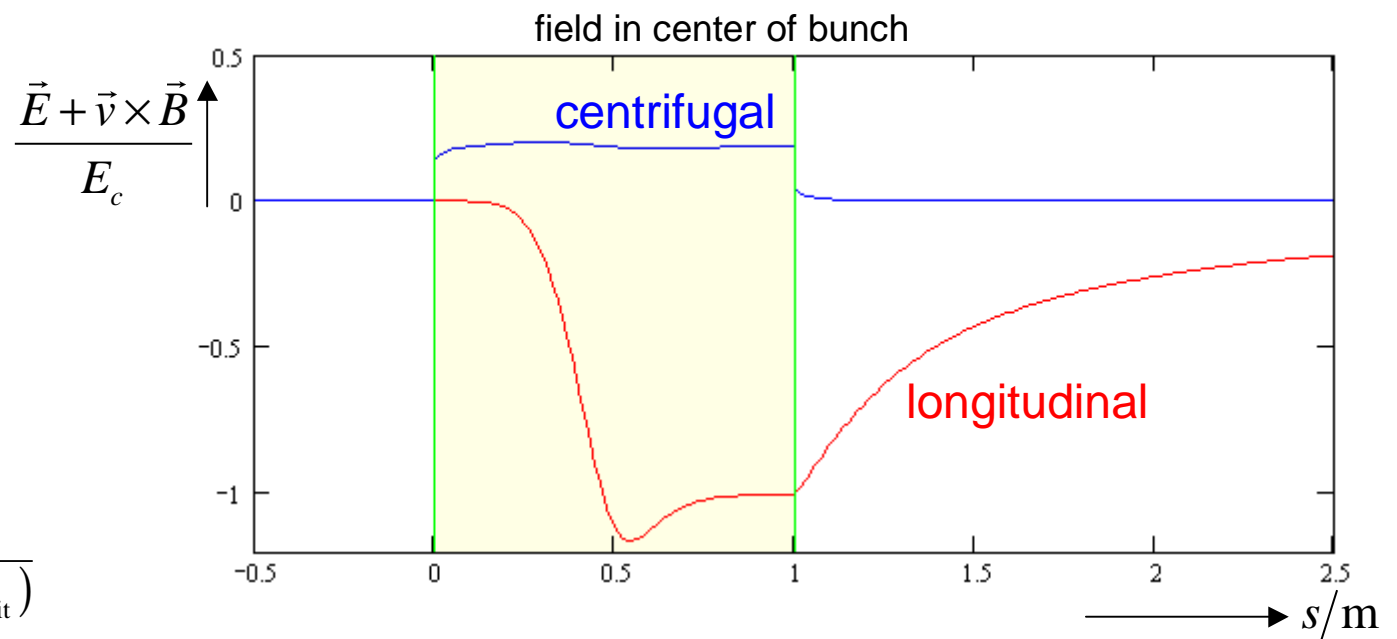
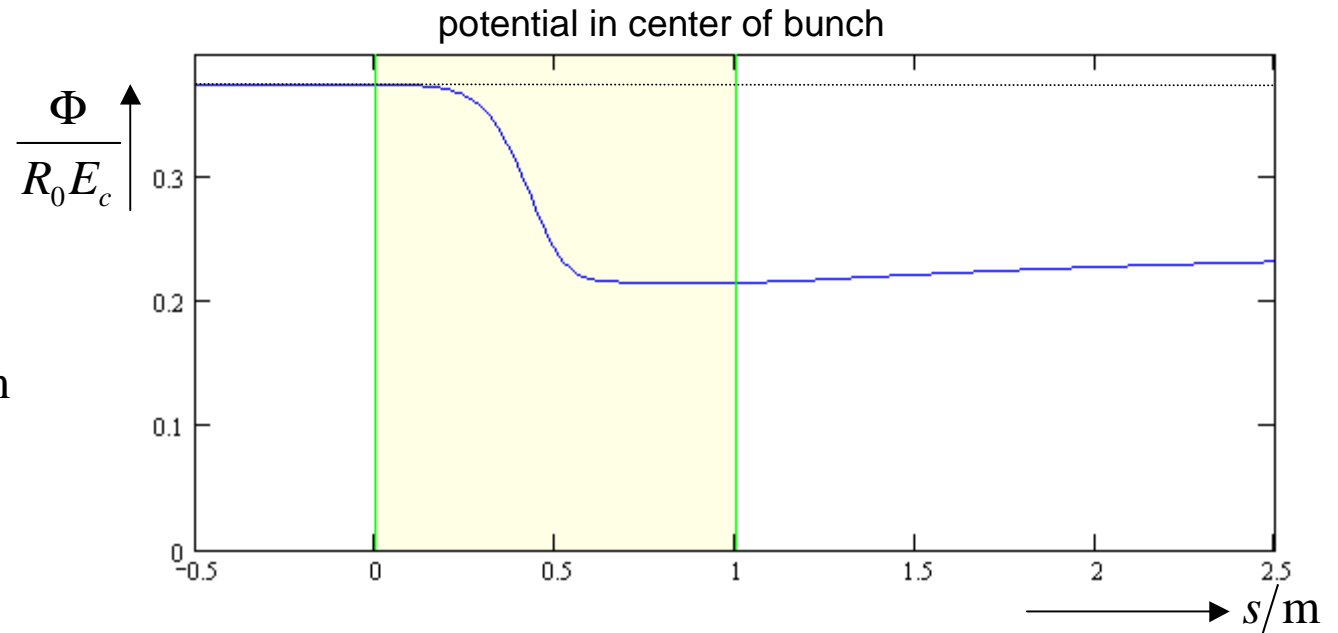
arc radius $R_0 = 10$ m

arc length $L = 1$ m

$\gamma = 10^4$

bunch length $\sigma = 100$ μm

bunch radius 10 μm

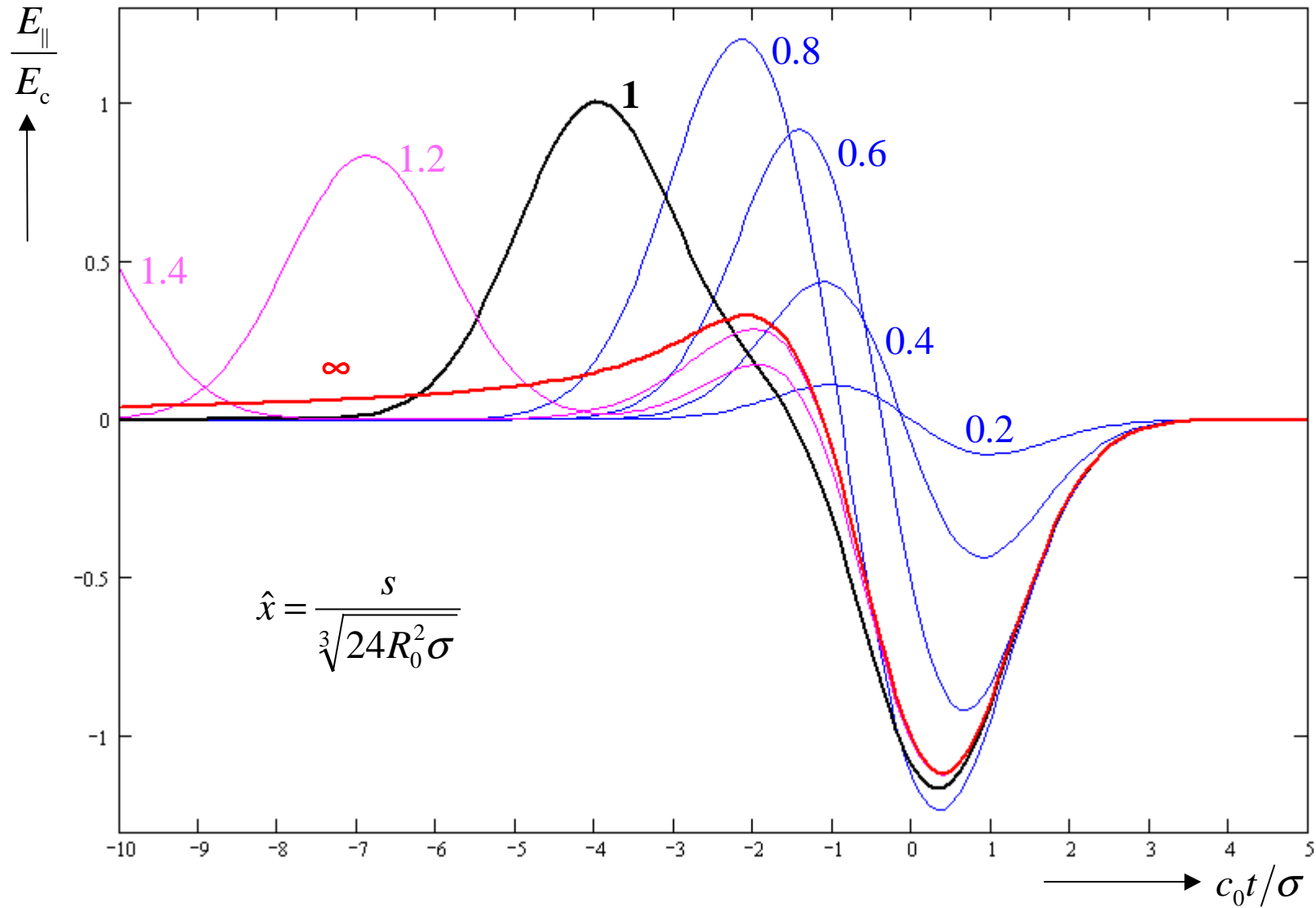


$$L_o = \sqrt[3]{24R_0^2\sigma} \approx 0.621 \text{ m}$$

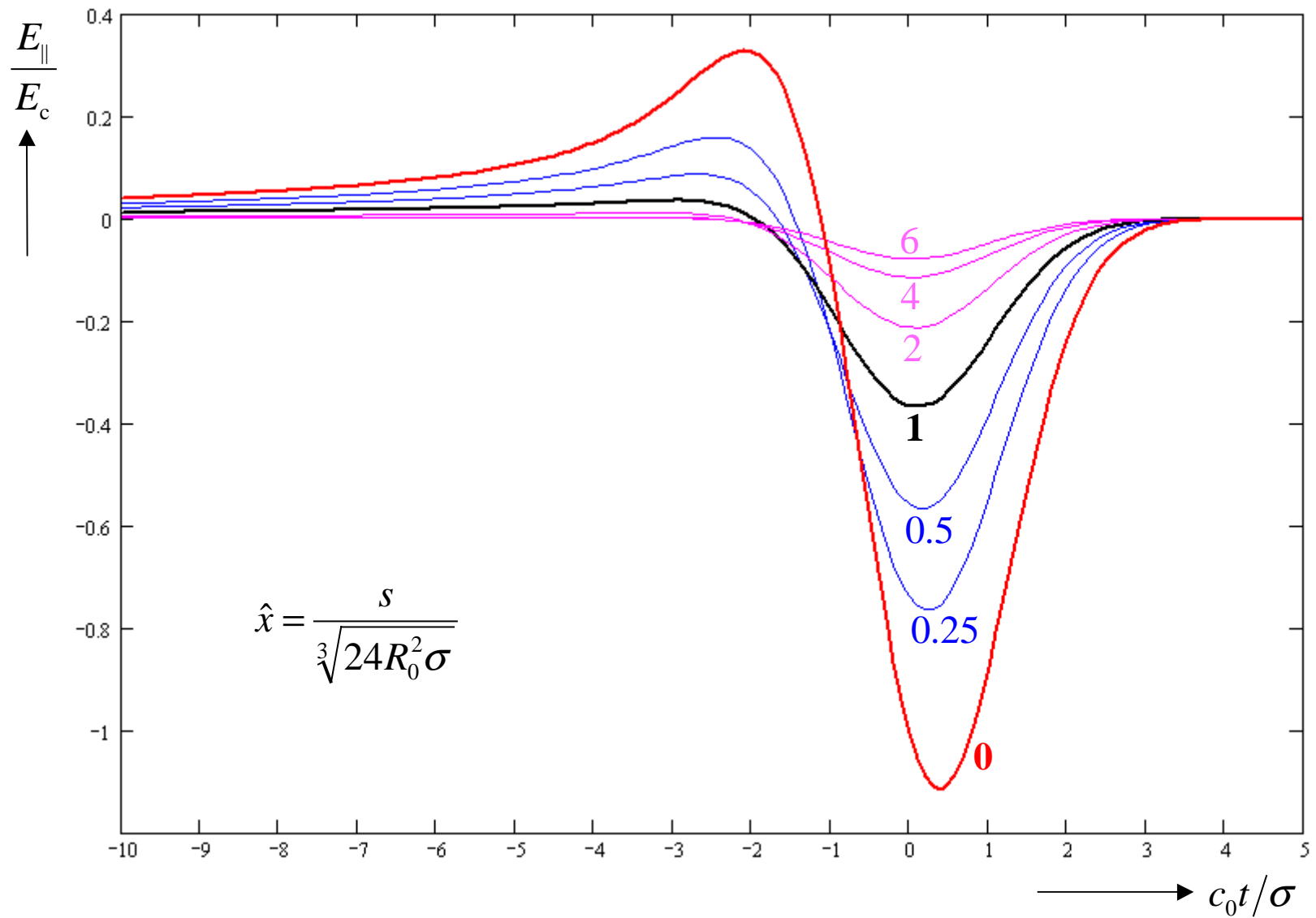
for $L_o \ll s - s_{\text{exit}} \ll L_\gamma$

$$E_{\parallel} \approx \frac{-\lambda}{4\pi\epsilon_0(s - s_{\text{exit}})}$$

Transient Longitudinal Field: Injection

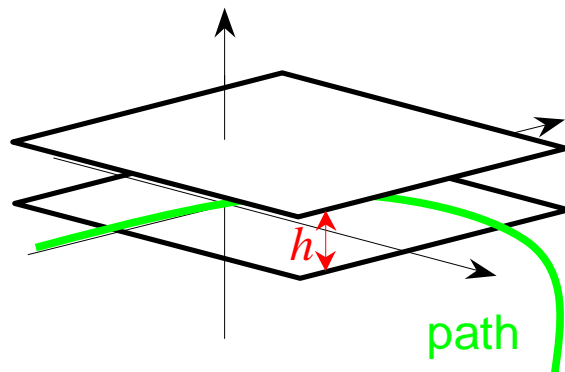


Transient Longitudinal Field: Ejection

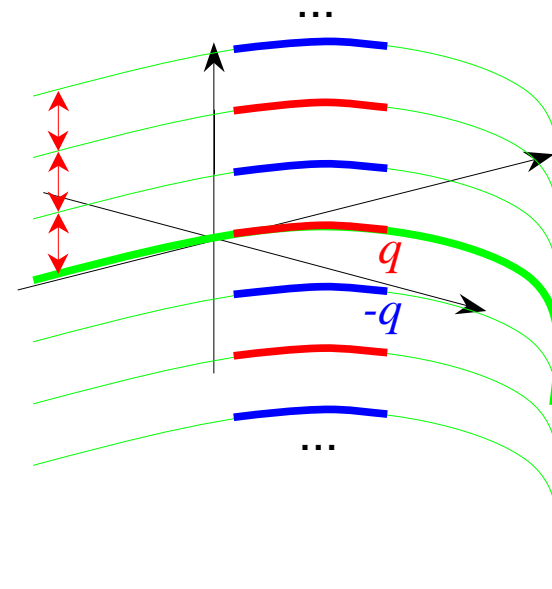


2.6 Shielding

a) shielding by horizontal conducting planes:

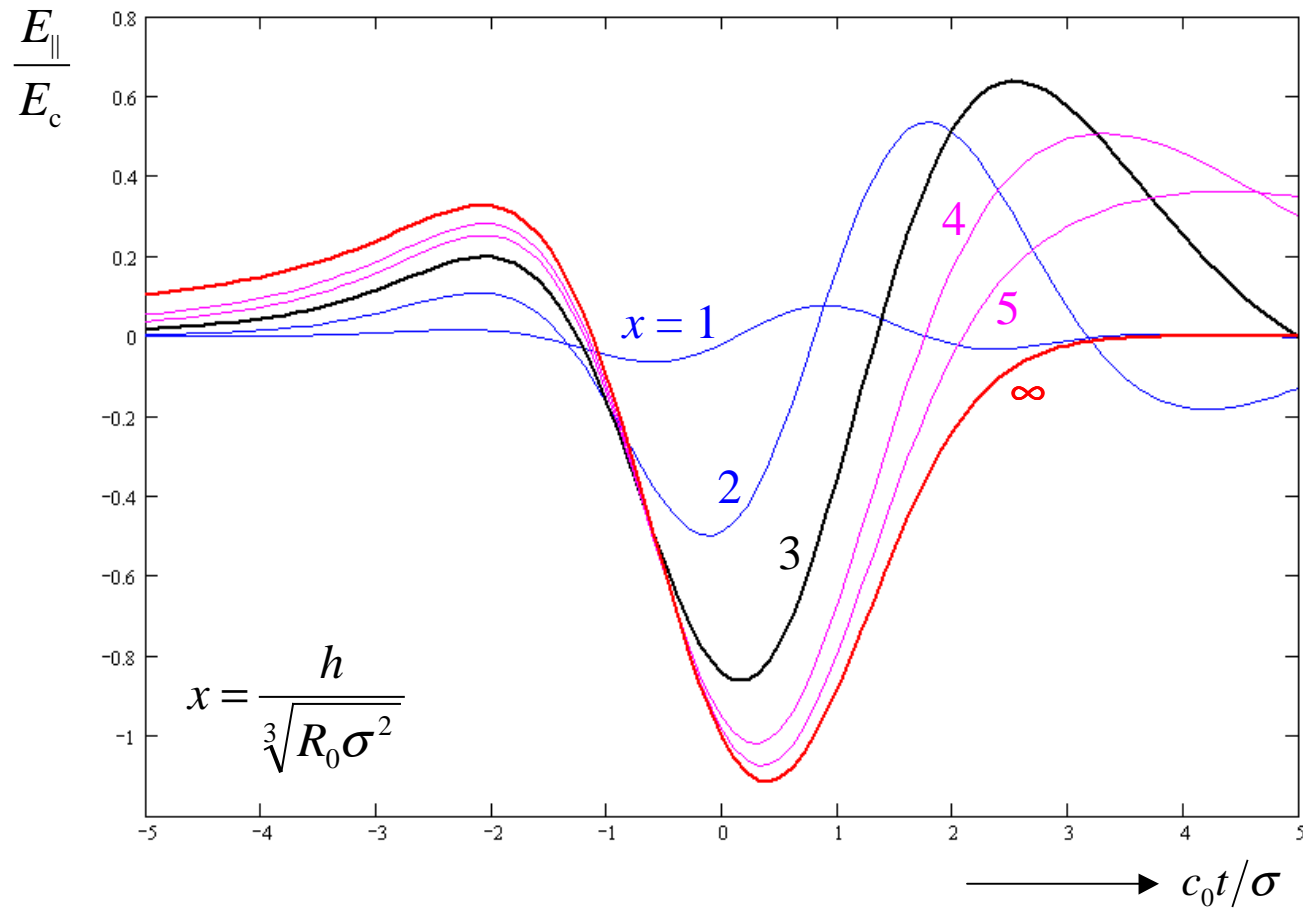


mirror charges:



general path
transient effects

Shielding: Longitudinal Field

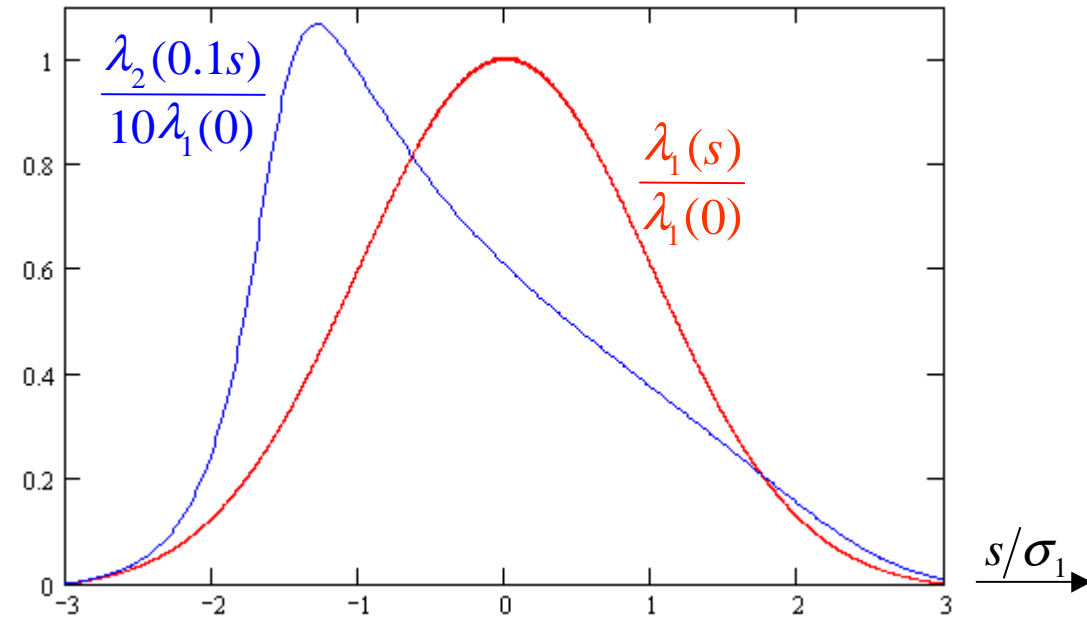


e.g.: $R_0 = 10\text{m}$, $\sigma = 100\mu\text{m}$, $h = 1\text{cm} \rightarrow \hat{h} = 2.1$

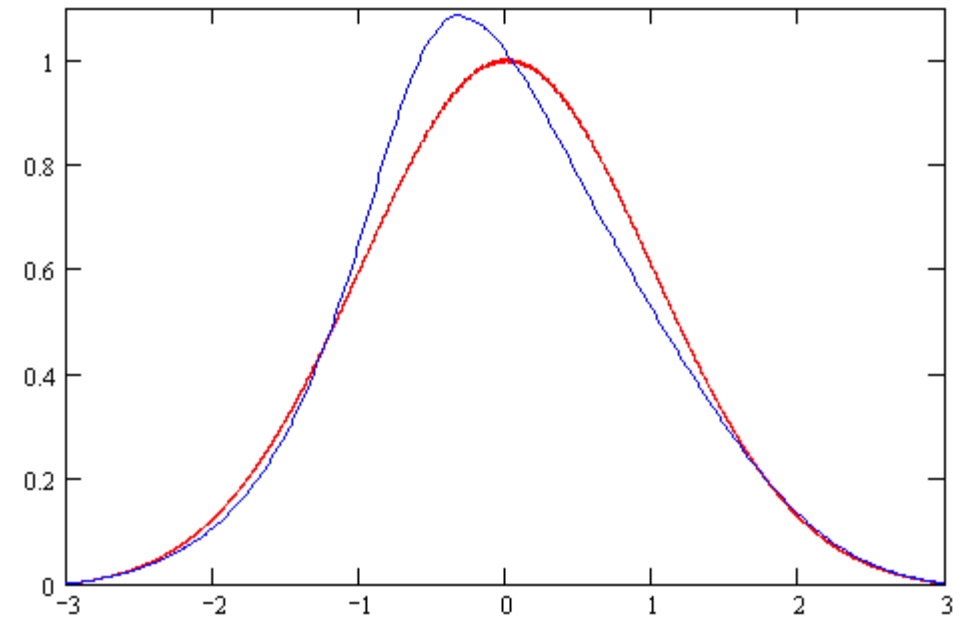
**c) benchmark example
(1D approach):**

gaussian bunch

$$q = 1 \text{ nC}, \mathcal{E} = 500 \text{ MeV}, \sigma_1 = 200 \mu\text{m} \\ \rightarrow \sigma_2 = 20 \mu\text{m}$$



shielding by horizontal
conducting planes, $h = 1 \text{ cm}$



gaussian bunch

$q = 1 \text{ nC}$, $\mathcal{E} = 500 \text{ MeV}$, $\sigma_1 = 200 \mu\text{m}$
 $\rightarrow \sigma_2 = 20 \mu\text{m}$

